# MATTER THEORY OF MAXWELL EQUATIONS 

WU SHENG-PING


#### Abstract

This article try to unified the four basic forces by Maxwell equations, the only experimental theory. Self-consistent Maxwell equation with the e-current from matter current is proposed. and is solved to four kinds of electrons and the structures of particles. The static properties and decay and scattering are reasoned, all meet experimental data. The equation of general relativity sheerly with electromagnetic field is discussed as the base of this theory. In the end the conformation elementarily between this theory and QED and weak theory is discussed compatible, except some bias in some analysis.


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## 1. Unit Dimension of sch

A rebuilding of units and physical dimensions is needed. Time $s$ is fundamental. We can define:
The unit of time: $s$ (second)
The unit of length: $c s$ ( $c$ is the velocity of light)
The unit of energy: $\hbar / s$ ( $h$ is Plank constant)

[^0]The unit dielectric constant $\epsilon$ is

$$
[\epsilon]=\frac{[Q]^{2}}{[E][L]}=\frac{[Q]^{2}}{\hbar c}
$$

The unit of magnetic permeability $\mu$ is

$$
[\mu]=\frac{[E][T]^{2}}{[Q]^{2}[L]}=\frac{\hbar}{c[Q]^{2}}
$$

We can define the unit of $Q$ (charge) as

$$
c \epsilon=c \mu=1
$$

then

$$
\begin{gathered}
{[Q]=\sqrt{\hbar}} \\
{[H]=[Q] /[L]^{2}=[c D]=[E]}
\end{gathered}
$$

Then

$$
\sqrt{\hbar}: C=\left(1.0546 \times 10^{-34}\right)^{1 / 2}
$$

$C$ is charge SI unit Coulomb.
For convenience we can define new base units by unit-free constants

$$
c=1, \hbar=1,[Q]=\sqrt{\hbar}
$$

then all physical unit are power of second $s^{n}$, the units are reduced.
Define

$$
\begin{gathered}
\text { UnitiveElectricalCharge }: \sigma=\sqrt{\hbar} \\
\sigma=1.03 \times 10^{-17} C \approx 64 e \\
e_{/ \sigma}=e / \sigma=1.57 \times 10^{-2} \approx 1 / 64
\end{gathered}
$$

The unit of charge can be reset by linear variation of charge-unit

$$
Q \rightarrow C Q, Q: \sigma / C
$$

We will use it without detailed explanation.

## 2. Self-consistent Electrical-magnetic Fields

Try equation for the free E-M field in mass-center frame

$$
\begin{gather*}
\partial \cdot \partial^{\prime} A=i A^{\nu *} \cdot \partial A_{\nu}=-J, \quad Q_{e}=1  \tag{2.1}\\
\partial_{\nu} \cdot A^{\nu}=0
\end{gather*}
$$

with definition

$$
\begin{gathered}
\left(A^{i}\right):=(V, \mathbf{A}),\left(J^{i}\right)=(\rho, \mathbf{J}),\left(J_{i}\right)=(-\rho, \mathbf{J}) \\
\partial:=\left(\partial_{i}\right):=\left(\partial_{t}, \partial_{x_{1}}, \partial_{x_{2}}, \partial_{x_{3}}\right) \\
\partial^{\prime}:=\left(\partial^{i}\right):=\left(-\partial_{t}, \partial_{x_{1}}, \partial_{x_{2}}, \partial_{x_{3}}\right)
\end{gathered}
$$

$Q_{e}$ is the absolute value of the charge of electron. It's deduced by using momentum to express e-current in a electron, and the rate between momentum and e-current is $k_{e}: e$. The equation 2.1 have symmetry

$$
C P T, c c . P T
$$

The energy of the field $A$ are

$$
\begin{equation*}
\varepsilon:=\int d V\left(E \cdot E^{*}+H \cdot H^{*}\right) / 2=-<A^{\nu}\left|\nabla^{2}\right| A_{\nu}>/ 2 \tag{2.2}
\end{equation*}
$$

It's gauge-invariant. As a convention the time-variant part is neglected.

