

MODIFIED ALCUBIERRE WARP DRIVE I:

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Abstract:

A solution of general relativity is presented that describes an Alcubierre [1] propulsion system in which it is possible to travel at superluminal speed while reducing the components of the energy-impulse tensor (thus reducing energy density) by an arbitrary value. This solution also allows us to reduce Hawking radiation which would otherwise lead to the explosion of the spaceship [7].

Introduction:

Alcubierre [1] in 1994 proposed a solution of the equations of general relativity which provides the only viable means to accelerate a spaceship up to superluminal velocities without using wormholes. A problem was soon identified: Pfenning [4] showed that the required energy is comparable to the total energy of the universe and that it is negative. Moreover he used quantum inequalities to show that this energy gets distributed at very short scale (about 100 times the Planck length) up to a multiplicative factor equal to the squared speed. This paper (part I), and the following one (part II) going to be published soon, investigate these problems, introducing a way to reduce the amount of energy involved and its spacial distribution within the warp bubble. Later Hiscock [10] proved the existence of an event horizon for superluminal travels which would imply the presence of Hawking radiation responsible for the rapid destruction of the spaceship, and this problem is also addressed and solved in this publication.

Note: In the following we adopt the notation used by Landau and Lifshitz in the second volume (“The Classical Theory of Fields”) of their well known Course of Theoretical Physics [12].

We propose the use of the metric

$$(1) \quad ds^2 = \left(1 - v^2 \frac{f(x, y, z - k(t))^2}{a(x, y, z - k(t))^2} \right) dt^2 + 2v \frac{f(x, y, z - k(t))}{a(x, y, z - k(t))} dt dz - dx^2 - dy^2 - dz^2$$

while the one used by Miguel Alcubierre [1] is

$$(2) \quad ds^2 = dt^2 - [dz - v f(x, y, z - k(t)) dt]^2 - dx^2 - dy^2$$

Our proposed metric (1) can be recast as [4]

$$ds^2 = dt^2 - \left[dz - v \frac{f(x, y, z - k(t))}{a(x, y, z - k(t))} dt \right]^2 - dx^2 - dy^2$$

Pfenning zone, i.e. the zone within the interval: $R - \frac{\Delta}{2} < r < R + \frac{\Delta}{2}$ where $\Delta \ll 1$

R radius of the Warp bubble Δ wall thickness of the Warp bubble $R \gg \Delta$

$$r = (x^2 + y^2 + (z - k(t))^2)^{\frac{1}{2}} \quad \text{and} \quad \frac{dk(t)}{dt} = v = \text{const}$$

In the Pfenning zone we let $a(x, y, z - k(t)) \gg 1$ (there is the source of esotic matter),

$$\text{where } z_0(t) = k(t) \quad \text{and} \quad \frac{dk(t)}{dt} = v = \text{const}$$

Einstein tensor in contravariant form in the Pfenning zone is:

$$\text{where } D_i = \frac{\partial}{\partial x^i}, x^i = x, y, z (i=1,2,3) \text{ e } D_{i,k} = \frac{\partial^2}{\partial x^i \partial x^k}, x^i, x^k = x, y, z (i, k=1,2,3)$$

Detailed computation:

$$G^{tt} = -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^4} (v^2 (f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 + f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2$$

$$+ D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 + D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2$$

$$- 2 D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t))$$

$$- 2 D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)))$$

$$G^{tx} = \frac{1}{2} \frac{1}{a(x, y, z - k(t))^3} (v (-a(x, y, z - k(t)) D_3(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t))$$

$$- a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t))$$

$$- a(x, y, z - k(t)) f(x, y, z - k(t)) D_{1,3}(a)(x, y, z - k(t))$$

$$+ 2 f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) + a(x, y, z - k(t))^2 D_{1,3}(f)(x, y, z - k(t)))$$

$$G^{ty} = \frac{1}{2} \frac{1}{a(x, y, z - k(t))^3} (v (-a(x, y, z - k(t)) f(x, y, z - k(t)) D_{2,3}(a)(x, y, z - k(t))$$

