

Mesons As Helmholtzian Factorizations

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Abstract: Mesons may be written as simple Helmholtzian factorizations.

The Helmholtzian operator factorization is:

$$\mathbf{J} \equiv D_B D_A \mathbf{f} = ((\square - |m|^2)) \mathbf{f}$$

where:

$$D_B \equiv \begin{pmatrix} D_0 & D_3^{\Rightarrow} & -D_2^{\Rightarrow} & -D_1 \\ -D_3^{\Rightarrow} & D_0 & D_1^{\Rightarrow} & -D_2 \\ D_2^{\Rightarrow} & -D_1^{\Rightarrow} & D_0 & -D_3 \\ -D_1^{\Downarrow} & -D_2^{\Downarrow} & -D_3^{\Downarrow} & D_0^{\Downarrow} \end{pmatrix} \quad \& \quad D_A \equiv \begin{pmatrix} D_0^{\Downarrow} & -D_3^{\Rightarrow} & D_2^{\Rightarrow} & -D_1 \\ D_3^{\Rightarrow} & D_0^{\Downarrow} & -D_1^{\Rightarrow} & -D_2 \\ -D_2^{\Rightarrow} & D_1^{\Rightarrow} & D_0^{\Downarrow} & -D_3 \\ -D_1^{\Downarrow} & -D_2^{\Downarrow} & -D_3^{\Downarrow} & D_0 \end{pmatrix}$$

and:

$$D_i^+ \equiv (\partial_i + m_i), \quad D_i^- \equiv (\partial_i - m_i)$$

$$D_i \equiv \begin{pmatrix} D_i^+ & 0 \\ 0 & D_i^- \end{pmatrix}, \quad D_i^{\Downarrow} \equiv \begin{pmatrix} D_j^- & 0 \\ 0 & D_i^+ \end{pmatrix}, \quad D_i^{\Rightarrow} \equiv \begin{pmatrix} 0 & D_i^- \\ D_i^+ & 0 \end{pmatrix}, \quad D_i^{\Rightarrow\Downarrow} \equiv \begin{pmatrix} 0 & D_i^+ \\ D_j^- & 0 \end{pmatrix}$$

and:

$$\mathbf{f} \equiv \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}, \quad f^i \equiv \begin{pmatrix} f_+^i \\ f_-^i \end{pmatrix}$$

$$f_+ \equiv \begin{pmatrix} \begin{pmatrix} f_+^1 \\ 0 \end{pmatrix} \\ \begin{pmatrix} f_+^2 \\ 0 \end{pmatrix} \\ \begin{pmatrix} f_+^3 \\ 0 \end{pmatrix} \\ \begin{pmatrix} f_+^0 \\ 0 \end{pmatrix} \end{pmatrix}, \quad f_- \equiv \begin{pmatrix} \begin{pmatrix} 0 \\ f_-^1 \end{pmatrix} \\ \begin{pmatrix} 0 \\ f_-^2 \end{pmatrix} \\ \begin{pmatrix} 0 \\ f_-^3 \end{pmatrix} \\ \begin{pmatrix} 0 \\ f_-^0 \end{pmatrix} \end{pmatrix}, \quad f \equiv \begin{pmatrix} \begin{pmatrix} f_+^1 \\ f_-^1 \end{pmatrix} \\ \begin{pmatrix} f_+^2 \\ f_-^2 \end{pmatrix} \\ \begin{pmatrix} f_+^3 \\ f_-^3 \end{pmatrix} \\ \begin{pmatrix} f_+^0 \\ f_-^0 \end{pmatrix} \end{pmatrix} = f_+ + f_-$$