## 3. Calculation of Recursive Re-Substitution

We can calculates the solution by recursive re-substitution for the two sides of the equation. For the equation

$$
\hat{P}^{\prime} B=\hat{P} B
$$

make the algorithm

$$
\hat{P}^{\prime}\left(\sum_{k \leq n} B_{k}+B_{n+1}\right)=\hat{P} \sum_{k \leq n} B_{k}
$$

One can write down a function initially and correct it by re-substitution. Here is the initial state

$$
V=V_{i} e^{-i k t}, A_{i}=V, \partial_{\mu} \partial^{\mu} A_{i}^{\nu}=0
$$

Substituting into equation 2.1

$$
\partial \cdot \partial^{\prime} A=i A^{\nu *} \partial A_{\nu}, Q_{e}=1
$$

with

$$
\partial_{\nu} \cdot A^{\nu}=0
$$

We calls the fields' correction $A_{n}$ with $n$ degrees of $A_{i}$ the n degrees correction. By the condition 2.1 to solve the dynamic process in the mass-center frame

$$
\begin{gathered}
\hat{P}^{\prime}\left(A-A_{i}\right)=\hat{P}\left(A-A_{i}\right) \\
\hat{P}(A):=F\left(i A^{\nu *} \partial A_{\nu}\right) \\
\hat{P}^{\prime}:=g_{i j} s^{i} s^{j}
\end{gathered}
$$

Then the dependence between initial state and later state is

$$
\begin{equation*}
<A_{\mu}\left|\hat{P}^{\prime} /\left(\hat{P}^{\prime}-\hat{P}\right)\right| A_{i}^{\mu}> \tag{3.1}
\end{equation*}
$$

because

$$
\int d A<A_{\mu}, \hat{P} /\left(\hat{P}^{\prime}-\hat{P}\right) A_{i}^{\mu}>=0
$$

The fraction of operator is comprehended as recursive re-substitution.
The decay to a stable state is calculated. The system is isolated,

$$
\begin{gathered}
<\nabla A^{\nu}, \nabla A_{\nu}>=<\partial_{t} A^{\nu}, \partial_{t} A_{\nu}> \\
A=A_{i t}+A_{f t}
\end{gathered}
$$

hence

$$
\begin{gathered}
-\partial_{t}\left(<\left(A_{i t}\right)^{\nu}, \nabla^{2}\left(A_{i t}\right)_{\nu}>+<\left(A_{f t}\right)^{\nu}, \nabla^{2}\left(A_{f t}\right)_{\nu}>\right) \\
=\partial_{t}\left(<\partial_{t}\left(A_{i t}\right)^{\nu}, \partial_{t}\left(A_{i t}\right)_{\nu}>+<\partial_{t}\left(A_{f t}\right)^{\nu}, \partial_{t}\left(A_{f t}\right)_{\nu}>\right) \\
=-\partial_{t}<\left(A_{i t}\right)^{\nu}, \partial_{t}^{2}\left(A_{f t}\right)_{\nu}>+c c \\
=\partial_{t}^{2}<A^{\nu}, \partial_{t} A_{\nu}>+c c
\end{gathered}
$$

and with

$$
\partial_{t}<\partial_{t}(\Re(A)-\Im(A))^{\nu},(\Re(A)-\Im(A))_{\nu}>=0
$$

and taking the charge normalization of $A$, to find

$$
\begin{aligned}
& e^{-C t}=\varepsilon / \varepsilon_{i} \\
& C=\varepsilon_{i}-\varepsilon_{f}
\end{aligned}
$$

## 4. Solution

Firstly

$$
\nabla^{2} A=-k^{2} A
$$

is solved. Exactly, it's solved in spherical coordinate

$$
0=r^{2}\left(\nabla^{2} f+k^{2} f\right)=k^{2} r^{2} f+\left(r^{2} f_{r}\right)_{r}+\frac{1}{\sin \theta}\left(\sin \theta f_{\theta}\right)_{\theta}+\frac{1}{\sin ^{2} \theta}\left(f_{\phi}\right)_{\phi}
$$

Its solution is

$$
\begin{gathered}
f=R \Theta \Phi=R_{l} Y_{l m} \\
\Theta=P_{l}^{m}(\cos \theta), \Phi=\cos (\alpha+m \phi) \\
R_{l}=N j_{l}(k r)
\end{gathered}
$$

$j_{l}(x)$ is spherical Bessel function.

$$
\begin{gathered}
j_{1}(x)=\frac{\sin (x)}{x^{2}}-\frac{\cos x}{x} \\
j_{1}(0)=0
\end{gathered}
$$

Contrary to the well-known result:

$$
\int_{0}^{\infty} x^{2} j_{1}(a x) j_{1}(b x) d x=\frac{1}{a} \delta(a-b)
$$

the functions $j_{1}(a x), j_{1}(b x)$ are not orthogonal (why? CV), because a direct calculation shows that.