$$- a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t))$$

$$- a(x, y, z - k(t)) D_3(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) + a(x, y, z - k(t))^2 D_{2,3}(f)(x, y, z - k(t))$$

$$+ 2 f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)))$$

$$G^{tz} = -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^5} (v (-2 a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_{2,2}(a)(x, y, z - k(t))$$

$$- 2 a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_{1,1}(a)(x, y, z - k(t))$$

$$+ 4 D_2(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 f(x, y, z - k(t))$$

$$\begin{aligned}
& -4 D_2(a)(x, y, z - k(t)) a(x, y, z - k(t))^3 D_2(f)(x, y, z - k(t)) \\
& + 4 D_1(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 f(x, y, z - k(t)) \\
& - 4 D_1(a)(x, y, z - k(t)) a(x, y, z - k(t))^3 D_1(f)(x, y, z - k(t)) \\
& + v^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t))^2 \\
& + v^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t))^2 + v^2 f(x, y, z - k(t))^3 D_1(a)(x, y, z - k(t))^2 \\
& + v^2 f(x, y, z - k(t))^3 D_2(a)(x, y, z - k(t))^2 + 2 a(x, y, z - k(t))^4 D_{2, 2}(f)(x, y, z - k(t)) \\
& + 2 a(x, y, z - k(t))^4 D_{1, 1}(f)(x, y, z - k(t)) \\
& - 2 v^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& - 2 v^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t))
\end{aligned}$$

$$\begin{aligned}
G^{xx} = & -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^4} \left(v \left(-v D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \right. \right. \\
& + 2 v D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& - v f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 + 4 v D_3(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - 4 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{3, 3}(f)(x, y, z - k(t)) + v D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& + v f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 + 4 v f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{3, 3}(f)(x, y, z - k(t)) \\
& + 4 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{3, 3}(a)(x, y, z - k(t)) \\
& \left. - 4 v f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{3, 3}(a)(x, y, z - k(t)) \right)
\end{aligned}$$

$$\begin{aligned}
& + 8 D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) \\
& - 8 D_3(a)(x, y, z - k(t))^2 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) \\
& + 12 v f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t))^2 \\
& - 16 v D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - 2 v D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \Big) \Big)
\end{aligned}$$

$$G^{xy} = \frac{1}{2} \frac{1}{a(x, y, z - k(t))^4} (v^2 (D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) - f(x, y, z - k(t)) D_1(a)(x, y, z - k(t))) (D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) - f(x, y, z - k(t)) D_2(a)(x, y, z - k(t))))$$

$$\begin{aligned}
G^{xz} = & - \frac{1}{2} \frac{1}{a(x, y, z - k(t))^4} \left(v \left(\left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{1,3}(f)(x, y, z - k(t)) \right. \right. \\
& - 2 v f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{1,3}(f)(x, y, z - k(t)) \\
& + 2 v f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{1,3}(a)(x, y, z - k(t)) \\
& - \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - 2 v D_3(f)(x, y, z - k(t)) a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t)) \\
& - \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{1,3}(a)(x, y, z - k(t)) \\
& \left. \left. - \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \right) \right)
\end{aligned}$$

$$\begin{aligned}
& + 4 \nu D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + 4 \nu f(x, y, z - k(t)) a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - 6 \nu f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + 2 D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \Big) \Big) \\
G^{yy} = & - \frac{1}{4} \frac{1}{a(x, y, z - k(t))^4} \left(\nu \left(-\nu D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \right. \right. \\
& + 2 \nu D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& - \nu f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 + 4 \nu D_3(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - 4 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{3,3}(f)(x, y, z - k(t)) + \nu D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& + \nu f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 + 4 \nu f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{3,3}(f)(x, y, z - k(t)) \\
& + 4 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{3,3}(a)(x, y, z - k(t)) \\
& - 4 \nu f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{3,3}(a)(x, y, z - k(t)) \\
& + 8 D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) \\
& - 8 D_3(a)(x, y, z - k(t))^2 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) \\
& + 12 \nu f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t))^2 \\
& - 16 \nu D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_3(a)(x, y, z - k(t))
\end{aligned}$$

