SunQM-4: Using full-QM deduction and \{N,n\} QM’s non-Born probability density 3D map to build a complete Solar system with orbital movement

Yi Cao
Ph.D. in Biophysics
e-mail: yicaojob@yahoo.com
© All rights reserved. Submitted to viXra.org on 3/25/2020.

Abstract

In SunQM-3 series, we studied Solar \{N,n\} QM within the boundary of the traditional Schrodinger equation/solution and Born’s rule. In SunQM-4 series, we start to relax that boundary. In the current paper, for a planet in nLL QM state doing circular orbital movement, we deduced out (a full-QM deduced) |Φ(φ)|^2 * |T(t)|^2 for a planet’s time-dependent probability density in φ-dimension. To satisfy the well-known QM rule that a matter wave’s group velocity equals to 2x of its phase velocity, we have to define a non-Born calculation as |T(t)|^2 ∝ (exp(-i * ω_{n,ph} * t))^2 where ω_{n,ph} is the phase angular frequency of a planet’s matter wave in φ-dimension. To obtain a physical meaningful |Φ(φ)|^2 * |T(t)|^2, we have to define a non-Born probability density calculation as |Φ(φ)|^2 ∝ Φ(φ), or its φ-dimensional probability density function is directly proportional to its matter wave function. Combining with SunQM-3s11’s result, we built a complete Solar system with time-dependent circular orbital movement using the full-QM deduced non-Born probability density 3D map. This 3D probability density described a Solar system not only at planet’s Eigen description level, but also at any level of resolution (down to proton level, or up to the whole universe level). Therefore, we propose that “Simultaneous-Multi-Eigen-Description (SMED)” is one of many nature attributes of QM. We believe that by adding the non-Born calculation to Born calculation, the QM will become more self-consistent and more complete.

Introduction

The SunQM series articles [1] - [16] have shown that the formation of Solar system (as well as each planet) was governed by its \{N,n\} QM. In SunQM-3 series, we studied Solar \{N,n\} QM by using the traditional Schrodinger equation/solution and Born’s probability. For example, in papers SunQM-3s6, -3s7, and -3s8, it has been shown that the formation of planet’s and star’s (radial) internal structure is governed by the planet’s or star’s radial QM. In papers SunQM-3s3 and -3s9, it has been shown that the surface mass (atmosphere) movement of Sun, Jupiter, Saturn, and Earth, etc., is governed by Star’s (or planet’s) θφ-2D dimension QM. In paper SunQM-3s4 and -3s10, it has been shown that the formation of either ring structures of a planet, or the belt structures in Solar system, is also governed by the \{N,n\} QM (the nLL effect). In paper SunQM-3s11, we have used \{N,n\} QM and Schrodinger equation’s solution to build a 3D probability density map for a complete Solar system with time-dependent orbital movement. In SunQM-4 series, we are going to study Solar \{N,n\} QM by using the method that more or less deviated from the traditional Schrodinger equation and Born’s rule. In the current paper, after using a non-Born probability calculation (where the probability is directly proportional to the wave function, not the conjugated-squared wave function), we have built a non-Born probability density 3D map for a complete Solar system with time-dependent orbital movement. Note: for \{N,n\} QM nomenclature as well as the general notes for \{N,n\} QM model, please see SunQM-1 section VII. Note: Microsoft Excel’s number format is often used in this paper, for example: x^2 = x^2, 3.4E+12 = 3.4*10^{12}, 5.6E-9 = 5.6*10^{-9}. Note: The reading sequence for SunQM series papers is: SunQM-1, 1s1, 1s2, 1s3, 2, 3, 3s1, 3s2, 3s6, 3s7, 3s8, 3s3, 3s9, 3s4, 3s10, 3s11, 4. Note: for all SunQM series papers, reader should check “SunQM-4s7: Updates and Q/A for SunQM series papers” for the most recent updates and corrections.
I. To build a time-dependent 3D probability density for a planet (in Solar system’s orbit) based on non-Born probability $|\Phi(\phi)|^2 \propto \exp(\text{i}m\phi)$