$$\Rightarrow \begin{pmatrix} -D_0 & D_3^{\Rightarrow} & -D_2^{\Rightarrow} & -D_1 \\ -D_3^{\Rightarrow} & -D_0 & D_1^{\Rightarrow} & -D_2 \\ D_2^{\Rightarrow} & -D_1^{\Rightarrow} & -D_0 & -D_3 \\ -D_1^{\Downarrow} & -D_2^{\Downarrow} & -D_3^{\Downarrow} & D_0^{\Downarrow} \end{pmatrix} \begin{pmatrix} -D_0^{\Downarrow} & -D_3^{\Rightarrow} & D_2^{\Rightarrow} & -D_1 \\ D_3^{\Rightarrow} & -D_0^{\Downarrow} & -D_1^{\Rightarrow} & -D_2 \\ -D_2^{\Rightarrow} & D_1^{\Rightarrow} & -D_0^{\Downarrow} & -D_3 \\ -D_1^{\Downarrow} & -D_2^{\Downarrow} & -D_3^{\Downarrow} & D_0 \end{pmatrix} \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}$$

$$= \begin{pmatrix} -D_0 & D_3^{\Rightarrow} & -D_2^{\Rightarrow} & -D_1 \\ -D_3^{\Rightarrow} & -D_0 & D_1^{\Rightarrow} & -D_2 \\ D_2^{\Rightarrow} & -D_1^{\Rightarrow} & -D_0 & -D_3 \\ -D_1^{\Downarrow} & -D_2^{\Downarrow} & -D_3^{\Downarrow} & D_0^{\Downarrow} \end{pmatrix} \begin{pmatrix} B_{\Downarrow}^1 + E^1 \\ B_{\Downarrow}^2 + E^2 \\ B_{\Downarrow}^3 + E^3 \\ -\nabla_{\Downarrow}^* \cdot \mathbf{f} \end{pmatrix}$$

Concerning the fundamental particles:

As shown in [2], a fundamental object is a vector components of which are mass-generalized electromagnetic field components, classified as follows:

$$e(i) \equiv \bar{\alpha}_i = \overline{(E^1, E^2, E^3)}_i, \quad v(i) \equiv \beta_i \equiv (B^1, B^2, B^3)_i$$

$$u_j(i) \equiv \phi_{ji} = (\eta_{j-1}(E^1), \eta_{j-2}(E^2), \eta_{j-3}(E^3))_i, \quad d_j(i) \equiv \bar{\psi}_{ji} = \overline{(\eta_{j-1}(B^1), \eta_{j-2}(B^1), \eta_{j-3}(B^1))}_i$$

where:

$$\eta_j(R_k^h) \equiv \begin{cases} R_k^h, & j \neq 0 \\ E_k^h, & j = 0, \quad \mathbf{R} = \mathbf{B} \\ B_k^h, & j = 0, \quad \mathbf{R} = \mathbf{E} \end{cases}, \quad \sigma_j(\mathbf{R}_k) \equiv \begin{pmatrix} \eta_{j-1}(R_k^1) \\ \eta_{j-2}(R_k^2) \\ \eta_{j-3}(R_k^3) \end{pmatrix}$$

(i denoting generation, j denoting color)

where:

$$\mathbf{E} = ((-D_0^{\Downarrow} f^1 - D_1 f^0), (-D_0^{\Downarrow} f^2 - D_2 f^0), (-D_0^{\Downarrow} f^3 - D_3 f^0), *)$$

$$\mathbf{B} = ((D_2 f^3 - D_3 f^2), (-D_1 f^3 + D_3 f^1), (D_1 f^2 - D_2 f^1), *)$$

$$\mathbf{E}_{\Downarrow} = ((-D_0^{\Rightarrow\Downarrow} f^1 - D_1^{\Rightarrow} f^0), (-D_0^{\Rightarrow\Downarrow} f^2 - D_2^{\Rightarrow} f^0), (-D_0^{\Rightarrow\Downarrow} f^3 - D_3^{\Rightarrow} f^0), *)$$

$$\mathbf{B}_{\Downarrow} = ((D_2^{\Rightarrow} f^3 - D_3^{\Rightarrow} f^2), (-D_1^{\Rightarrow} f^3 + D_3^{\Rightarrow} f^1), (D_1^{\Rightarrow} f^2 - D_2^{\Rightarrow} f^1), *)$$

(where * denotes a gauge component)

So, in particular (written horizontally as vectors for brevity):

