## GRAVITOMAGNETICS, THE BASICS OF A SIMPLER APPROACH. H. Ron Harrison* September 2013


#### Abstract

Galileo studied bodies falling under gravity and Tycho Brahe made extensive astronomical observations which led Kepler to formulate his three famous laws of planetary motion. All these observations were of relative motion. This led Newton to propose his theory of gravity which could just as well have been expressed in a form that does not involve the concept of force. The approach in this paper extends the Newtonian theory and the Special Theory of Relativity by including relative velocity by comparison with electromagnetic effects and also from the form of measured data. This enables the non-Newtonian effects of gravity to be calculated in a simpler manner than by use of the General Theory of Relativity (GR). Application to the precession of the perihelion of Mercury and the gravitational deflection of light gives results which agree with observations and are identical to those of GR. It also gives the accepted expression for the Schwarzschild Radius. This approach could be used to determine non-Newtonian variations in the trajectories of satellites. An extra term is then added to the initial basic equation which acts in the direction of the relative velocity. The amended basic equation now predicts a change in the speed of light and derives the accepted result for the Shapiro time delay. It also gives the accepted value for the Last Stable Orbit. Because the extra term is a function of $(v / c)^{3}$ the previously mentioned predictions are not significantly changed *Formerly Senior Lecturer of the Department of Mechanical Engineering and Aeronautics, City University, London.


### 1.0 THE BASICS

### 1.1 Newtonian Gravity

Galileo studied bodies falling to Earth under gravity and concluded that all bodies fell with the same acceleration independent of size and material. Tycho Brahe made extensive astronomical observations which led Kepler to formulate his three famous laws of planetary motion relative to the Sun. All of these observations were of relative motion but the mass of one body was, in each case, much greater than that of the other. These led Newton to propose his theory of gravity using the concept of force and yielding an equation which gives the acceleration of a body relative to the centre of mass. He could just as well have presented it in the form

$$
\begin{equation*}
a_{B / A}=-\frac{G\left(m_{A}+m_{B}\right)}{r_{B / A}^{2}} \tag{1}
\end{equation*}
$$

without invoking the concept of force, and only requiring one definition of mass.
That is, the acceleration of body $B$ relative to $A$, in the radial direction, is proportional to the sum of their masses and inversely proportional to the square of their separation. $G$ is the gravitational constant.

### 1.2 Gravitomagnetics

It is now proposed that equation (1) be extended to include the relative velocity. The axioms are.
(a) It is assumed that in mass-free space light travels in straight lines. This defines a non-rotating frame of reference.
(b) Because all motion is relative there are no other restrictions on the frame of reference.
(c) Gravity propagates at the same speed as light.
(d) Mass, or rest mass, is simply the quantity of matter and is regarded as constant. It could be a count of the number of basic particles

The initial proposed equation is based on comparisons with electromagnetics. This equation gives results which agree with the measured results of the precession of the perihelion of Mercury and with the deflection of light grazing the Sun. Also it gives the correct definition for the Schwarzschild Radius. However, it suggests that the speed of light is constant. As a result it does not predict the Shapiro Time Delay. An extra term is then added which gives agreement with the time delay and also generates the accepted value for the Last Stable Orbit. See equations (2a), (3a), (4a) and (5a).

$$
\begin{align*}
& \boldsymbol{a}=-\frac{K}{r^{2}}\left(1-\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K}{r^{2} c^{2}} \boldsymbol{v} \times\left(\boldsymbol{v} \times \boldsymbol{e}_{r}\right)  \tag{2}\\
& \boldsymbol{a}=-\frac{K}{r^{2}}\left(1+\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K v_{r}}{r^{2} c^{2}} \boldsymbol{v} \tag{3}
\end{align*}
$$

where $\boldsymbol{a}=$ acceleration of body B relative to body A , $\boldsymbol{v}=$ the relative velocity, $\quad r=$ the separation and $\boldsymbol{e}_{r}=$ the unit vector from body A to body B. Also $c=$ speed of light, $K=G\left(m_{A}+m_{B}\right)$ and $v_{\mathrm{r}}$ is the radial component of velocity. Note that $G$ is a constant which could be incorporated into the definition of the quantity of matter. These equations reduce to equation (1) when $v \ll c$.