In the SunQM-3 series papers, we have demonstrated that our Solar system can be described by Schrodinger equation

$$i\hbar \frac{\partial}{\partial t} \Psi(r, \theta, \phi, t) = \left[ -\frac{\hbar^2}{2m} \nabla^2 + V(r, \theta, \phi, t) \right] \Psi(r, \theta, \phi, t) \quad \text{eq-1}$$

Under certain physics condition (e.g., plane wave, or hydrogen atom, etc.), Schrodinger equation (as a linear partial differential equation) can be solved by separating the variables so that we can find solutions that are simple products of

$$\Psi(r, \theta, \phi, t) = R(r) \Theta(\theta) \Phi(\phi) T(t) \quad \text{eq-2}$$

Because of the rotation diffusion (or RotaFusion, or RF, see SunQM-2 for details) in $\theta\phi$-2D-dimension, the two functions of $\Psi$ in $\theta$- and $\phi$-dimensions are usually grouped together as spherical harmonics

$$\Theta(\theta) \Phi(\phi) = Y(l, m) \quad \text{eq-3}$$

Although the traditional time-independent probability density formula for Schrodinger equation is written as

$$r^2 |\Psi(r, \theta, \phi, t)|^2 = r^2 |R(r)|^2 |Y(l, m)|^2 \quad \text{eq-4}$$

where $R(r) = R(n)$, the general time-dependent probability density formula for Schrodinger equation should be something like

$$r^2 |\Psi(r, \theta, \phi, t)|^2 = r^2 |R(r)|^2 |\Theta(\theta)|^2 |\Phi(\phi)|^2 |T(t)|^2 \quad \text{eq-5}$$

In SunQM-3s11 sections II-a, II-b and II-c, planet’s probability density in $r$-dimension $r^2 \* |R(r)|^2$, in $\theta$-dimension $|\Theta(\theta)|^2$, and in $\phi$-dimension $|\Phi(\phi)|^2$ were all deduced by using full-QM. However, for the time-dependency (of $\phi$-dimensional orbit movement), we simply replaced $\phi$ by $\phi$-ot. This makes the $\phi$-dimension’s time-dependent probability $|\Phi(\phi)|^2 \* |T(t)|^2$ become a semi-QM deduced function. So in the current paper, we try to deduce $|\Phi(\phi)|^2$ and $|T(t)|^2$ separately, and then combine them through a simple production as shown in eq-5.

I-a. Defining $|\Phi(\phi)|^2 \propto \exp(\text{i}m\phi)$, or $\exp(-\text{i}m\phi)=1$, for a planet’s $\phi$-dimensional probability density in Solar system

In the traditional QM, for the nLL QM state (where $l=n-1$, $m=n-1$, see SunQM-3s1), by default, the $\phi$-dimension’s Born probability density

$$|\Phi(\phi)|^2 \propto |e^{\text{i}m\phi}|^2 = e^{-\text{i}m\phi}e^{\text{i}m\phi} = 1 \quad \text{eq-6}$$

It can be explained as its $+\phi$ directional wave cancels out its own $-\phi$ directional wave. Now we try to use the general physics to explain why this is not correct for a planet’s orbital movement in the $\{N,n\}$ QM (because all QM has to degenerated back to classical physics when the quantum number $n \to \infty$). For the Solar $\{N,n\}$ QM, let’s define that $\exp(\text{i}m\phi)$ correlates to the (planet mass’ or matter wave’s) eastward circular orbital movement, and $\exp(-\text{i}m\phi)$ correlates to the (planet mass’ or matter wave’s) westward circular orbital movement. From the result of SunQM-3s8 Table 6, we understand that Solar $\{N,n\}$ QM’s mass movement (or matter wave movement) is directly correlated to the mass’ orbital velocity (notice that matter wave is a wave packet, so its orbital velocity is group velocity). Inside the Sun (where most atoms are not in nLL QM state), this orbital velocity mass movement is transformed into the micro thermal movement, and it is always in RF (notice that Sun’s self-spin
movement is ignored here, because Sun’s spin velocity is too slow in comparison to Sun atoms’ thermal movement velocity). Outside the Sun in the planet region, the same orbital velocity mass movement is transformed into planet’s (macro) orbital movement and this orbital movement is neither the micro thermal movement, nor in RF.

