#### Interpret Quantum Mechanics by uncertain complex waves

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# Abstract

This paper suggests a quantum reality which are uncertain complex waves. These waves are described completely by an axiomatic system. This axiomatic system solves the measurement problem elegantly. From the uncertain complex waves, I revalue the classical picture.

## I. Introduction

Quantum mechanics has a long history and it has been used to explain many experiment results. Quantum mechanics passed through countless testing experiment [1][2], even the most rigorous experiment [3]. However, the controversy of the foundation of quantum mechanics has been never stopped [1][2]. The most outstanding debate is how the wave function collapse happen [4][5]? What properties are observable in a particular experiment [6]? What is a quantum reality [7]? Many works were done to answer these questions [8]. However, no approach is agreed widely because each has own difficulty [8].

If axioms about the operator in orthodox quantum mechanics are removed then the Schrodinger equation can indicate the observable properties by itself. So in this paper, I suggest to modify the axiomatic system of the orthodox quantum mechanics. Axioms about operators are removed. The observable properties are indicated to operators which their eigenvalues appear in solutions of the Schrodinger equation.

On the other hand, if we know a general rule about possible states before and after measurement then we can understand the wave function collapse. So I propose an axiom about the possible states. This suggestion shows that discrete potential energy makes the change of state uncertain and suddenly. The wave function collapse is a simple consequence of this process.

From suggested modification, the quantum reality are regarded as uncertain complex waves. The quantum reality is counterintuitive but it can drive the classical picture by averaging. On other words, the classical reality isn't regarded as real reality.

# II. Theory and discussion

## The uncertain complex waves

We know that the world is composed of micro-particles. So I believe that it exits a quantum reality of micro-particles. N. Bohr considered the quantum reality is counterintuitively [9]. I suggest a viewpoint which the quantum reality are uncertain complex waves. In this paper, I give an axiomatic system to describe completely these waves. The axiomatic system of the orthodox quantum mechanics is stated by E. G. Harris and it is showed in [10]. It is modified to become a new axiomatic system for uncertain complex waves. The new axiomatic system includes two axioms from Harris's statement.

They are the Born rule and the Schrodinger equation. Axioms about operators of the orthodox quantum mechanics are removed. The axiom about vector in Hilbert space is modified to become new form. And I proposed an axiom about possible states. It can be used for observed system as well as any other case. The first axiom in the suggested axiomatic system is:

Quantum system is described by a set of vectors which is in Hilbert space. They are called state vectors or wave functions of system. The vector  $\psi$  and  $\lambda$ .  $\psi$  ( $\lambda$  is a complex number) describe the same state. In general,  $\psi$  is normalized to the unit.

Here, quantum system is described by not only vector in Hilbert space as the orthodox quantum mechanics. The first axiom requires many vectors for description. These vectors are determined by axiom about possible states. Now, in the Hilbert space, we can define operators. Example, coordinates operators:  $\hat{x}=x$ ;  $\hat{y}=y$ ;  $\hat{z}=z$ . Momentum operators:  $p_x=-i$ .  $\hbar$ .  $\frac{d}{dx}$ ;  $p_y=-i$ .  $\hbar$ .  $\frac{d}{dy}$ ;  $p_z=-i$ .  $\hbar$ .  $\frac{d}{dz}$ . Hamiltonian operator for single particle system:  $H=\frac{p_x^2+p_y^2+p_z^2}{2.m}+V(x,y,z,t)$ . The quantity V(x,y,z,t) is potential energy operator. It is a function of coordinates and time. m is mass of the particle. Relativity Hamiltonian operator for single particle has form:  $H=\alpha_x.p_x+\alpha_y.p_y+\alpha_z.p_z+\beta.m.c^2$ . Here  $\alpha_x=\begin{pmatrix}0&\sigma_x\\-\sigma_x&0\end{pmatrix}$ ,  $\alpha_y=\begin{pmatrix}0&\sigma_y\\-\sigma_y&0\end{pmatrix}$ ,  $\alpha_z=\begin{pmatrix}0&\sigma_z\\-\sigma_z&0\end{pmatrix}$ ,  $\beta=\begin{pmatrix}l_2&0\\0&-l_2\end{pmatrix}$ .  $\sigma_x,\sigma_y,\sigma_z$  are Pauli matrices.  $I_2$  is unit matrix  $\begin{pmatrix}1&0\\0&1\end{pmatrix}$ . The second axiom in E. G. Harris's statement is:

Wave functions of system satisfy the Schrodinger equation: i.  $\hbar$ .  $\frac{d}{dt}\psi = H\psi$ 

The wave functions also satisfy two packs of conditions. The first are initial condition and boundary conditions which are Dirichlet and Neumann conditions. The first pack is called tight condition. The second pack includes integrable square, single-valued and continuous conditions. They pack is called open condition. Wave function which satisfies tight condition exists only. It is called tight solution, tight wave function or tight state vector. Wave functions which satisfy the open condition are called open solutions, open wave functions or open state vectors.