$e^- = e(1) = \overline{(E^1, E^2, E^3)}_1$	$\mu^- = e(2) = \overline{(E^1, E^2, E^3)}_2$	$\tau^- = e(3) = \overline{(E^1, E^2, E^3)}_3$
$v_e = v(1) = (B^1, B^2, B^3)_1$	$v_\mu = v(2) = (B^1, B^2, B^3)_2$	$v_\tau = v(3) = (B^1, B^2, B^3)_3$
$u_R = u_1(1) = (B^1, E^2, E^3)_1$	$c_R = u_1(2) = (B^1, E^2, E^3)_2$	$t_R = u_1(3) = (B^1, E^2, E^3)_3$
$u_G = u_2(1) = (E^1, B^2, E^3)_1$	$c_G = u_2(2) = (E^1, B^2, E^3)_2$	$t_G = u_2(3) = (E^1, B^2, E^3)_3$
$u_B = u_3(1) = (E^1, E^2, B^3)_1$	$c_B = u_3(2) = (E^1, E^2, B^3)_2$	$t_B = u_3(3) = (E^1, E^2, B^3)_3$
$d_R = d_1(1) = \overline{(E^1, B^2, B^3)}_1$	$s_R = d_1(2) = \overline{(E^1, B^2, B^3)}_2$	$b_R = d_1(3) = \overline{(E^1, B^2, B^3)}_3$
$d_G = d_2(1) = \overline{(B^1, E^2, B^3)}_1$	$s_G = d_2(2) = \overline{(B^1, E^2, B^3)}_2$	$b_G = d_2(3) = \overline{(B^1, E^2, B^3)}_3$
$d_B = d_3(1) = \overline{(B^1, B^2, E^3)}_1$	$s_B = d_3(2) = \overline{(B^1, B^2, E^3)}_2$	$b_B = d_3(3) = \overline{(B^1, B^2, E^3)}_3$

Mesons may be written as a simple Helmholtzian factorization.

For example:

$$\begin{aligned}
 u_R &= u_1(1) = (B^1, E^2, E^3)_1 = ; \quad d_R = d_1(1) = \overline{(E^1, B^2, B^3)}_2 \\
 u_R &= u_1(1) = \left(\begin{pmatrix} B_{\hat{\downarrow}}^1 \\ E^2 \\ E^3 \\ * \end{pmatrix}, \mathbf{0}, \mathbf{0} \right) ; \quad \overline{d_R} = \overline{d_1(1)} = \left(\begin{pmatrix} E^1 \\ B_{\hat{\downarrow}}^2 \\ B_{\hat{\downarrow}}^3 \\ * \end{pmatrix}, \mathbf{0}, \mathbf{0} \right) \\
 \Rightarrow u_R : \overline{d_R} &= \left(\begin{pmatrix} B_{\hat{\downarrow}}^1 + E^1 \\ B_{\hat{\downarrow}}^2 + E^2 \\ B_{\hat{\downarrow}}^3 + E^3 \\ * \end{pmatrix}, \mathbf{0}, \mathbf{0} \right) = \left(\begin{pmatrix} (B_{\hat{\downarrow}} + E)^1 \\ (B_{\hat{\downarrow}} + E)^2 \\ (B_{\hat{\downarrow}} + E)^3 \\ * \end{pmatrix}, \mathbf{0}, \mathbf{0} \right) = (B_{\hat{\downarrow}} + E)_1 = \pi^+ . \\
 &= \left(\begin{pmatrix} -D_0 & -D_3^{\leftrightarrow} & D_2^{\leftrightarrow} & -D_1 \\ D_3^{\leftrightarrow} & -D_0 & -D_1^{\leftrightarrow} & -D_2 \\ -D_2^{\leftrightarrow} & D_1^{\leftrightarrow} & -D_0 & -D_3 \\ -D_1^{\hat{\downarrow}} & -D_2^{\hat{\downarrow}} & -D_3^{\hat{\downarrow}} & D_0 \end{pmatrix} \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}, \mathbf{0}, \mathbf{0} \right) = D_A(\mathbf{f}, \mathbf{0}, \mathbf{0}) \\
 \Rightarrow J(u_R : \overline{d_R}) &= ((\square - |m|^2))(\mathbf{f}, \mathbf{0}, \mathbf{0}) = D_B D_A(\mathbf{f}, \mathbf{0}, \mathbf{0}) = D_B((B_{\hat{\downarrow}} + E)_1) = D_B(\pi^+)
 \end{aligned}$$