For the Solar \([N,n]\) QM, the traditional QM probability calculation is still correct for the space region with \(-100\%\) mass occupancy (while the slow macro movement can be ignored). Because there (inside Sun, ignoring the Sun self-spin, only consider the micro thermal movement) all mass is doing RF, so (in \(\varphi\)-dimension) at any time, there is always equal amount of mass (or matter waves) doing the eastward \(\exp(\im\varphi)\) thermal movement as that doing the westward \(\exp(-\im\varphi)\) thermal movement, so that the \(\varphi\)-dimension’s matter wave related Born probability \(\left|\Phi(\varphi)\right|^2\) has to combine these two opposite direction’s matter waves in form of eq-6, and it becomes a delocalized standing wave and \(\varphi\)-dimensional evenly distributed probability. However, for the space region with \(<1\%\) mass occupancy in the solar system (mainly the planet region, now we have to use planet’s orbital velocity), all mass (or matter waves) has only the eastward \(\exp(+\im\varphi)\) macro movement, no westward \(\exp(-\im\varphi)\) macro movement, then what is the value of \(\exp(-\im\varphi)\)? After many tries, we were forced to choose

\[
e^{-\im\varphi} = 1
eq-7
\]

The physical meaning of eq-7 is explained in section I-h’s discussion item 1). Thus for a planet in a \(|nlm\rangle = |nLL\rangle = |n,l=n-1,m=n-1\rangle\) QM state, we have the \(\varphi\)-dimensional wave function

\[
\Phi(\varphi) \propto e^{\im\varphi}
eq-8
\]

and the \(\varphi\)-dimensional probability density

\[
\left|\Phi(\varphi)\right|^2 \propto \left|e^{\im\varphi}\right|^2 = e^{-\im\varphi}e^{\im\varphi} = 1 \times e^{\im\varphi} = e^{\im\varphi}
\]

So \(\exp(-\im\varphi) = 1\) can also be explained as that for an eastward orbital moving planet’s matter wave, its westward wave component is zero, or its westward wave component makes no contribution to the probability density.

Then how to alter a complex numbered probability in eq-9 into a real numbered probability? Here let’s first use a citizen-scientist level method: for the QM probability of \([\exp(ix)]^2 = [\cos(x) + \im\sin(x)]^2\), we simply ignore the imaginary part, only use the real part, so that \([\exp(ix)]^2 = [\cos(x) + \im\sin(x)]^2 \rightarrow [\cos(x)]^2\). Using this method, we have

\[
\left|\Phi(\varphi)\right|^2 \propto e^{\im\varphi} = \left(e^{\im\varphi / 2}\right)^2m = \left\{\left[(\cos(\varphi / 2) + \im\sin(\varphi / 2))\right]^m \rightarrow \left\{[\cos(\varphi / 2)]^2\right\}^m = \left[1 + \cos(\varphi / 2)\right]^m = \left[\frac{1 + \cos(\varphi)}{2}\right]^m = \left[\frac{1 + \cos(\varphi)}{2}\right]^{(n-1)}
\]

where \(m = +l = n-1\). The 2nd last step in eq-10 is from the general trigonometric relationship:

\[
[\cos(\varphi / 2)]^2 = \frac{1 + \cos(\varphi)}{2}
eq-11
\]

In this way, we altered a complex numbered probability into a real numbered probability. The power index \((n-1)\) in eq-10 has the function to make \([1 + \cos(\varphi)] / 2\) curve peak become narrower. Therefore eq-10 equivalent to a true wave packet description for a planet in the \(\varphi\)-1D-dimension.