The equation can also be written in terms of the Newtonian part plus the gravitomagnetic part

$$
\begin{equation*}
\boldsymbol{a}=\boldsymbol{a}_{N}+\boldsymbol{a}_{V}=-\frac{K}{r^{2}} \boldsymbol{e}_{r}+\frac{K}{r^{2}}\left(\frac{v}{c}\right)^{2} \boldsymbol{e}_{2 \phi} \tag{4}
\end{equation*}
$$

where $\phi$ is the angle between the velocity and the radius.
It should be noted that the velocities of the individual bodies do not appear in these equations, only the relative velocity. For two isolated bodies the relative motion is the only measureable value.

A convenient definition of force is

$$
\begin{equation*}
\boldsymbol{P}=\mu \boldsymbol{a}=-\frac{G m_{A} m_{B}}{r^{2}}\left(1-\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{\boldsymbol{r}}+\frac{2 G m_{A} m_{B}}{r^{2} c^{2}} \boldsymbol{v} \times\left(\boldsymbol{v} \times \boldsymbol{e}_{\boldsymbol{r}}\right) \tag{5}
\end{equation*}
$$

where $\quad \mu=m_{A} m_{B} /\left(m_{A}+m_{B}\right)$, the reduced mass.
By definition of the centre of mass (or the centre of momentum) the total momentum is zero with reference to the centre of mass. It is now proposed that the motion of the centre of mass of two bodies is not affected by collision. From this it follows that for a group of particles the motion of the centre of mass is unaffected by internal impacts.

The relative acceleration is only radial when the relative velocity is either radial or tangential. In general the moment of momentum can be shown to be a function of the relative position. So, for an elliptic orbit it remains within bounds. Conservation of moment of momentum results from Newton's third law, but this is not true for electromagnetic or gravitomagnetic reactions. So this result should not be a surprise.

General inferences from equation (2) .
Reverts to Newtonian form when $v \ll c$.
The second term of (2) is normal to the velocity.
If $v=c$ the first term of (2) vanishes so that there is no change of speed.
Moment of velocity ( or moment of momentum per total quantity of matter) is not conserved. It is shown to be a function of $r$.
The equivalence of inertial mass to gravitational mass does not arise.

### 1.3 Modified Equations.

An extra term is added in the direction of the relative velocity. This will affect the speed of light but not its deflection. As the term is a function of $c^{3}$ it only has a very small effect on the motion of large bodies in Solar orbits.

The new equation is

$$
\begin{equation*}
\boldsymbol{a}=\boldsymbol{a}_{N}+\boldsymbol{a}_{V}+\boldsymbol{a}_{T}=-\frac{K}{r^{2}} \boldsymbol{e}_{r}+\frac{K}{r^{2}}\left(\frac{v}{c}\right)^{2} \boldsymbol{e}_{2 \phi}+\frac{K}{r^{2}} \frac{v_{r}}{c}\left(\frac{v}{c}\right)^{2} \boldsymbol{t} \tag{4a}
\end{equation*}
$$

where $t$ is the unit vector in the direction of $\boldsymbol{v}$ and $\phi$ is the angle between the velocity and the radius.
Note that $\boldsymbol{t}=\frac{\boldsymbol{v}}{|v|}$. Also note that $c$ is the speed of light in a gravity free vacuum.

This equation will be considered to be the basic for Post Newtonian Gravity. Justification will come from agreement with verified experimental data.


Figure 1

Equation (4a) may be rewritten as

$$
\begin{equation*}
\boldsymbol{a}=-\frac{K}{r^{2}}\left(1-\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K}{r^{2} c^{2}} \boldsymbol{v} \times\left(\boldsymbol{v} \times \boldsymbol{e}_{r}\right)+\frac{K}{r^{2}} \frac{v_{r}}{c}\left(\frac{v}{c}\right)^{2} \boldsymbol{t} \tag{2a}
\end{equation*}
$$

or

$$
\begin{equation*}
\boldsymbol{a}=-\frac{K}{r^{2}}\left(1+\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K v_{r}}{r^{2} c^{2}} \boldsymbol{v}+\frac{K}{r^{2}} \frac{v_{r}}{c}\left(\frac{v}{c}\right)^{2} \boldsymbol{t} \tag{3a}
\end{equation*}
$$