If the potential energy operator is discontinuous then the Schrodinger equation must be investigated in many space domains. We divide space into the as large as possible domains  $D_i$  (i=1,2,...) in which the potential energy is continuous in each domain. These domains are called continuous domains of the system. The Schrodinger equation has one the largest symmetry group  $G_i$  in each domain  $D_i$ . It is called basic group of the system in  $D_i$ . Open solutions of the Schrodinger equation in  $D_i$  is a set  $U_i$  which is a linear space.  $T_{in}$  (i=1,2,..;n=1,2,..) is an irreducible representation of  $G_i$  in  $U_i$ . In the space of open solutions,  $U_{in}$  is a subspace which is invariant with  $T_{in}$ . Space  $U_{in}$  is called irreducible space of  $G_i$ . The axiom about the possible states is suggested by following:

Possible states of system in a continuous domain  $D_i$  are vectors which belong to irreducible spaces  $U_{in}$  of the basic group  $G_i$ .

The basis of each space  $U_{in}$  are orthogonal, the basis of each space  $U_i$  are too. From the first axiom, we can consider these bases are orthonormal.  $\psi_{in}$  is a vector in  $U_{in}$ .  $\psi_i$  is a vector in space  $U_i$ , we have:  $\psi_i = \sum_n c_{in} \cdot \psi_{in}$ . Vector  $\psi_i$  which is an open solution of the Schrodinger equation, in general, doesn't belong to any irreducible space of  $G_i$ . So  $\psi_i$  isn't a possible state. So the axiom about the possible states and the superposition principle aren't compatible.

With stationary systems, each space  $U_{in}$  corresponds to a stationary energy level. Vectors in  $U_{in}$  correspond to the same energy level. The degenerative degree of this energy level equals the number of dimensions of  $U_{in}$ . Of course, with a stationary system, combinations of two possible states which aren't the same energy level are not a possible state. Note  $D_i$  is a 4-dimensions domain. So the axiom about possible states is also used for time-dependent possible states.

If  $\psi$  is the tight solution, it can be expanded:  $\psi = \sum_n c_{in}. \psi_{in}.$  The Born rule is stated:

In general, in a continuous domain, system doesn't belong any particular possible state. The possibility of each possible state  $\psi_{in}$  is  $|c_{in}|^2$ .

Here, I suggest that the uncertainty is natural property of quantum system. And it doesn't depend on restriction of human and apparatus.

# 2. Observable properties

We can divide the measurement into two types. The first, the system acts directly on detector. Experiments Stern-Gerlach [11] and double slits [12] belong to this type. The second, the system acts on a middle system then it acts on the detector. Example the measurement of the radiative spectrum of the atom [13], the radiative field is the middle system.

The distribution function of possibility density of the system can be got from the first type. We must combine the measurement results and a suitable mathematical model to get other information about the system. We must combine the measurement results and solutions of the Schrodinger equation to get quantum properties of the system. Example, we can find the energy level's structure of an atom by combining the its experimental spectrum and perturbation solution of the Schrodinger equation. If the Schrodinger equation of the system is difficult to analyze then we can only drive a little information from the experiment results. Example, we only drive a little information from the visible radiative spectrum of solid materials [14] because its Schrodinger equation is very difficult to analyzing.

Qualitative quantum properties can be gotten from the qualitative solutions of the Schrodinger equation. Example, it may be driven to the atom's number of possible states from qualitative solutions of the Schrodinger equation.

Quantitative quantum properties can be gotten only from the combination between experiment result and quantitative solutions of the Schrodinger equation. The quantitative properties are divided into two types. The first are constants as mass and charge of micro-particle. The second, in a stationary system, are eigenvalues of the operator A which commutes with Hamiltonia. We can drive the value of the eigenvalues

 $a_i$  of operator A because they appear in the quantitative solutions of the Schrodinger equation. So Observable quantities in a particular experiment are indicated clearly by only Schrodinger equation without axioms of operators. They correspond to operators which their eigenvalues appear in the quantitative solutions of the Schrodinger equation. And we can define: quantitative quantum properties of the system can be described by operators, possible values of each quantity are eigenvalues of the operator. These operators need to be hermitic because their possible values are real.

# 3. The wave function collapse

If the potential energy is discontinuous then space is divided into many continuous domains  $D_i$ . Symmetry groups of the Schrodinger equation and possible states are different in each domain  $D_i$ . So the state of the system changes when it shifts to the next continuous domain. An observed system is acted by apparatus. We describe this effect by a measurement potential energy  $V_{meas}$ . It is made by the apparatus and like every other interaction. The total potential energy of system is  $V_{total} = V + V_{meas}$ . The total potential energy of system usually is discontinuous because of the measurement potential energy. So space is usually divided into continuous domains because of the measurement. The shifting between continuous domains of system makes the change of state. This process is suddenly and uncertain. And it is like the wave function collapse [5].