The solution of $l=1, m=1, Q=e_{/ \sigma}$ is calculated or tested for electron,

$$
V=-N R_{1}(k r) Y_{1,-1} e^{-i k t}
$$

## 5. Electrons and Their Symmetries

Some states of electrical field $A$ are defined as the core of the electron, it's the initial function $A_{i}=V$ that is electrical, for the re-substitution to get the whole electron function:

$$
\begin{gathered}
e_{r}^{+}: N R_{1}\left(k_{e} r\right) Y_{1,1} e^{i k_{e} t} \\
e_{l}^{+}: N R_{1}\left(k_{e} r\right) Y_{1,-1} e^{i k_{e} t} \\
e_{r}^{-}:-N R_{1}\left(k_{e} r\right) Y_{1,-1} e^{-i k_{e} t} \\
e_{l}^{-}:-N R_{1}\left(k_{e} r\right) Y_{1,1} e^{-i k_{e} t}
\end{gathered}
$$

$r, l$ is the direction of the spin. We use these symbols $e$ to express the complete potential field $A$ or the abstract particle.

By mainly the second rank of correction $A_{2}$, a static field, we have

$$
\begin{gathered}
<e_{c \mu}\left|-i \partial_{t}\right| e_{c}^{\mu}>\approx Q_{c}, \sigma=1,\left|Q_{c}\right|=Q_{e} \\
\quad-<\nabla e_{c \mu} \mid \nabla e_{c}^{\mu}>\approx k_{e} e_{/ \sigma}, \sigma=1
\end{gathered}
$$

We can use this to normalize the electron functions.
The magnetic dipole moment of electron is calculated as the first rank of proximation

$$
\begin{gathered}
\mathbf{r} \times \partial \cdot \partial^{\prime} A / 2 \\
\mu_{z}=<A_{i}\left|-i \partial_{\phi}\right| A_{i}>/ 2 \\
=\frac{Q_{e}}{2 k_{e}}, \sigma=1
\end{gathered}
$$

The spin is

$$
\begin{gathered}
S_{z}=<A\left|\partial_{t} \partial_{\phi}\right| A>/ 2, Q_{e}=1 \\
=1 / 2
\end{gathered}
$$

The correction of the equation 2.1 is

$$
\begin{gathered}
A_{n}=A_{n-1} \cdot\left(\partial\left(i A_{i}-i A_{i}^{*}\right) / 2\right) * u \\
=A_{i}^{*}\left(i \partial_{t} A_{i}\right)\left(\partial_{t}\left(i A_{i}-i A_{i}^{*}\right) / 2\right)^{n-3} \partial\left(i A_{i}-i A_{i}^{*}\right) / 2 \\
u=\delta(t-r)) /(4 \pi r)
\end{gathered}
$$

The convolution is made in 4 -d space.
The function of $e_{r}^{+}$is decoupled with $e_{l}^{+}$

$$
\begin{aligned}
& \quad<\nabla\left(e_{r}^{+}\right)^{\nu}+\nabla\left(e_{l}^{+}\right)^{\nu}, \nabla\left(e_{r}^{+}\right)_{\nu}+\nabla\left(e_{l}^{+}\right)_{\nu}>/ 2 \\
& -<\nabla\left(e_{r}^{+}\right)^{\nu}, \nabla\left(e_{r}^{+}\right)_{\nu}>/ 2-<\nabla\left(e_{l}^{+}\right)^{\nu}, \nabla\left(e_{l}^{+}\right)_{\nu}>/ 2=0
\end{aligned}
$$

The increment of field energy $\varepsilon$ on the coupling of $e_{r}^{+}, e_{r}^{-}$mainly between $A_{2}$ is

$$
\varepsilon_{e} \approx-e_{/ \sigma}^{3} k_{e}=-\frac{1}{1.66 \times 10^{-16} s}
$$

This value of increments on the coupling of electrons are

| $\varepsilon_{e}$ | $e_{r}^{+}$ | $e_{r}^{-}$ | $e_{l}^{+}$ | $e_{l}^{-}$ |
| :---: | :---: | :---: | :---: | :---: |
| $e_{r}^{+}$ | + | - | 0 | 0 |
| $e_{r}^{-}$ | - | + | 0 | 0 |
| $e_{l}^{+}$ | 0 | 0 | + | - |
| $e_{l}^{-}$ | 0 | 0 | - | + |