From which it is seen that the additional term is negligible when $(v / c)^{3}$ is small compared to unity
Again, noting that $\boldsymbol{t}=\boldsymbol{v} /|v|$ means that (3a) may be written as

$$
\boldsymbol{a}=-\frac{K}{r^{2}}\left(1+\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K v_{r}}{r^{2} c^{2}}\left[1+\frac{|v|}{2 c}\right] \boldsymbol{v}
$$

or $\quad \boldsymbol{a}=-\frac{K}{r^{2}}\left(1+\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K v_{r}}{r^{2} c^{2}} Q \boldsymbol{v} \quad$. where $\quad Q=\left[1+\frac{|v|}{2 c}\right]$

And
$\boldsymbol{P}=\mu \boldsymbol{a}=-\frac{G m_{A} m_{B}}{r^{2}}\left(1-\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 G m_{A} m_{B}}{r^{2} c^{2}} \boldsymbol{v} \times\left(\boldsymbol{v} \times \boldsymbol{e}_{r}\right)+\frac{G m_{A} m_{B}}{r^{2} c^{2}} \frac{v_{r}}{c}\left(\frac{v}{c}\right)^{2} \boldsymbol{t}$
where $\quad \mu=m_{A} m_{B} /\left(m_{A}+m_{B}\right)$, the reduced mass.

Equations (2-5) will account for the precession of the perihelion of Mercury and will predict the observed value for the deflection of light grazing the Sun. These results were heralded as confirmation of Einstein's General Theory of Relativity. They also give the accepted value for the Schwarzschild Radius.

Equations ( $2 \mathrm{a}-5 \mathrm{a}$ ) will generate the same results as mentioned above but will now give the accepted value for the Last Stable Orbit. It also agrees with the measured value of the Shapiro Time delay because the speed of light is now affected by gravity.

### 2.0 APPLICATION TO TWO MASS PROBLEM

### 2.1 Polar Coordinates

The following development is based on the conventional treatment of the two body gravitational problem. For the dynamics of bodies in Solar orbits the modified equations are not required. Here, $r$ is the separation and $\boldsymbol{e}_{r}$ is the unit vector in the direction of body 2 as seen from body $1 . \theta$ is the orientation of
the unit vector with respect to the 'fixed' stars and $\boldsymbol{e}_{\boldsymbol{\theta}}$ is the unit vector normal to $\boldsymbol{e}_{\boldsymbol{r}}$ in the plane of the motion.
$\begin{array}{ll}\text { Now } & \boldsymbol{a}=\left(\ddot{r}-r \dot{\theta}^{2}\right) \boldsymbol{e}_{r}+(r \ddot{\theta}+2 \dot{r} \dot{\boldsymbol{\theta}}) \boldsymbol{e}_{\theta} \\ \text { and } & \boldsymbol{v}=r \boldsymbol{e}_{r}+(r \dot{\theta}) \boldsymbol{e}_{\theta}\end{array}$

Equation (3a) can now be expressed in component form

$$
\begin{align*}
& \ddot{r}-r \dot{\theta}^{2}=-\frac{K}{r^{2}}\left[1+\left(\frac{v}{c}\right)^{2}\right]+\frac{2 K v_{r}^{2}}{r^{2} c^{2}} Q  \tag{6}\\
& r \ddot{\theta}+2 \dot{r} \dot{\theta}=\frac{1}{r} \frac{d}{d t}\left(r^{2} \dot{\theta}\right)=\frac{2 K v_{r} v_{\theta}}{r^{2} c^{2}} Q \tag{7}
\end{align*}
$$

but for low values of $(v / c)$ the term $Q$ will be taken to be unity.
Define $h=r(r \dot{\theta})$, the moment of momentum per reduced mass, and $u=1 / r$. so that $h=\dot{\theta} / u^{2}$ thus

$$
\begin{array}{r}
\dot{r}=-\frac{1}{u^{2}} \frac{d u}{d \theta} \frac{d \theta}{d t}=-h \frac{d u}{d \theta} \\
\ddot{r}=-u^{2} h^{2} \frac{d^{2} u}{d \theta^{2}}-\dot{h} \frac{d u}{d \theta}
\end{array}
$$

and

$$
\ddot{\theta}=-\frac{2 h \dot{r}}{r^{3}}+\frac{\dot{h}}{r^{2}}=2 u^{3} h^{2} \frac{d u}{d \theta}+u^{2} \dot{h}
$$