$$\left. - 2 v D_1(f(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t))) \right)$$

$$\begin{aligned}
G^{yz} = & -\frac{1}{2} \frac{1}{a(x, y, z - k(t))^4} \left(v \left(\left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{2,3}(f)(x, y, z - k(t)) \right. \right. \\
& + 2 v f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{2,3}(a)(x, y, z - k(t)) \\
& - 2 v f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{2,3}(f)(x, y, z - k(t)) \\
& - \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{2,3}(a)(x, y, z - k(t)) \\
& - \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& - \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - 2 v D_3(f)(x, y, z - k(t)) a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t)) \\
& + 4 v D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& + 4 v f(x, y, z - k(t)) a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - 6 v f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& \left. \left. + 2 D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \right) \right)
\end{aligned}$$

$$\begin{aligned}
G^{zz} = & -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^6} (v^2 (4 f(x, y, z - k(t)) a(x, y, z - k(t))^4 D_{1,1}(f)(x, y, z - k(t)) \\
& + 4 f(x, y, z - k(t)) a(x, y, z - k(t))^4 D_{2,2}(f)(x, y, z - k(t))
\end{aligned}$$

$$\begin{aligned}
& + 11 f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& + 11 f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - 4 f(x, y, z - k(t))^2 a(x, y, z - k(t))^3 D_{1,1}(a)(x, y, z - k(t)) \\
& - 4 f(x, y, z - k(t))^2 a(x, y, z - k(t))^3 D_{2,2}(a)(x, y, z - k(t)) + 3 D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^4 \\
& + 3 D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^4 \\
& - 2 v^2 f(x, y, z - k(t))^3 a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& - 2 v^2 f(x, y, z - k(t))^3 a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& + v^2 f(x, y, z - k(t))^2 a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t))^2 \\
& + v^2 f(x, y, z - k(t))^2 a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t))^2 + v^2 f(x, y, z - k(t))^4 D_1(a)(x, y, z - k(t))^2 \\
& + v^2 f(x, y, z - k(t))^4 D_2(a)(x, y, z - k(t))^2 \\
& - 14 D_2(f)(x, y, z - k(t)) a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& - 14 D_1(f)(x, y, z - k(t)) a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_1(a)(x, y, z - k(t))
\end{aligned}$$

Because of the value $a = a(x, y, z - k(t)) \gg 1$ (extremely large) the components of T^{ik} can be

reduced by an arbitrary value in the Pfenning zone

$$\text{Einstein Equation: } G^{ik} = \frac{8\pi G}{c^4} T^{ik}$$

G^{ik} proportional to T^{ik} (energy-impulse tensor)

The function $f(r)$ assumes the following values:

$$r = (x^2 + y^2 + (z - k(t))^2)^{\frac{1}{2}} \quad \text{where} \quad \frac{dk(t)}{dt} = v$$

$$f = f(r) = 0 \quad \text{for every } r \text{ such that } r > R + \frac{\Delta}{2}$$

$$0 < f = f(r) < 1 \quad \text{for every } R - \frac{\Delta}{2} < r < R + \frac{\Delta}{2}, \text{ e.g., } f = -\frac{(r - R - \frac{\Delta}{2})}{\Delta} \quad [4] \text{ (Pfenning zone)}$$

$$f = f(r) = 1 \quad \text{for every } r \text{ such that } 0 < r < R - \frac{\Delta}{2}$$

The function $a(r)$ assumes the following values:

$$a = a(r) = 1 \quad \text{for every } r \text{ such that } r > R + \frac{\Delta}{2}$$

$$a = a(r) = a(x, y, z - k(t)) \gg 1 \quad \text{extremely large for every } r \text{ such that } R - \frac{\Delta}{2} < r < R + \frac{\Delta}{2}$$

(Pfenning zone)

$$a = a(r) = 1 \quad \text{for every } r \text{ such that } 0 < r < R - \frac{\Delta}{2}$$

HISCOCK SOLUTION AND HAWKING RADIATION [10]:

With few calculations, the Alcubierre metric (2) becomes $ds^2 = A(r) d\tau^2 - \frac{dr^2}{A(r)}$ [10] which

using our modified metric (1) turns out to imply $A(r) = 1 - v^2 \left[1 - \frac{f}{a} \right]^2$

(event horizon at $v \geq 1$) .