The increment of field energy $\varepsilon$ on the coupling of $e_{r}^{+}, e_{l}^{-}$mainly between $A_{4}$ is

$$
\varepsilon_{x} \approx-\frac{1}{4} e_{/ \sigma}^{7} k_{e} \approx-\frac{1}{2.18 \times 10^{-8} s}
$$

The calculations referenced to latter theorems. This value of increments on the coupling of electrons are

$$
\begin{array}{ccccc}
\varepsilon_{x} & e_{r}^{+} & e_{r}^{-} & e_{l}^{+} & e_{l}^{-} \\
e_{r}^{+} & + & 0 & 0 & - \\
e_{r}^{-} & 0 & + & - & 0 \\
e_{l}^{+} & 0 & - & + & 0 \\
e_{l}^{-} & - & 0 & 0 & +
\end{array}
$$

## 6. Anti-Particle Symmetry

Antimatter causes by the symmetry $P T$, we define

$$
\overline{A(x)}:=A(-x)
$$

Hence the core of electron functions made by $P T$ are

$$
\begin{gathered}
\overline{e_{r}^{+}}:-N R_{1}\left(k_{e} r\right) Y_{1,1} e^{-i k_{e} t} \\
\overline{e_{l}^{+}}:-N R_{1}\left(k_{e} r\right) Y_{1,-1} e^{-i k_{e} t} \\
\overline{e_{r}^{-}}: N R_{1}\left(k_{e} r\right) Y_{1,-1} e^{i k_{e} t} \\
\overline{e_{l}^{-}}: N R_{1}\left(k_{e} r\right) Y_{1,1} e^{i k_{e} t}
\end{gathered}
$$

We find

$$
\overline{e_{r}^{+}}=e_{l}^{-}
$$

The both are the same particle except strong coupling, which evinces that antiparticle has meaning beyond Maxwell Equations. The equation 2.1 reserves of cause.

## 7. Propagation and Movement

Because field $F$ is additive, the group of electrons are express by:

$$
F=\sum_{i} f_{i} * \nabla e_{i},<f_{i} \mid f_{i}>=1
$$

It's called propagation. The convolution is made only in space:

$$
f * g=\int d^{3} x f(t, x) g(t, y-x)
$$

Each $f_{i}$ is normalized to 1 . We always use

$$
\sum_{i} f_{i} * e_{i}, \sum_{i} f_{i} * \nabla e_{i}
$$

to express its abstract construction and the field. The following are stable propagation:

| particle | electron | photon | neutino |
| :---: | :---: | :---: | :---: |
| notation | $e_{r}^{+}$ | $\gamma_{r}$ | $\nu_{r}$ |
| structure | $e_{r}^{+}$ | $\left(e_{r}^{+}+e_{r}^{-}\right)$ | $\left(e_{r}^{+}+e_{l}^{-}\right)$ |

Define

$$
\begin{aligned}
\varsigma_{k, l, m}(x):= & R_{l}(k r) Y_{l, m}, \varsigma_{k}=\varsigma_{k}^{ \pm}(x):=\varsigma_{k, 1, \pm 1}(x) \\
& <\varsigma_{k}(x), \varsigma_{k}(x)>=\delta(0)
\end{aligned}
$$

It meets the following results

## Theorem 7.1.

$$
\int d^{3} \mathbf{x} R\left(\varsigma_{k}^{ \pm}(x)\right) \varsigma_{k}^{*}(y-x)=0, y \neq O
$$

$R$ is any rotation.
Proof. Use the limit

$$
\lim _{k^{\prime} \rightarrow k} \int d V \varsigma_{k}^{ \pm}(x) \varsigma_{k^{\prime}}^{*}(x-y)
$$

and the identity

$$
h \nabla^{2} g-g \nabla^{2} h=\nabla \cdot(h \nabla g-g \nabla h)
$$

Theorem 7.2.

$$
\begin{gathered}
\varsigma_{1} * \frac{1}{4 \pi r}=\varsigma_{1} \\
\left.\varsigma_{1} e^{i t} * \delta(t-r)\right) /(4 \pi r)=\varsigma_{1} e^{i t} \\
\varsigma_{1}^{n} \varsigma_{1}^{* n^{\prime}} * \frac{1}{4 \pi r}=\varsigma_{1}^{n} \varsigma_{1}^{* n^{\prime}} \\
\left.\varsigma_{1}^{n} e^{i n t} \varsigma_{1}^{* n^{\prime}} e^{-i n^{\prime} t} * \delta(t-r)\right) /(4 \pi r)=\varsigma_{1}^{n} e^{i n t} \varsigma_{1}^{* n^{\prime}} e^{-i n^{\prime} t}
\end{gathered}
$$