Equations $(6,7)$ may now be written

$$
\begin{equation*}
\frac{d^{2} u}{d \theta^{2}}+\frac{d u}{d \theta} \frac{\dot{h}}{u^{2} h^{2}}+u=\frac{K}{h^{2}}+\frac{K u^{2}}{c^{2}}-\frac{K}{c^{2}}\left(\frac{d u}{d \theta}\right)^{2} \tag{8}
\end{equation*}
$$

and

$$
\begin{equation*}
u \dot{h}=-\frac{2 K}{c^{2}} u^{3} h^{2} \frac{d u}{d \theta} \tag{9}
\end{equation*}
$$

Since $\dot{h}=\frac{d h}{d \theta} \dot{\theta}=\frac{d h}{d \theta} u^{2} h \quad$ combining with equation (9) gives, $\quad \frac{d h}{h}=\left(\frac{2 K d i}{c^{2}}\right)$
Integrating equation (10) leads to

$$
h=h_{0} e^{-2 K\left(u-u_{0}\right) / c^{2}}
$$

Therefore, for small variations $\quad h^{2} \approx h_{0}^{2}\left(1-4 K\left(u-u_{0}\right) / c^{2}\right)$
where the suffix 0 refers, in this case, to the position $\theta=\pi / 2$ measured from the periapsis.
Substituting in equation (8) for $h$, using equation (12), we obtain

$$
\begin{equation*}
\frac{d^{2} u}{d \theta^{2}}+u=\frac{K}{h_{0}^{2}}+\frac{K}{c^{2}}\left[\frac{4 K\left(u-u_{0}\right)}{h_{0}^{2}}+\left(\frac{d u}{d \theta}\right)^{2}+u^{2}\right] \tag{13}
\end{equation*}
$$

### 2.2 Precession of the Periapsis.

Equation (4) is very much easier to apply. This equation is equally applicable to the prediction of satellite trajectories..
The equation which was developed in reference [19] for calculating the precession of the perihelion of Mercury per orbit is

$$
\begin{equation*}
\theta_{P}=\frac{6 \pi G\left(m_{1}+m_{2}\right)}{c^{2} a\left(1-e^{2}\right)} \tag{14}
\end{equation*}
$$

where $a$ is the semi-major axis and $e$ is the eccentricity. This generates $42.89 \operatorname{arcsec} / c e n t u r y$.
For the binary pulsar PSR 1913+16, which was discovered by Hulse and Taylor in 1974, (see reference [22]), the accepted data is that the masses of the two stars are 1.441 and 1.387 times the mass of the Sun, the semi-major axis is $1,950,100 \mathrm{~km}$, the eccentricity is 0.617131 and the orbital period is 7.751939106 hr . Using equation (14) we obtain the result $4.22 \mathrm{deg} / \mathrm{yr}$, which is in agreement with the measured value and that predicted by General Relativity. The orbital decay, or inward spiralling, of binary pulsars is said to be simply due to energy loss caused by gravitational wave emission. This may be the case but energy loss alone will not account for the phenomenon. The loss of mass alone would cause outward spiralling as do most cases of tidal drag.

### 2.3 Moment of Momentum

If the additional term is negligible then it can be shown that the moment of momentum is

$$
h_{2}=h_{1} \exp \left(\frac{2 K}{c^{2}}\left(\frac{1}{r_{1}}-\frac{1}{r_{2}}\right)\right)
$$

which depends on separation but is constant when $c=$ infinity.

### 2.4 Schwarzschild Radius

For a constant radius $v_{r}=0$ and $v=v_{\theta}$ so equation (3) or (3a) becomes

$$
-\frac{K}{r^{2}}\left(1+\frac{v_{\theta}^{2}}{c^{2}}\right) \boldsymbol{e}_{r}=-\frac{v_{\theta}^{2}}{r} \boldsymbol{e}_{r} \quad \text { so if } v_{\theta}=c \quad \text { then } r=\frac{2 K}{c^{2}}=\frac{2 G(M+0)}{c^{2}}=r_{g}
$$

which is known as the Schwarzschild Radius.