$$\begin{aligned}
 d_G &= d_2(1) = \overline{(B_{\hat{\downarrow}}^1, E^2, B_{\hat{\downarrow}}^3)}_1 ; \quad c_G = u_2(2) = (E^1, B_{\hat{\downarrow}}^2, E^3)_2 \\
 \Rightarrow d_G : \overline{c_G} &= \left(\begin{pmatrix} \overline{B_{\hat{\downarrow}}^1} \\ \overline{E^2} \\ \overline{B_{\hat{\downarrow}}^3} \\ * \end{pmatrix}, \begin{pmatrix} \overline{E^1} \\ \overline{B_{\hat{\downarrow}}^2} \\ \overline{E^3} \\ * \end{pmatrix}, \mathbf{0} \right) = \left(\begin{pmatrix} (\overline{B_{\hat{\downarrow}} + E})^1 \\ (\overline{B_{\hat{\downarrow}} + E})^2 \\ (\overline{B_{\hat{\downarrow}} + E})^3 \\ * \end{pmatrix}, \mathbf{0}, \mathbf{0} \right) = \overline{(B_{\hat{\downarrow}} + E)_{1/2}}
 \end{aligned}$$

$$\Rightarrow J(d_G : \overline{c_G}) = ((\square - |m|^2))(\overline{\mathbf{f}}, \mathbf{0}, \mathbf{0}) = D_B D_A(\overline{\mathbf{f}}, \mathbf{0}, \mathbf{0}) = D_B((\overline{B_{\hat{\downarrow}} + E})_{1/2})$$

$$\begin{aligned}
 c_R &= u_1(2) = (B_{\hat{\downarrow}}^1, E^2, E^3)_2 ; \quad d_R = d_1(1) = \overline{(E^1, B_{\hat{\downarrow}}^2, B_{\hat{\downarrow}}^3)}_1 \\
 \Rightarrow c_R : \overline{d_R} &= \left(\begin{pmatrix} E^1 \\ B_{\hat{\downarrow}}^2 \\ B_{\hat{\downarrow}}^3 \\ * \end{pmatrix}, \begin{pmatrix} B_{\hat{\downarrow}}^1 \\ E^2 \\ E_{\hat{\downarrow}}^3 \\ * \end{pmatrix}, \mathbf{0} \right) = \left(\mathbf{0}, \begin{pmatrix} (B_{\hat{\downarrow}} + E)^1 \\ (B_{\hat{\downarrow}} + E)^2 \\ (B_{\hat{\downarrow}} + E)^3 \\ * \end{pmatrix}, \mathbf{0} \right) = (B_{\hat{\downarrow}} + E)_2 = D^+
 \end{aligned}$$

$$= \left(\mathbf{0}, \begin{pmatrix} -D_0 & -D_3^{\leftrightarrow} & D_2^{\leftrightarrow} & -D_1 \\ D_3^{\leftrightarrow} & -D_0 & -D_1^{\leftrightarrow} & -D_2 \\ -D_2^{\leftrightarrow} & D_1^{\leftrightarrow} & -D_0 & -D_3 \\ -D_1^{\hat{\downarrow}} & -D_2^{\hat{\downarrow}} & -D_3^{\hat{\downarrow}} & D_0 \end{pmatrix} \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}, \mathbf{0} \right) = D_A(\mathbf{0}, \mathbf{f}, \mathbf{0})$$

$$\Rightarrow J(c_R : \overline{d_R}) = ((\square - |m|^2))(\mathbf{0}, \mathbf{f}, \mathbf{0}) = D_B D_A(\mathbf{0}, \mathbf{f}, \mathbf{0}) = D_B((B_{\hat{\downarrow}} + E)_2) = D_B(D^+)$$

So, clearly, mesons of the same generation may be written simple Helmholtzian factorizations.
(mesons of differing generations are something like 'split-level' Helmholtzian factorizations).

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