Proof. Calculate

$$
\begin{gathered}
\nabla^{2}\left(\varsigma_{1} * \frac{1}{4 \pi r}\right)=\nabla^{2} \varsigma_{1} \\
\left.\nabla^{2}\left(\varsigma_{1} e^{i t} * \delta(t-r)\right) /(4 \pi r)\right)=\nabla^{2}\left(\varsigma_{1} e^{i t}\right) \\
F_{\lambda_{i}}\left(\varsigma_{1}^{n} \varsigma_{1}^{* n^{\prime}} * \frac{1}{4 \pi r}\right) * e^{i \lambda_{i} x_{i}}
\end{gathered}
$$

and use the identity

$$
\delta(t-r) /(4 \pi r)=\delta(t) /(4 \pi r) * \delta(t-r)
$$

## Theorem 7.3.

$$
\nabla\left(\varsigma_{k} * \varsigma_{k^{\prime}}\right)=\left(\nabla \varsigma_{k}\right) * \varsigma_{k^{\prime}}+\varsigma_{k} * \nabla\left(\varsigma_{k^{\prime}}\right)
$$

Use this condition to prove:

$$
\varsigma_{k}(x-c) * \varsigma_{k^{\prime}}(x+c)=C \varsigma_{k}(x-c) \varsigma_{k^{\prime}}(x+c) * \delta(x-2 c)
$$

## Theorem 7.4.

$$
<\varsigma_{1}^{n}, \varsigma_{1}^{n}>=<\varsigma_{1}, \varsigma_{1}>
$$

Prove in Fourier space.

## Theorem 7.5.

$$
\left(\nabla \varsigma_{k}\right) * \varsigma_{1}=k \varsigma_{k} * \nabla \varsigma_{1}
$$

The movement of the propagation is called Movement, ie. the third level wave:

$$
F=f * \sum f_{i} * \nabla e_{i}
$$

The static MDM (magnetic dipole moment) of a group of electrons is

$$
\begin{gathered}
\mathbf{r} \times \partial \cdot \partial^{\prime}\left(\int d \mathbf{x} \cdot e_{x} * \sum_{j} \nabla e_{j}^{\mu}\left(x_{i}\right)\right) / 2 \\
\mu=<\int d \mathbf{x} \cdot e_{x} * \sum_{j} \nabla e_{j}^{\mu}\left(x_{i}\right)|i \mathbf{r} \times \partial| \int d \mathbf{x} \cdot e_{x} * \sum_{j} \nabla e_{j \mu}\left(x_{i}\right)>/ 2, Q_{e}=1
\end{gathered}
$$

MDM is coupling between electrons
Calculating the following coupling system

$$
e_{x} * \nabla A, A=\sum_{i} e_{i}
$$

with the condition 2.1

$$
\begin{align*}
& \partial \cdot \partial^{\prime} e_{x} \approx 0 \\
& e_{x}=e^{i N t} \varsigma_{N} \tag{7.1}
\end{align*}
$$

With the condition of charge

$$
\begin{gather*}
<\int d \mathbf{x} \cdot e_{x} * \nabla A^{\nu}\left|-i \partial_{t}\right| \int d \mathbf{x} \cdot e_{x} * \nabla A_{\nu}>=Q_{x}, Q_{e}=1 \\
N \approx<A^{\nu}, A_{\nu}>/ Q_{x}, Q_{e}=1 \tag{7.2}
\end{gather*}
$$

## 8. Conservation Law and Balance Formula

No matter in E-M fields level or in movement (the third) level, the conservation law is conservation of momentum and conservation of angular momentum. A balance formula for a reaction is the equivalent formula in positive matter, ie. after all antimatter is shifted to the other side of the reaction formula. Balance formula is suitable for the analysis of the energy transition of decay. The invariance of electron itself in reaction is also a conservation law.