### 2.5 Last Stable Orbit

Numerical integration of equation (3a) shows that the Last Stable Orbit occurs when the radius of the orbit is three times the Schwarzschild Radius, to an accuracy of $1 \%$. This is the accepted result based on General Relativity. If equation (3) is used then a value of $2.62 r_{g}$ may be calculated algebraically. However if $Q$ is not unity, as shown in equation (3a), then
equation (12), with $Q$ included, is $\quad h^{2} \approx h_{0}^{2}\left(1-2 r_{g}\left(u-u_{0}\right) Q\right)$
where the suffix 0 refers to circular motion when $u=u_{o}$..
Substituting in equation (8) for $h$, using equation (12a), equation (13) becomes

$$
\begin{equation*}
\frac{d^{2} u}{d \theta^{2}}+u=\frac{K}{h_{0}^{2}}+\frac{K}{c^{2}}\left[\frac{4 K\left(u-u_{0}\right) Q}{h_{0}^{2}}+\left(\frac{d u}{d \theta}\right)^{2}+u^{2}-\left(\frac{d u}{d \theta}\right) 2(Q-1)\right] \tag{13a}
\end{equation*}
$$

If $u=u_{o}+\varepsilon \quad$ then, for small variations

$$
\begin{equation*}
\varepsilon^{\prime \prime}+\varepsilon\left(1-r_{g} u_{o}(2 Q+1)+\left(r_{g} u_{o}\right)^{2} Q\right)=0 \tag{13b}
\end{equation*}
$$

For circular motion it can be shown that $\quad\left(\frac{v}{c}\right)^{2}=\frac{u_{o} r_{g}}{2-u_{o} r_{g}}$ so $Q=\left[1+\frac{1}{2} \sqrt{\frac{u_{o} r_{g}}{2-u_{o} r_{g}}}\right]$
For a stable near circular orbit then $\left(1-r_{g} u_{o}(2 Q+1)+\left(r_{g} u_{o}\right)^{2} Q\right)>0$ so when the factor of $\varepsilon$ in equation (13b) is zero, algebraic manipulation of (13b) gives $r_{d} / r_{g}=3.0$.

### 2.6 Deflection of Light

In equation (3a) terms 2 and 3 are parallel to he velocity so the component normal to the velocity is

$$
\boldsymbol{a} \cdot \boldsymbol{n}=-\frac{K}{r^{2}}\left(1+\left(\frac{v}{c}\right)^{2}\right) \boldsymbol{e}_{r} \cdot \boldsymbol{n}=v \frac{d \psi}{d t} \quad . \text { For small variation of the speed of light assume }
$$

that $v=c$. Also, for small deflections the scalar product of $\boldsymbol{e}_{\boldsymbol{r}}$ and $\boldsymbol{n}$ can be seen from Figure(2) to be $R_{s} / r$. Therefore
$\boldsymbol{a} \cdot \boldsymbol{n}=-\frac{2 K}{r^{2}} \frac{R_{s}}{r} \quad$ now $\quad r=\sqrt{x^{2}+R_{s}^{2}}$ so
$\boldsymbol{a} \cdot \boldsymbol{n}=-\frac{2 K R_{s}}{\left(x^{2}+R_{s}^{2}\right)^{3 / 2}}=c \frac{d \psi}{d t} c \frac{d \psi}{d x} \frac{d x}{d t} \quad$ or, as $d x / d t$ is approximately $c$
$d \psi=-\frac{2 K}{c^{2}} R_{s} \frac{1}{\left(x^{2}+R_{s}^{2}\right)^{3 / 2}} d x \quad$ integrating gives
$\psi=\frac{2 K}{c^{2} R_{s}}\left(\frac{x}{\sqrt{\left(x^{2}+R_{s}^{2}\right)}}\right) . \quad$ When $x=0 \psi=0$ and when $x$ goes to infinity

$$
\begin{equation*}
\psi=-\frac{2 K}{c^{2} R_{s}} \tag{15}
\end{equation*}
$$

Therefore the total deflection $\delta=2 \psi$ so $\quad \delta=\frac{4 K}{c^{2} R_{s}}$

With $K=M_{s u n} G$ and $R_{s}$ being the radius of the Sun the deflection is 1.75 arcsec . This value agrees with the measured value and with General Relativity. This confirms the assumption that the deflection of light grazing the Sun is small.