Therefore Hawking temperature T is: $T = \frac{A'(r=R+\frac{\Delta}{2})}{4\pi} = v^2 \left(\frac{f}{a} \right)' = v^2 \frac{((\frac{f'}{a}) - \frac{fa'}{a^2})}{2\pi}$ with

$(\prime = \frac{d}{dr})$ where $\frac{a'}{a^2} \ll 1$ for $a(x, y, z - k(t)) \gg 1$ and both $f(r)$ and $f'(r)$

obtained for $r = R + \frac{\Delta}{2}$. If $a' \simeq a$ then Hawking radiation is practically irrelevant. As a matter of fact in the general case $A' = -v^2(1 - \frac{f}{a})(\frac{f}{a})'$ $\frac{f}{a} \ll 1$ in the region of esotic matter (Pfenning zone) for superluminal travels there is already a event horizon at $r > R - \frac{\Delta}{2}$ and even in this more general case Hawking radiation is irrelevant.

INTERNAL METRIC OF THE WARP BUBBLE (which contains the spaceship):

$$0 < r < R - \frac{\Delta}{2}$$

$$ds^2 = (1 - v^2) dt^2 + 2v dt dz - dx^2 - dy^2 - dz^2 \quad \text{or} \quad ds^2 = dt^2 - (dz - v dt)^2 - dx^2 - dy^2$$

Moving with velocity v (multiple of the speed of light c) along the z -axis.

METRIC OUTSIDE OF THE BUBBLE BEYOND THE PFENNING ZONE:

$$r > R + \frac{\Delta}{2}$$

$$ds^2 = dt^2 - dx^2 - dy^2 - dz^2$$

i.e., Minkowski space.

Conclusions: These calculations show that the modified Alcubierre propulsion system can achieve superluminal speeds without giving rise to problems in the energy density and in the components of the energy-impulse tensor. Hawking radiation becomes irrelevant and the problem which would lead to the explosion of the warp bubble [7] does not occur.

Appendix I:

The components of the Einstein tensor proportional to the energy-impulse tensor have been calculated with reference to an observer whose gravitational field is very weak and whose speeds

are far lower than the speed of light, observing the spaceship and the warp bubble moving at speed v , i.e., an inertial reference frame in which the spaceship is moving at speed v). If we want to calculate in the Eulerian reference frame, that is moving with the spaceship, we get for each component of the energy-impulse tensor in implicit form the following:

$$(\text{energy density}) = k G^{tt}$$

$$(\text{impulse density } x) = -k G^{tx}$$

$$(\text{impulse density } y) = -k G^{ty}$$

$$(\text{impulse density } z) = k \left[G^{tz} v \left(\frac{f}{a} \right) - G^{tz} \right]$$

$$(\text{stress } xx) = k G^{xx}$$

$$(\text{stress } yy) = k G^{yy}$$

$$(\text{stress } zz) = k \left[G^{zz} v^2 \left(\frac{f}{a} \right)^2 - 2v \left(\frac{f}{a} \right) G^{tz} + G^{zz} \right]$$

$$(\text{stress } xy) = k G^{xy}$$

$$(\text{stress } xz) = k \left[-v \left(\frac{f}{a} \right) G^{xt} + G^{xz} \right]$$

$$(\text{stress } yz) = k \left[-v \left(\frac{f}{a} \right) G^{yt} + G^{yz} \right]$$

The various stress xx , stress yy , stress zz , stress xy , stress xz , stress yz and their symmetric counterparts are the components of the stress tensor [12], and $k = \frac{c^4}{8\pi G}$; c is the speed of

light and G is Newton gravitational constant; $a = a(x, y, z - k(t)) \gg 1$ in Pfenning zone where the energy-impulse tensor is not zero.

