# The nature of quantum particles

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Abstract: In this paper we will discuss a recent double-slit experiment with quantum particles where an interference pattern is observed even when which-path information is available, and argue that this result is in clear contradiction with orthodox quantum mechanics but perfectly understandable in the framework of a nonlinear quantum physics.

Keywords: nonlinear quantum physics, orthodox quantum mechanics, complementarity principle, double-slit experiment

### 1. Introduction

As is well known, if in a simple double-slit experiment with quantum particles a which-path measurement is made, i.e., if we observe which slit the particle goes through, then no interference pattern is observed on the far screen. That is, it is not possible to have which-way information and an interference pattern at the same time.

According to the orthodox interpretation of quantum mechanics<sup>1</sup>, this can be "explained" by applying Niels Bohr's complementarity principle: there are pairs of features (concepts, observables, etc.) which can be observed individually but never at the same time, and both must be used if we want to obtain a full description of a phenomenon. This happens because the wave  $\psi$  that is divided into two as it reaches the double slit simply represents the probability of observing the corresponding particle in a certain state. The resulting waves  $\psi_1$  and  $\psi_2$  therefore represent the probability of observing the particle coming through the upper and lower slits, respectively, and so if we observe the particle just outside one of the slits then the wave going through the remaining one has to collapse, and so no interference pattern will appear on the screen.

One could, however, argue that such a measurement necessarily disturbs the particle whose state is being observed, e.g., by placing a detector just outside one of the slits, and thus it is natural to expect, even without invoking any particular principle, that the interference pattern is simply washed out. In fact, this simple and more intuitive explanation is the one obtained in the framework of a nonlinear quantum physics<sup>2</sup> inspired in the early ideas proposed by Louis de Broglie<sup>3</sup>. According to this approach the wave  $\psi$  actually exists as a physical entity, and guides its localized corpuscle preferentially along a path where its intensity is higher. Therefore, as a particle reaches the double slit, the corpuscle takes only one of the paths but its guiding wave, which is divided into two, takes both. As a consequence, one of the resulting waves,  $\psi_1$  or  $\psi_2$ , will travel alone towards the screen. In this situation, if a detector is placed just outside one of the slits then the wave going

through it will either be absorbed or have its phase altered, and thus the interference pattern will naturally disappear.

In order to distinguish between these two interpretations we would, therefore, need an experiment where which-path information is available but the particles going through the double slit are not disturbed in any way in their path to the far screen. In this paper, whose discussion can be found in a previous work<sup>4</sup>, we will look at a recent, more sophisticated double-slit experiment where for the first time this has actually been achieved, and where, as a consequence, an interference pattern is still observed. We will argue that this result is in clear contradiction with orthodox quantum mechanics but is perfectly understandable if we are willing to take a nonlinear approach.

# 2. A two-photon double-slit experiment

We will now present and discuss the double-slit experiment recently performed by Menzel, Puhlmann, Heuer and Schleich<sup>5</sup> in 2012 and repeated by the same authors one year later. The setup is shown in Fig. 1.

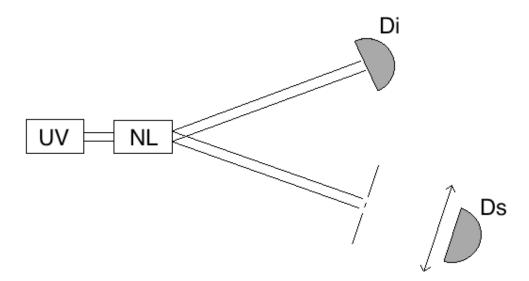


Fig. 1 – Menzel's double-slit experiment.

The source in this experiment is a UV pump beam composed of two distinct intensity maxima. Therefore, each photon  $\psi$  can, as soon as it leaves the pump, be represented by a superposition of two waves, one corresponding to its upper maximum,  $\psi_1$ , and one representing its lower maximum,  $\psi_2$ .

The beam is injected onto a nonlinear crystal NL that transforms an incoming photon  $\psi$  into a pair of photons with half the frequency of the first – one called idler and represented by the wave  $\psi_i$  and the other one called signal,  $\psi_s$ . Only one pair of photons, correlated in space and time, is produced at a time. Thus, in this scheme each photon produced at the crystal can either be detected at its own upper maximum or at its lower maximum, and so

the wave-functions of the idler and signal photons can be written as  $\psi_i = \psi_{i1} + \psi_{i2}$  and  $\psi_s = \psi_{s1} + \psi_{s2}$ , respectively.

The idler beam is directly incident onto detector  $D_i$  and the signal one is directed onto a double slit, in a way that the beam's upper maximum  $\psi_{s1}$  is injected onto the upper slit and its lower maximum  $\psi_{s2}$  is incident on the lower slit. The outcoming waves then diffract and are incident onto detector  $D_s$ , which is scanned along a direction perpendicular to the signal photon's trajectory. The signal at  $D_s$  can then be counted in coincidence with the detections at  $D_i$ .