Figure 2

### 2.7 Shapiro Time Delay

When $v=c$ equation (2a) gives $\boldsymbol{a} \cdot \boldsymbol{t}=\frac{K}{r^{2}} \frac{v_{r}}{c_{0}}=\frac{\mathrm{d} c}{\mathrm{~d} t} \quad$ where $c_{0}$ is the speed of light in empty space.
Therefore $\quad \frac{K}{r^{2}} \frac{\dot{r}}{c_{0}}=\frac{d c}{d r} \frac{d r}{d t} \quad$ or $\quad \frac{K}{r^{2} c_{0}} d r=d c$

Integration leads to

$$
[c]_{1}^{\infty}=\left[-\frac{K}{r c_{0}}\right]_{1}^{\infty}, \quad c_{0}-c=0+\frac{K}{r_{1} c_{0}}
$$

$$
\text { or } \quad c=c_{0}\left(1-\frac{G M}{r c_{0}^{2}}\right)
$$

This is as quoted by Einstein in 1911. [32]
Consider the case of light grazing the Sun at a radius $R_{s}$ and calculate the journey time. As an approximation assume the path to be a straight line. Then, since $c=\frac{d x}{d t}$
$c_{0} d t=\frac{d x}{1-G M /\left(r c_{0}^{2}\right)} \quad$ where $\quad r=\sqrt{x^{2}+R_{s}^{2}} \quad$ also as $\frac{G M}{r c_{0}^{2}} \ll 1$

$$
\int c_{0} d t=\int\left(1+\frac{G M / c_{o}^{2}}{\sqrt{x^{2}+R_{s}{ }^{2}}}\right) d x=x+\frac{G M}{c_{0}^{2}} \ln \left[2\left(\sqrt{x^{2}+R_{s}{ }^{2}}+x\right)\right]
$$

Between the limits $\mathrm{x}=0$ to x and $t=0$ to $t$ we have the total time
$t=\frac{x}{c_{0}}+\frac{G M}{c_{0}^{3}} \ln \left(2 \sqrt{x^{2}+R_{s}^{2}}+2 x\right)-\frac{G M}{c_{0}^{3}} \ln \left(2 R_{s}\right)$
so the additional time due to the gravitational effect is

$$
\Delta t=\frac{G M}{c_{0}^{3}} \ln \left[\frac{\sqrt{x^{2}+R_{s}^{2}}+x}{R_{s}}\right]
$$

This gives a time delay of $232 \mu$ s for the double transit time from Earth to Venus. This is exactly the same as quoted by Bertotti, reference made to C. M. Will. The value for Mars is $247 \mu \mathrm{~s}$, which is as quoted by Reasenberg, Shapiro et. al., is also given by the above equation.

This shows that the speed of light is reduced by gravity. In the case of light grazing the Sun the reduction in speed $/ \mathrm{c}_{\mathrm{o}}$ is $\frac{G M_{s}}{c_{0}^{2} R_{s}} \approx 2 \times 10^{-6}$, that is, only two parts per million. The variable speed of light has been incorporated into the basic equation however this would not have any major effect on the motion of material bodies. The variation of the deflection of light is not changed. Because the light path is slightly curved the increase in path length will add to the delay but as the speed change is so small the additional time is less than $1 \%$.

### 3.0 GRAVITOMAGNETICS APPLIED TO ROTATING BODIES

### 3.1 Basic Equations

When equation (2) is applied to two body systems the equation generated is identical to the de Sitter form and agrees with the measurement of precession of the perihelion of Mercury and of the Binary Pulsar PSR $1913+16$. The equation is equally applicable if one spherical body is large and non-rotating. Note that the additional term, which is a function of $c^{3}$, is negligible for Solar dynamics.

$$
\begin{equation*}
\boldsymbol{a}=-\frac{K}{r^{2}}\left(1-\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 K}{r^{2} c^{2}} \boldsymbol{v} \times\left(\boldsymbol{v} \times \boldsymbol{e}_{r}\right) \tag{2rpt}
\end{equation*}
$$

where $\quad c$ is the speed of light, $\quad \boldsymbol{a}$ is relative acceleration, $v$ is relative velocity and $r$ is relative position. Also $\boldsymbol{e}_{r}$ is the unit vector from body A to body B.
$K=G\left(m_{A}+m_{B}\right)$ where $G$ is the gravitational constant