## 9. Muon

Generally, there are kinds of energy increments.
Weak coupling

$$
W:<\nabla\left(e_{r}^{+}\right)^{\nu} \mid \nabla\left(e_{l}^{-}\right)_{\nu}>
$$

Light coupling

$$
L:<\nabla\left(e_{r}^{+}\right)^{\nu} \mid \nabla\left(e_{r}^{-}\right)_{\nu}>
$$

Weak side coupling

$$
W s:<e_{x} * \nabla\left(e_{r}^{+}\right)^{\nu}\left|e_{x} * \nabla\left(e_{l}^{-}\right)_{\nu}>-<e_{x} * \nabla\left(e_{r}^{+}\right)^{\nu}\right| e^{i p x} * \nabla\left(e_{l}^{-}\right)_{\nu}>
$$

Light side coupling

$$
L s:<e_{x} * \nabla\left(e_{r}^{+}\right)^{\nu}\left|e_{x} * \nabla\left(e_{r}^{-}\right)_{\nu}>-<e_{x} * \nabla\left(e_{r}^{+}\right)^{\nu}\right| e^{i p x} * \nabla\left(e_{r}^{-}\right)_{\nu}>
$$

Strong coupling

$$
S:<\nabla\left(e_{r}^{+}\right)^{\nu} \mid \nabla\left(e_{r}^{+}\right)^{\nu}>
$$

$\mu$ is composed of

$$
\mu_{r}^{+}: e_{\mu} *\left(e_{r}^{+}+\overline{\nu_{r}}\right)
$$

From the equation 15.1, 7.2 and 7.1 and the other deductive, $\mu$ is with mass $3 m_{e} / e_{/ \sigma}=3 \times 64 k_{e}$, $\operatorname{spin} 1 / 2$, MDM $\mu_{B} k_{e} / k_{\mu}$.

The main channel of decay

$$
\begin{gathered}
\mu_{r}^{+} \rightarrow M_{l}^{+}+\overline{\nu_{l}} \\
M_{r}^{+}=e_{M} *\left(\overline{e_{l}^{-}}+\nu_{l}\right) \\
e_{\mu} * e_{r}^{+}+e^{-i p_{1} x} * e_{M}^{*} * e_{l}^{-}+\overline{e^{i p_{2} x}} * \nu_{l} \rightarrow \overline{e_{\mu}} * \nu_{r}+\overline{e^{i p_{1} x}} * \overline{e_{M}} \nu_{l}
\end{gathered}
$$

The outer waves $e_{\mu}$ and $e^{-i p_{1} x} * e_{M}^{*}, e_{\mu}$ and $e^{-i p_{2} x}$ are coupling. The energy difference is kind of $W s$, the interacting field is between $A_{4}$.

$$
2<e_{\mu} * \nabla\left(e_{r}^{+}\right)^{\nu}\left|e^{-i p_{1} x} * \nabla\left(e_{l}^{-}\right)_{\nu}>-2<e_{\mu} * \nabla\left(e_{r}^{+}\right)^{\nu}\right| e_{\mu} * \nabla\left(e_{l}^{-}\right)_{\nu}>+c c
$$

Sum up in spectrum of $p_{1}$ that with the emission positive:

$$
\begin{gathered}
\left.=2 \frac{k_{e}}{k_{\mu}+k_{e}}<\nabla\left(e_{\mu} *\left(e_{r}^{+}, A_{4}\right)^{\nu}\right)\left|2 e^{i n t} \nabla\left(e_{l}^{-}, A_{4}\right)_{\nu} / k_{e}>-2<e_{\mu} * \nabla e_{r}^{+}\right| e_{\mu} * \nabla\left(e_{l}^{-}\right)_{\nu}\right)>+c c . \\
=\left(2 \frac{k_{\mu}}{k_{\mu}+k_{e}}-2\right) \varepsilon_{x} \\
=-\frac{2 \varepsilon_{x} k_{e}}{k_{\mu}+k_{e}}
\end{gathered}
$$

The emission of decay is

$$
=\frac{1}{2.1 \times 10^{-6} s} \quad\left[2.1970 \times 10^{-6} s\right][1]
$$

The data in square bracket is experimental data. The decay of particle $M$ is like a scattering with no energy emission

$$
M_{r}^{+} \rightarrow \overline{e_{l}^{-}}+\nu_{l}
$$

## 10. Pion Positive

Pion positive is

$$
\pi_{l}^{-}: e_{\pi} *\left(\overline{e_{r}^{+}}+e_{l}^{-}\right)+e_{\pi}^{*} * e_{r}^{+}
$$