Now, this experiment is divided in two parts:

- 1 Detector  $D_s$  is first placed just outside the double slit for a near-field detection. In this case, the spatial correlation between the signal and the idler photons is observed: when an idler photon arrives through the upper maximum, which corresponds to the wave  $\psi_{i1}$ , a signal photon is detected in coincidence outside the upper slit,  $\psi_{s1}$ , and if an idler photon arrives through the lower maximum  $\psi_{i2}$  then a signal photon is detected outside the lower slit,  $\psi_{s2}$ .
- $2 D_s$  is then moved away from the double slit for a far-field detection. In this case an interference pattern is always observed when both slits are open, independently on whether the detections at  $D_s$  are counted alone or in coincidence with the ones at  $D_i$ , but disappears when one of the slits is closed.

We will now discuss the results from the point view of orthodox quantum mechanics and then in terms of a nonlinear physics.

# 2.1. Orthodox interpretation

According to the orthodox interpretation of quantum mechanics, the photon  $\psi$  leaving the UV source is in a superposition of two states  $\psi_1$  and  $\psi_2$  corresponding to the equal probabilities of the particle being detected at the upper and lower maxima, respectively.

When the incoming photon arrives at the nonlinear crystal NL, an outcoming pair of idler  $\psi_i$  and signal  $\psi_s$  photons is produced in an entangled state. This means that if an idler photon is detected by  $D_i$  at the upper maximum then detector  $D_s$  will receive a corresponding signal photon through the upper slit, which will make its wave-function collapse to the upper state, i.e.,  $\psi_s \to \psi_{s1}$ , and when an idler photon is detected at the lower maximum a signal photon will be observed through the lower slit, which means that a collapse will occur to the lower state,  $\psi_s \to \psi_{s2}$ .

Now, as  $D_s$  is moved away from the double slit to allow for the far-field detection, measuring the position of the idler photon gives us the information about which slit the corresponding signal photon goes through. In this situation, as described above, the latter's wave-function will collapse to one of the upper or lower states, and thus  $\psi_s \to \psi_{s1}$  if it

comes through the upper slit or  $\psi_s \to \psi_{s2}$  if it arrives through the lower slit. As a consequence, and contrary to the experimental results, no interference pattern should appear at the far screen.

In sum, the which-path information obtained due to the entanglement of the photon pair did not collapse the system's wave-function – which had to happen if the wave simply represented the probability of a particle being observed in a certain state – and so did not avoid the observation of an interference pattern. We have therefore shown that, in this particular case, there has been a clear violation of the orthodox interpretation of quantum mechanics, as Bohr's complementarity cannot be invoked in the way that it has been done in previous double-slit experiments.

# 2.2. Nonlinear approach

In a nonlinear approach, each photon  $\psi$  emitted by the UV pump is composed of a corpuscle that is guided by a wave with two maxima, an upper one represented by  $\psi_1$  and a lower one  $\psi_2$ . As the corpuscle leaves the source along the path of either maxima of its guiding wave  $\psi = \psi_1 + \psi_2$ , it will tend to remain there and avoid the intensity minimum at the centre of the wave, which means that the other maximum will travel alone towards the nonlinear crystal NL.

As the incoming photon enters the crystal, an idler photon  $\psi_i$  and a signal photon  $\psi_s$  are produced, each one composed of its own corpuscle and guiding wave with the two maxima now represented by  $\psi_{i1}$  and  $\psi_{i2}$  for the idler photon and by  $\psi_{s1}$  and  $\psi_{s2}$  for the signal photon. Moreover, each corpuscle produced will preferentially remain, as it travels along its path, in the same (upper or lower) maximum as the original one.

Now, when they both reach their corresponding detectors in the first part of the experiment, the idler and signal photons will naturally be detected at the same upper or lower maximum, which agrees with the experimental results.

When detector  $D_s$  is moved away to the far field, in the second part of the experiment, the signal photon's wave  $\psi_s$  reaching the double slit will be composed of the same upper  $\psi_{s1}$  and lower  $\psi_{s2}$  maxima. Therefore, it is obvious that an interference pattern will appear independently on whether we detect the signal alone or in coincidence with the idler photons. Naturally, if one of the slits is closed then the wave incident on it will be blocked and thus the interference pattern will disappear.

We thus see that, from a nonlinear point of view, it is easy to understand why an interference pattern can be present even if which-path information is available to the observer. Contrary to previous setups, in this particular one the particles going through the double slit have not been disturbed in any way in their path to the detector placed in the far field, and thus there is no reason for the interference pattern to be washed out.

#### 3. Conclusion

In this paper we have shown that the experimental results obtained by Menzel *et al* are in clear contradiction with the orthodox interpretation of quantum mechanics, as in their setup the observation of an interference pattern after a double slit does not depend on whether which-path information is available to the observer. Thus, the incoming wave  $\psi$  that splits into two as it reaches the slits cannot simply represent a probability wave.

Moreover, we have made it clear why in this experiment, according to a nonlinear approach, the interference pattern does not disappear when which-way information is available. Essentially, a simple and most natural explanation is that there is a real physical wave guiding the corpuscle along its path, and therefore there is no reason for it to collapse even if we know which slit the corpuscle goes through. Finally, if we accept that the wave is physical then it is as well natural that the interference does not occur when one of the slits is closed.

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