Now let’s use a formal mathematical method to explain eq-10. First, let’s define an equivalent wave function

\[
\Phi(\varphi)_{\text{equivalent}} \propto e^{\im\varphi / 2}
eq-12
\]

so that
Notice that eq-12 and eq-13 do have the real physical meaning (see SunQM-4s1 section I-b). Second, let’s use eq-31-SunQM-3s11 (meaning SunQM-3s11’s eq-31) to construct a new real numbered wave function, (Notice that the m (or l) in eq-31-SunQM-3s11 is replaced by m/2 (or l/2) here)

\[ Y\left(\frac{l}{2} = n - 1, \pm \frac{m}{2} = n - 1\right) = \cos(m\phi/2) \left[\sin(\theta)\right]^{1/2} \]

Where \(l/2 = m/2 = n-1\), and its probability function

\[ \left|Y\left(\frac{l}{2} = n - 1, \pm \frac{m}{2} = n - 1\right)\right|^2 = \left[\cos(m\phi/2)\right]^2 \left[\sin(\theta)\right]^m \]

Third, let’s use a linear combination of eq-15 and follow the method from eq-33-SunQM-3s11 through eq-43-SunQM-3s11 (notice that the \(m \rightarrow m/2\) is replaced by \(\phi \rightarrow \phi/2\) here, and we can NOT do both at the same time), we obtain

\[ \left\{ \frac{1}{1 + 2\delta} \sum_{\delta} \cos((m + \delta)\phi/2) \right\}^{1/2} \rightarrow \cos(\phi/2)^{2m} \]

or, the left side of eq-16 is infinitely approaching to the right side of eq-16 when the integer \(\delta \ll m\) value is high enough. 

Eq-16 bridges the complex numbered probability of \([\exp(i \phi)]^2 = [\cos(x) + i\sin(x)]^2\) to its corresponding real numbered probability \([\cos(x/2)]^2\) in eq-10. For better understanding, please read SunQM-3s11 section III-c, and SunQM-4s1.

**I-b. How a planet’s \(n\) in \([N,n/6]\) QM relates to the \(n\) in Bohr-de Broglie’s 1D circular orbit QM in the formula \(n\lambda = 2\pi r\)?**

This is a pre-required knowledge for section I-c’s study. In the Bohr-de Broglie’s 1D circular orbit QM (abbreviated as “circular 1D QM”)’s formula \(n\lambda = 2\pi r\), let’s rename the quantum number \(n\) to be \(j\), so it becomes

\[ j \lambda = 2\pi r \]

where \(j = 1, 2, \ldots\) is a positive integer number (at least for now). Using Earth-corelated \(\{1.5/6\}\) orbit as the example, it has orbit \(r = 1.565E+11\) m. So using circular 1D QM’s eq-17, at \(j=1, \lambda = 2\pi/j = 9.84E+11\) m. At \(j=5, \lambda = 2\pi/5 = 1.97E+11\) m. At \(j=30, \lambda = 2\pi/30 = 3.28E+10\) m. Now comparing to \([N,n/q]\) QM’s orbit \(\lambda\) calculation (in SunQM-2 Table 1), we see that \(\{1.5/6\}\) with \(r_1 = \{1,1/6\}\) and \(n=5\) has \(\lambda = 1.97E+11\) m, exactly equals to circular 1D QM’s \(j=5\). Also, we see that \(\{0.3/6\}\) with \(r_1 = \{0,1/6\}\) and \(n = 5^6^1 = 30\) has \(\lambda = 3.28E+10\) m, exactly equals to circular 1D QM’s \(j=30\). It is obvious that at least for \(nLL\) QM state (therefore for all planets or even for all belts in Solar system), \([N,n/q]\) QM’s quantum number \(n\) is a special set of quantum number \(j\) in circular 1D QM which satisfy

\[ j = n q^{(1-N)} \]