It's with mass $5 \times 64 m_{e}$, spin $1 / 2$ and MDM $\mu_{B} k_{e} / k_{\pi^{+}}$.
Decay Channels:

$$
\pi_{l}^{-} \rightarrow \mu_{l}^{-}+\nu_{r}
$$

It's with balance formula

$$
e_{\pi}^{*} * e_{r}^{+}+e_{\pi} * e_{l}^{-}+\overline{e^{i p_{1} x}} * \overline{e_{\mu}} * \nu_{r} \rightarrow \overline{e_{\pi}} * e_{r}^{+}+e^{i p_{1} x} * e_{\mu} * e_{l}^{-}+e^{i p_{2} x} * \nu_{r}
$$

The the emission of energy, ie. the cross energy, is kind of $W$

$$
-\varepsilon_{x}=\frac{1}{2.18 \times 10^{-8} s} \quad\left[\left(2.603 \times 10^{-8} s\right][1]\right.
$$

## 11. Pion Neutral

Pion neutral is atom-like particle

$$
\pi^{0}: e_{\pi^{0}} *\left(e_{r}^{+}+e_{l}^{+}\right)+e_{\pi^{0}}^{*} *\left(e_{r}^{-}+e_{l}^{-}\right)
$$

It has mass $4 \times 64 m_{e}$, zero spin and zero MDM. Its decay modes are

$$
\pi^{0} \rightarrow \gamma_{r}+\gamma_{l}
$$

The loss of energy is kind of $L$

$$
-2 \varepsilon_{e}=\frac{1}{8.3 \times 10^{-17} s} \quad\left[8.4 \times 10^{-17} s\right][1]
$$

> 12. TAU
$\tau$ maybe that

$$
\tau_{l}^{+}: e_{\tau} *\left(5 e_{r}^{+}+\overline{5 e_{r}^{+}+e_{r}^{-}}\right)
$$

Its mass $51 \times 64 m_{e}$, spin $1 / 2, \operatorname{MDM} \mu_{B} k_{e} / k_{\tau}$. It has decay mode

$$
\begin{aligned}
\tau_{l}^{+} & \rightarrow \mu_{l}^{+}+\nu_{l}+\overline{\nu_{l}} \\
e_{\tau} * 5 e_{r}^{+}+\overline{e^{i p_{1} x}} * \overline{e_{\mu}} * \nu_{l}+\overline{e^{i p_{2} x}} * \nu_{l} & \rightarrow \overline{e_{\tau}} * 5 e_{r}^{+}+\overline{e_{\tau}} * e_{r}^{-}+e^{i p_{1} x} * e_{\mu} * e_{l}^{+}+e^{i p_{3} x} * \nu_{l}
\end{aligned}
$$

The energy gap is kind of $L s$

$$
\begin{aligned}
& 5<\overline{e_{\tau}} * \nabla\left(e_{r}^{+}\right)^{\nu} \mid e^{-i p_{1} x} * \nabla\left(e_{r}^{-}\right)_{\nu}>-5<\overline{e_{\tau}} * \nabla\left(e_{r}^{+}\right)^{\nu} \mid \overline{e_{\tau}} * \nabla\left(e_{r}^{-}\right)_{\nu}>+c c . \\
& \approx-\frac{5 \varepsilon_{e}}{k_{\tau} / k_{e}} \\
&=\frac{1}{1 \times 10^{-13} s} \quad\left[2.9 \times 10^{-13} s ; B R . \quad 0.17\right][1]
\end{aligned}
$$

Depending on this kinds of particle including

$$
q_{r}^{n+}:=n\left(e_{r}^{+}+\overline{e_{r}^{+}}\right)
$$

we can construct particles of great mass decaying without strong emission (light radiative), for example

$$
e_{L} *\left(q_{r}^{n+}+\overline{e_{l}^{-}}\right)
$$

This series of particle include $\mu, \tau$ and in fact almost all light radiative particles are of this kind, they are created in colliding. Another condition is possibly that, in the collision, the created light radioactive particle $\left(q_{r}^{n+}, \overline{e_{l}^{-}}\right)$with different $n$ is mixed to some rates as to the detector can't distinguish them.

Because the channel width decides the channel branch rates, obviously the most experimental data violate this rule. So that the channels listing after the same name of a particle in fact belong to different particles.

The particle $K^{+}$possibly is

$$
K^{+}=\left(q_{r}^{3+}+\overline{\nu_{l}+e_{l}^{-}}\right) \rightarrow \mu_{r}^{+}+\overline{\nu_{l}}
$$

It has emission of $L s$.