Appendix II: In the non-constant case $v = v(t)$ the metric becomes:

$$ds^2 = \left(1 - v(t)^2 \frac{f(x, y, z - k(t))^2}{a(x, y, z - k(t))^2} \right) dt^2 + 2v(t) \frac{f(x, y, z - k(t))}{a(x, y, z - k(t))} dt dz - dx^2 - dy^2 - dz^2$$

which can be recast in implicit form as:

$$ds^2 = dt^2 - \left[dz - v(t) \frac{f(x, y, z - k(t))}{a(x, y, z - k(t))} dt \right]^2 - dx^2 - dy^2$$

Using this metric we can calculate the Einstein tensor as follows:

$$\begin{aligned} G^{tt} = & -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^4} (v(t)^2 (f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 + f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 \\ & + D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 + D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\ & - 2 D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\ & - 2 D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t))) \end{aligned}$$

$$\begin{aligned} G^{tx} = & -\frac{1}{2} \frac{1}{a(x, y, z - k(t))^3} (v(t) (a(x, y, z - k(t)) f(x, y, z - k(t)) D_{1,3}(a)(x, y, z - k(t)) \\ & + a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\ & + a(x, y, z - k(t)) D_3(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) - a(x, y, z - k(t))^2 D_{1,3}(f)(x, y, z - k(t)) \\ & - 2 f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) D_3(a)(x, y, z - k(t))) \end{aligned}$$

$$\begin{aligned} G^{ty} = & \frac{1}{2} \frac{1}{a(x, y, z - k(t))^3} (v(t) (-a(x, y, z - k(t)) D_3(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\ & - a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\ & - a(x, y, z - k(t)) f(x, y, z - k(t)) D_{2,3}(a)(x, y, z - k(t)) \\ & + 2 f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) + a(x, y, z - k(t))^2 D_{2,3}(f)(x, y, z - k(t))) \end{aligned}$$

$$\begin{aligned}
G^{tz} = & -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^5} (v(t) (4 D_2(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 f(x, y, z - k(t)) \\
& - 4 D_2(a)(x, y, z - k(t)) a(x, y, z - k(t))^3 D_2(f)(x, y, z - k(t)) \\
& - 2 a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_{2, 2}(a)(x, y, z - k(t)) \\
& - 4 D_1(a)(x, y, z - k(t)) a(x, y, z - k(t))^3 D_1(f)(x, y, z - k(t)) \\
& + 4 D_1(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 f(x, y, z - k(t)) \\
& - 2 a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_{1, 1}(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t))^2 + v(t)^2 f(x, y, z - k(t))^3 D_1(a)(x, y, z - k(t))^2 \\
& + v(t)^2 f(x, y, z - k(t))^3 D_2(a)(x, y, z - k(t))^2 + 2 a(x, y, z - k(t))^4 D_{2, 2}(f)(x, y, z - k(t)) \\
& + 2 a(x, y, z - k(t))^4 D_{1, 1}(f)(x, y, z - k(t)) + v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t))^2))
\end{aligned}$$

$$\begin{aligned}
G^{xx} = & \frac{1}{4} \frac{1}{a(x, y, z - k(t))^4} \left(v(t)^2 D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \right. \\
& - 2 v(t)^2 D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + v(t)^2 f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 - 4 \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^3 D_3(f)(x, y, z - k(t)) \\
& - 4 v(t)^2 D_3(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 - v(t)^2 D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - v(t)^2 f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 \\
& + 4 \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& + 4 v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{3, 3}(f)(x, y, z - k(t)) \\
& - 4 v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{3, 3}(f)(x, y, z - k(t)) - 12 v(t)^2 f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t))^2 \\
& + 4 v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{3, 3}(a)(x, y, z - k(t)) \\
& - 4 v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{3, 3}(a)(x, y, z - k(t)) \\
& \left. - 8 v(t) D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) \right)
\end{aligned}$$