Here, the apparatus isn't restricted in the frame of classical law to understand the wave function collapse [15]. It isn't only the measurement potential energy, the state of the system but also may change suddenly because of any other interaction. The state of the system after measurement doesn't depend on the observer's mind. It only depends on the Hamiltonian and the measurement potential energy.

Here, there are no registered properties. This is difference from orthodox quantum mechanics. This difference may be cleared because the orthodox quantum mechanics can't indicate clearly observed operator in a particular experiment.

#### 1. Reality and the classical picture

From the Born rule, it is easy to find the average value of the operator A:  $\langle A \rangle = \langle \psi | A \psi \rangle = \int \psi^* A \psi . \, dx_\mu$ ,  $\mu = 0,1,2,3$ . Here, the integration is executed in 4-dimensional space. We define average in 3-dimensions space  $\langle A \rangle_t = \int \psi^* A \psi . \, dV$ . From the Schrodinger equation we have the Heisenberg equation for the 3-dimension average:

$$i. \, \hbar. \frac{d}{dt} \langle A \rangle_t = \frac{\partial \langle A \rangle_t}{\partial t} + \langle [H, A] \rangle_t$$

The Heisenberg equation for coordinate average is:

$$\frac{d}{dt}\langle x\rangle_t = \frac{\langle p_x\rangle_t}{m}$$

The Heisenberg equation for momentum average is:

$$\frac{d}{dt}\langle p_x \rangle_t = -\frac{\partial \langle V \rangle_t}{\partial x}$$

Two equations above are the motion equations in classical mechanics. It shows clearly a viewpoint that the classical properties are averages of the quantum properties. And the laws of the macro-world are obtained by averaging the quantum laws. It is

suitable if the classical picture is only considered averaging the uncertain complex waves. The classical picture is only a reflection image of quantum reality on human senses. In general, each sensing structure gets a correspondent reflection image of reality. However, the uncertain complex wave is the unique origin of all image. The uncertain complex waves are strange things because of restriction of human sensing. But no reason prevents these waves to become a real reality.

## III. Conclusion

Paper describes the quantum reality are uncertain complex waves. Although these waves are counterintuitive but they can offer the classical picture by averaging. An axiomatic system is suggested to describe these waves completely. They are obtained by modifying the axiomatic system of the orthodox quantum mechanics. Then the Schrodinger equation without axioms about operators can indicate the observable properties by itself. While we can understand the wave function collapse by using an axiom about possible states.

### References

- [1] P. Pickl, Interpretations of Quantum Mechanics, EPJ Web of Conferences 71, 00110 (2014)
- [2] A. Bassi et al, *Models of Wave-function Collapse, Underlying Theories, and Experiment Test*, arXiv: 1024.4325v3 [quant-ph] (2012)
- [3] L. Hardy, Nonlocality for two particles without inequalities for almost all entangled states, Phy.Rev.Lett **71** (1993) 1665-1668
  - [4] N. Gisin, Collapse. What else?, arXiv: 1701.08300v2 [quant-ph] (2017)
  - [5] S. Shrapnel et al, Updating the Born rule, New Journal of Physics, Vol 20, (2008)
  - [6] R. B. Griffiths, Consistent Quantum Measurements, arXiv: 1501.04813v2 [quant-ph] (2015)
  - [7] A. Hobson, Quantum realism is consistent with quantum facts, arXiv: 1810.08638 [quant-ph] (2018)
  - [8] D. K. Lazarou, Interpretation of Quantum Theory An Overview, arXiv: 0712.3466 [quant-ph] 2009
- [9] H. Zinkernagel et al, *Are we living in a quantum world? Bohr and quantum fundamentalism*, arXiv: 1603.00342 [quant-ph]
- [10] W. Greiner, *Quantum Mechanics: An Introduction*, pp 438-441,  $4^{th}$  ed (2001), Publisher: Springer, ISBN-13: 978-3540674580
- [11] W. Greiner, *Quantum Mechanics: An Introduction*, pp 468-471, 4<sup>th</sup> ed (2001), Publisher: Springer, ISBN-13: 978-3540674580
- [12] W. Greiner, *Quantum Mechanics: An Introduction*, pp 460-462, 4<sup>th</sup> ed (2001), Publisher: Springer, ISBN-13: 978-3540674580
- [13] F. Stienkemeier et al, Spectroscopy of alkali atoms (Li, Na, K) attached to large helium cluster, Z.Phys.D **38**, 253-263 (1996)
- [14] P. D. Cooper et al, A review of possible optical absorption features of oxygen molecules in the icy surfaces of outer solar system bodies, Planetary and Space Science **51** 183-192 (2003)

[15] M. Dorato, *Bohr's Realation Holism and the classical-quantum Interaction*, arXiv: 1608.00205 [quan-ph]