The calculations are made easier for multi-body systems by the use of a defined force as shown in equation (5).

$$
\begin{equation*}
\boldsymbol{P}=\mu \boldsymbol{a}=-\frac{G m_{A} m_{B}}{r^{2}}\left(1-\frac{v^{2}}{c^{2}}\right) \boldsymbol{e}_{r}+\frac{2 G m_{A} m_{B}}{r^{2} c^{2}} \boldsymbol{v} \times\left(\boldsymbol{v} \times \boldsymbol{e}_{r}\right) \tag{5rpt}
\end{equation*}
$$

$\mu$ is the reduced mass $m_{A} m_{B} /\left(m_{A}+m_{B}\right)$

### 3.2 Gravity Probe B

Gravity Probe B is the study of the precession of a gyroscope in a polar orbit about the Earth.
When the new theory is applied to Gravity Probe B the following equations are derived algebraically using equation (5).

$$
\begin{aligned}
& \dot{\phi}_{x}=\frac{G m_{E} \dot{\alpha}}{c^{2} R_{c}} \quad \text { Geodesic or de Sitter. (North toSouth) } \\
& \dot{\phi}_{z}=\frac{G I_{E} \Omega}{2 c^{2} R_{c}^{3}}\left(2-3 \cos ^{2} \alpha\right) \quad \text { frame dragging or Lense - Thirring. (West to East) } \\
& \text { Radius of the Earth } \\
& \text { Mass of the Earth } \\
& \text { Moment of inertia of the Earth } \\
& \text { Height of orbit } \\
& \text { Radius of orbit } \\
& \text { Gravitational constant } \\
& \text { Speed of light } \\
& R_{E}=6378000 \mathrm{~m} \\
& m_{\mathrm{E}}=5.978 \times 10^{24} \mathrm{~kg} \\
& I_{E}=8.04 \times 10^{37} \mathrm{~kg} \mathrm{~m}^{2} \\
& =650 \mathrm{~km} \\
& R_{c}=6378000+650000=7028000 \mathrm{~m} \\
& G=6.673 \times 10^{-11} \mathrm{~N} \mathrm{~m}^{2} \mathrm{~kg}^{-2} \\
& c=2.998 \times 10^{8} \mathrm{~m} \mathrm{~s}^{-1}
\end{aligned}
$$

$\alpha$ is the location of the satellite and $\Omega$ is the angular velocity of the Earth.
These equations yield,
$\dot{\phi}_{x}=4.4 \pm 0.0 \quad \mathrm{arcsec} / \mathrm{yr} \quad(6.6 \quad \pm 0.0 \quad$ based on general theory by Schiff $)$
$\dot{\phi}_{z}=0.02 \pm 0.06 \mathrm{arcsec} / \mathrm{yr} \quad(0.042 \pm 0.126 \quad " \quad " \quad " \quad " \quad " \quad ")$

### 3.3 Precession of the Pericentre of a small body orbiting a large rotating mass

Consider a test body in orbit around a spherical body rotating at an angular speed of $\Omega$. The test body is performing an elliptical orbit with a period of $T$ in a plane that has a inclination (inc) relative to the equatorial plane of the rotating body.

The rate of precession of the pericentre, as seen from the plane of the orbit, in radians per orbit, is

$$
\begin{equation*}
\Delta \phi_{P}=\frac{6 \pi G M}{c^{2} a\left(1-e^{2}\right)}-\frac{2 G I \Omega}{c^{2} a^{3}\left(1-e^{2}\right)^{3 / 2}} \cos (\text { inc }) T \tag{16}
\end{equation*}
$$

Where $I$ is the moment of inertia, $e$ is the eccentricity and $a$ is the semi-major axis..
The first term is the de Sitter precession and has been derived algebraically from equation (2). It agrees exactly with the generally accepted form and agrees with the measured results for the precession of the perihelion of Mercury and for the binary pulsar PSR 1913+16. However, the second term, the LenseThirring term, justified by numerical integration, is only half of the generally accepted value.