$$\begin{aligned}
& + 2 v(t)^2 D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& + 16 v(t)^2 D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& + 8 v(t) D_3(a)(x, y, z - k(t))^2 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) \Big)
\end{aligned}$$

$$\begin{aligned}
G^{xy} = & -\frac{1}{2} \frac{1}{a(x, y, z - k(t))^4} (v(t)^2 (D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) - f(x, y, z - k(t)) D_1(a)(x, y, z - k(t))) (\\
& -D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) + f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)))
\end{aligned}$$

$$\begin{aligned}
G^{xz} = & -\frac{1}{2} \frac{1}{a(x, y, z - k(t))^4} \left(-\left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^3 D_1(f)(x, y, z - k(t)) \right. \\
& + 2 v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{1,3}(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{1,3}(f)(x, y, z - k(t)) \\
& + 4 v(t)^2 D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + 4 v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - 6 v(t)^2 f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 D_3(f)(x, y, z - k(t)) a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t)) \\
& - v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& - v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& - v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{1,3}(a)(x, y, z - k(t)) \\
& + 2 v(t) D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& \left. + v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{1,3}(f)(x, y, z - k(t)) \right)
\end{aligned}$$

$$\begin{aligned}
G^{yy} = & \frac{1}{4} \frac{1}{a(x, y, z - k(t))^4} \left(v(t)^2 D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \right. \\
& \left. - 2 v(t)^2 D_2(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \right)
\end{aligned}$$

$$\begin{aligned}
& + v(t)^2 f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 - v(t)^2 D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - 4 \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^3 D_3(f)(x, y, z - k(t)) - 4 v(t)^2 D_3(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - v(t)^2 f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 \\
& + 4 \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& + 4 v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{3,3}(f)(x, y, z - k(t)) \\
& - 4 v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{3,3}(f)(x, y, z - k(t)) - 12 v(t)^2 f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t))^2 \\
& + 4 v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{3,3}(a)(x, y, z - k(t)) \\
& - 4 v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{3,3}(a)(x, y, z - k(t)) \\
& - 8 v(t) D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) \\
& + 2 v(t)^2 D_1(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& + 16 v(t)^2 D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& + 8 v(t) D_3(a)(x, y, z - k(t))^2 \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) \Big) \\
G^{yz} = & - \frac{1}{2} \frac{1}{a(x, y, z - k(t))^4} \left(- \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^3 D_2(f)(x, y, z - k(t)) \right. \\
& + 2 v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t)) D_{2,3}(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t))^2 D_{2,3}(f)(x, y, z - k(t)) \\
& + \left(\frac{d}{dt} v(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& + v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^3 D_{2,3}(f)(x, y, z - k(t)) \\
& - 2 v(t)^2 D_3(f)(x, y, z - k(t)) a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t)) \\
& + 4 v(t)^2 f(x, y, z - k(t)) a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& + 4 v(t)^2 D_3(f)(x, y, z - k(t)) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& \left. - 6 v(t)^2 f(x, y, z - k(t))^2 D_3(a)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \right)
\end{aligned}$$

$$\begin{aligned}
& -v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t)) D_3(a)(x, y, z - k(t)) \\
& -v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 D_3(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& -v(t) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t))^2 f(x, y, z - k(t)) D_{2,3}(a)(x, y, z - k(t)) \\
& + 2 v(t) D_3(a)(x, y, z - k(t)) \left(\frac{d}{dt} k(t) \right) a(x, y, z - k(t)) f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \Big)
\end{aligned}$$