## 13. Proton

Proton may be like
$p_{l}^{-}: e_{p} *\left(e_{r}^{+}+e_{r}^{-}+e_{l}^{+}+e_{l}^{-}+\overline{e_{r}^{+}+e_{r}^{-}+2 e_{l}^{+}}\right)+e_{p}^{*} *\left(\overline{e_{r}^{+}+e_{r}^{-}+e_{l}^{+}+e_{l}^{-}}+e_{r}^{+}+e_{l}^{+}+e_{l}^{-}\right)$
The mass is $31 \times 64 m_{e}$ that's very close to the real mass. The MDM is calculated as $3 \mu_{N}$, spin is $1 / 2$. The proton thus designed is eternal because if it decay even to the finest blocks the energy of emission is negative.

## 14. Scattering

The scattering can be calculated as dynamic electromagnetic mechanical theory, ie. the magnitude of scattering is calculated by the formula 3.1. the first correction is

$$
2 j^{\mu}=i e_{/ \sigma} A_{f \mu}^{*} \partial A_{i}^{\mu}-i e_{/ \sigma} A_{f}^{\mu} \partial A_{i \mu}^{*}
$$

The reaction is like

$$
\begin{gathered}
e^{i p_{1} x} * e+e^{i p_{2} x} * e \rightarrow e^{i p_{3} x} * e+e^{i p_{4} x} * e \\
A_{i}=A_{1}+A_{2}, A_{f}=A_{3}+A_{4}
\end{gathered}
$$

the transfer is

$$
<\left(j_{f \mu}\left|1 / \partial^{2}\right| j^{\mu}>=i \mu \approx \frac{C\left(p_{1}^{\prime}+p_{2}^{\prime}\right)^{\nu}\left(p_{3}^{\prime}+p_{4}^{\prime}\right)_{\nu}}{2\left(p_{1}^{\prime}-p_{3}^{\prime}\right)^{2}}\right.
$$

The $p_{i}^{\prime}$ is the cap momentum relative to $p_{i}$.

$$
i C=e^{2}=\frac{\varepsilon_{e}}{k_{e} e}
$$

The difference with the classical theory caused by different normalization, the unitary volume has particles $k$. After the sum of spins, This calculation is the same result of QED.

## 15. Great Unification

The mechanic feature of the electromagnet fields is

$$
T_{i j}=F_{i}^{k *} F_{k j}-g_{i j} F_{\mu \nu} F^{\mu \nu^{*}} / 4
$$

$T$ is stress-energy tensor,

$$
T_{i j}=\sum m \beta_{i} \beta_{j}, \beta_{0}=1, \beta=v / c
$$

$T_{00}$ is quantum expression of the energy, by Lorentz transform it's easy to get the quantum expression of momentum. The General Theory of Relativity is

$$
\begin{equation*}
R_{i j}-\frac{1}{2} R g_{i j}=8 \pi G T_{i j} / c^{4} \tag{15.1}
\end{equation*}
$$

Firstly we redefine the unit second as $S$ to simplify the equation 15.1

$$
R_{i j}-\frac{1}{2} R g_{i j}=T_{i j}, Q_{e}=1
$$

We observe that the co-variant curvature is

$$
R_{i j}=F_{i k}^{*} F_{j}^{k}+g_{i j} F_{\mu \nu}^{*} F^{\mu \nu} / 8, Q_{e}=1
$$

## 16. Conclusion

Fortunately this model explained all the effects in the known world: strong, weak and electromagnetic effects, and even subclassify them further if not being to add new ones. In this model the only field is electromagnetic field, and this stands for the philosophical with the point of that unified world is from an unique source. All depend on a simple fact: the current of matter in a system is time-invariantly zero in mass-center frame, then we can devise current of matter to analysis the e-charge current.

Except electron function my description of particles in fact is compatible with Quantum Electromagnetic Mechanics, and they two should reach the similar result. But my theory isn't compatible to the theory of quarks, if not it is calculated in the style of Quantum Electromagnetic Mechanics. In fact, The electron function is a good promotion for the experimental model of proton that went up very early.

Underlining my calculations a fact is that the electron has the same phase (electron resonance), which the BIG BANG theory would explain, all electrons are generated in the same time and place, the same source.

## References

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E-mail address: sunylock@139.com
8 Hanbei Rd, Tianmen, Hubei Province, The People's Republic of China. Postcode: 431700


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