### 3.4 Nodal Precession

Nodal precession is the rate of precession of the line of the intersection of the plane of the orbit and the equatorial plane, as seen from the equatorial plane. Again, in radians per orbit

$$
\begin{equation*}
\Delta \phi_{N}=\frac{G I \Omega}{c^{2} a^{3}\left(1-e^{2}\right)^{3 / 2}}\left[1-\frac{1}{2}\left(\frac{\Omega}{\omega^{*}}\right) \cos (\text { inc })\right] T \tag{17}
\end{equation*}
$$

Where $\omega^{*}=\frac{2 \pi}{T}=\sqrt{\frac{G M}{a^{3}}}$ is the mean angular velocity of the test body.
Compared with the accepted result the first term is, again, only half its value. The second term does not even appear in the accepted equation.

### 4.0 DISCUSSION

Equation (4a) is easier to apply than GR and therefore leaves less room for misinterpretation. That force is a secondary quantity was strongly advocated by H. R. Hertz who regarded force as "a sleeping partner". Force is to dynamics as money is to commerce. Once force has been demoted to a defined quantity then force fields and inertia are also defined quantities, similarly for work and energy. Equation (2) is loosely modelled on the Lorentz force but this relationship is for guidance only in the same way that Maxwell used a mechanical model to form his equations. However, he abandoned the reference in his final paper on the subject once he had established that his equations predicted the then known observations.

As shown above, when the new approach is applied to two body systems it agrees with the well verified observations of the precession of the perihelion of Mercury and the definition of the Schwarzschild Radius. Both agree with the results obtained from the General Theory of Relativity. The third term in equation (4a) was added as it generates the accepted value for the Last Stable Orbit and agrees with the measurements of the Shapiro Time Delay.

The de Sitter effect agrees with the accepted results of analysis whether algebraically or by numerical integration for two body systems or large non-rotating bodies. This is true whether using equation (2) or equation (5).

However, for the Lense-Thirring terms there is an unresolved factor of two which affects both nodal and pericentre precession. Also, an extra term appears in the expression for the nodal precession which is not in the accepted result.

The published test on the Earth satellites LAGEOS I \& II, see reference [28], appear to agree with the accepted theory. The inclination of the satellites is approximately $90^{\circ}+/-20^{\circ}$. The reason for this is that the accepted Lense-Thirring term does not depend on the inclination but all other effects do and therefore can be cancelled out. See also references [10] and [29]. The extra term in the present theory also depends on the inclination and therefore cancels out as well. In this configuration it is small compared to the main term.

The Gravity Probe B experiment testing the precession of gyroscopes in Earth orbit displays two equations, one for the geodesic term and one for the frame-dragging effect. The geodesic term does not involve the rotation of the Earth but the frame-dragging term does. The same form of equations have been generated algebraically using equation (5) but the frame-dragging term is half of the published value. However, the geodesic term is two thirds of the published value.

The gravitational effect on the speed of light is still discussed but apart from the Shapiro Time Delay the effect is negligible when dealing with the motion of bodies. The decrease of the speed of light grazing the Sun is only 2 parts per million.

It has proved to be impossible, so far, to find any modification to equation (4) such that it gives the generally accepted value for the Lense-Thirring effect without changing the de Sitter effect applications. The de Sitter results have been obtained by several observations but the Lense-Thirring effect is very small compared to other effects. In the LAGEOS experiments for the precession of the pericentre the Lense-Thirring effect is less than $1 \%$ of the de Sitter effect, which makes it more difficult to evaluate. The GP-B test results have recently been published, reference [27]. There are four gyroscopes, two of which have original frame-dragging results which are close to that predicted by the new theory. The geodesic results are, on average, close to the that of the accepted value. Nevertheless, over a one month period two of the gyroscopes precess at a rate close to the new theory predictions.

It is widely stated that the inward spiralling of a binary star system is due to gravitational radiation. The loss of energy alone is not the cause of this effect. Energy loss can be related to outward spiralling, as is the case for the Earth Moon system. However, radiation pressure could be the cause.

When general relativity is applied to multiple body systems several authors have produced slightly different results. Some results even do not return to the Newtonian form when the velocities are zero but only if the speed of light is taken to be infinite. There are also comments made that the frame dragging effect has been verified by the binary pulsar mentioned above when, in fact, it is a verification of the de Sitter effect.

This new approach does not undermine the General Theory of Relativity but because it is a simpler method it leaves less room for misinterpretation. Many of the extensions of GR are very complex mathematically, making errors more likely.

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