$$\begin{aligned}
G^{zz} = & -\frac{1}{4} \frac{1}{a(x, y, z - k(t))^6} (v(t)^2 (4 f(x, y, z - k(t)) a(x, y, z - k(t))^4 D_{1,1}(f)(x, y, z - k(t)) \\
& + 4 f(x, y, z - k(t)) a(x, y, z - k(t))^4 D_{2,2}(f)(x, y, z - k(t)) \\
& - 4 f(x, y, z - k(t))^2 a(x, y, z - k(t))^3 D_{1,1}(a)(x, y, z - k(t)) \\
& + 11 f(x, y, z - k(t))^2 D_1(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 \\
& - 4 f(x, y, z - k(t))^2 a(x, y, z - k(t))^3 D_{2,2}(a)(x, y, z - k(t)) \\
& + 11 f(x, y, z - k(t))^2 D_2(a)(x, y, z - k(t))^2 a(x, y, z - k(t))^2 + 3 D_2(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^4 \\
& + 3 D_1(f)(x, y, z - k(t))^2 a(x, y, z - k(t))^4 + v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t))^2 D_2(f)(x, y, z - k(t))^2 \\
& + v(t)^2 f(x, y, z - k(t))^2 a(x, y, z - k(t))^2 D_1(f)(x, y, z - k(t))^2 + v(t)^2 f(x, y, z - k(t))^4 D_1(a)(x, y, z - k(t))^2 \\
& + v(t)^2 f(x, y, z - k(t))^4 D_2(a)(x, y, z - k(t))^2 \\
& - 14 D_2(f)(x, y, z - k(t)) a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& - 14 D_1(f)(x, y, z - k(t)) a(x, y, z - k(t))^3 f(x, y, z - k(t)) D_1(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 f(x, y, z - k(t))^3 a(x, y, z - k(t)) D_2(f)(x, y, z - k(t)) D_2(a)(x, y, z - k(t)) \\
& - 2 v(t)^2 f(x, y, z - k(t))^3 a(x, y, z - k(t)) D_1(f)(x, y, z - k(t)) D_1(a)(x, y, z - k(t)))
\end{aligned}$$

with $a(x, y, z - k(t)) \gg 1$ in the Pfenning zone there is again an arbitrary reduction of the components of the energy-impulse tensor and consequently of the energy density. The equations in the Eulerian system are presented in implicit form in appendix I, we only need to replace the equations in appendix II to find the same results. The same boundary conditions that apply when $v = cost$ apply also in this case.

BIBLIOGRAPHY:

- [1] M. Alcubierre, *Classical and Quantum Gravity* **11**, L73 (1994).
- [2] C. Barcelo, S. Finazzi, and S. Liberati, ArXiv e-prints (2010), arXiv:1001.4960 [gr-qc].
- [3] C. Clark, W. A. Hiscock, and S. L. Larson, *Classical and Quantum Gravity* **16**, 3965 (1999).
- [4] M. J. Pfenning and L. H. Ford, *Classical and Quantum Gravity* **14**, 1743 (1997).
- [5] F. S. N. Lobo and M. Visser, *Classical and Quantum Gravity* **21**, 5871 (2004).
- [6] F. S. N. Lobo, ArXiv e-prints (2007), arXiv:0710.4474 [gr-qc].
- [7] Finazzi, Stefano; Liberati, Stefano; Barceló, Carlos (2009). "Semiclassical instability of dynamical warp drives". *Physical Review D* **79** (12): 124017. arXiv:0904.0141
- [8] Van den Broeck, Chris (1999). "On the (im)possibility of warp bubbles". arXiv:gr-qc/9906050
- [9] C. Van Den Broeck, *Class. Quantum Grav.* 16 (1999) 3973
- [10] Hiscock, William A. (1997). "Quantum effects in the Alcubierre warp drive spacetime". *Classical and Quantum Gravity* **14** (11): L183–L188. arXiv gr-qc/9707024
- [11] Paper "Modified Alcubierre warp drive II", [viXra:1507.0193](https://arxiv.org/abs/1507.0193)
- [12] L D Landau and E M Lifshitz "The Classical Theory of Fields", Fourth Edition: Volume 2 (Course of Theoretical Physics Series)