if we set \(r_1\) at \([N,1/6]\). Notice that (in eq-18) \(j\) is \(\phi\)-1D-dimension-only quantum number, while \(n\) is \(r\phi\) 3D-dimension quantum number (see SunQM-4s1 section I-a for more detailed figure explanation). For example: if we set \(r_1\) at \([N=1,1/6]\),
then a \{1.5//6\} orbit has circular 1D QM’s \(j = 5^*6^*(1-1) = 5\). If we set \(r_1\) at \(\{N=0,1//6\}\), then a \{1.5//6\} = \{0,30//6\} orbit has circular 1D QM’s \(j = 5^*6^*(1-0) = 30\). If we set \(r_1\) at \(\{N=2,1//6\}\), then a \{1.5//6\} = \{2.5//6\} orbit has circular 1D QM’s \(j = 5^*6^*(1-2) = 5/6\), (notice that \(j=5/6\) is not a integer, but a pseudo-integer number for \(\{N,n/q\}\) QM, because in \(\{N,n/q\}\) QM, all \(n/q\) numbers are treated as pseudo-integer number). Hence, we can move \(r_1\) around to change the circular 1D QM’s \(j\). In this way, we also unified the \(\{N,n/q\}\) QM with the Bohr-de Broglie’s 1D circular orbital movement QM (at least for the nLL QM state, like all planets in the Solar system).

**Statement-1**: In circular 1D QM, the ground state is \(j = 1\), so it correlates to \(r_1\) at \(\{N,n/q\}\) orbit in \(\{N,n/q\}\) QM. For example, circular 1D QM’s \(j = 1\) ground state for \(\{1,5//6\}\) orbit is \(r_1\) at \(\{1,5//6\}\). This knowledge is crucial in the next section’s full-QM deduction.

I-c. For a circular orbital moving planet, deduce its relationship between \(E_{\phi,j}/h_{\text{gen}}, r_n, v_n, v_{n,\text{ph}}, v_{n,gr}, \text{and} \ 0n,gr, \text{and} \ 0n,\text{ph} \)

Here we use (de Broglie’s) free particle plane wave (which is forced to do the circular orbital moving in \(\phi, 1D\) dimension) to describe a planet’s orbital movement (see SunQM-3s11 section III-c, and SunQM-4s1 for why this is valid). The standard wave function of a plane wave (as Schrodinger equation’s solution) is \cite{17}

\[
\Psi_k(x,t) = A \exp\left[i(kx - \frac{\hbar}{2m}k^2t)\right] = A \exp\left[i(kx - \frac{E}{\hbar})t\right] = A \exp[i(px - \text{Et})/\hbar]
\]

for \(\text{eq-19}\)

where \(k\) is the wave number, \(h\) is the Planck constant divided by \(2\pi\), \(m^*\) is the mass of a particle/planet (not a quantum number), and the momentum \(p = h * k\). Here we only interest in the time-dependency (of a planet’s orbital movement) in \(\phi, 1D\)-dimension, so \(E\) is the energy only in \(\phi, 1D\)-dimension. Therefore, we need to use quantum number \(n\) (rather than quantum number \(j\)). It is

\[
E_{\phi,j} = K_{\phi,j} + V_{\phi,j} = K_{\phi,j} + \frac{1}{2} m^* v_j^2
\]

for \(\text{eq-20}\)

because \(V_{\phi,j} = 0\) in \(\phi, 1D\)-dimension (for a point-central gravity-force field). Notice that it should not be

\[
E_n = E_{r\theta\phi,n} = -\frac{1}{2} m^* v_n^2
\]

for \(\text{eq-21}\)

because \(\text{eq-21}\) includes the \(E_{r\theta\phi}\)’s contribution (see SunQM-4s4 section VII for more detailed explanation). Then according to \(\text{eq-19}\), the time portion (or the time dependency) of a plane wave’s wave function is

\[
T(t) \propto e^{-i\frac{E_{\phi,j}}{\hbar}}
\]

for \(\text{eq-22}\)

According to SunQM-2’s result, we should use \(h_{\text{gen}}\) (the general Planck constant divided by \(2\pi\)) to replace \(h\) in the Solar \(\{N,n/q\}\) QM. So \(\text{eq-22}\) become

\[
T(t) \propto e^{-i\frac{E_{\phi,j}}{h_{\text{gen}}}}
\]

for \(\text{eq-23}\)

Also according to SunQM-2’s result, the wave form of \(\text{eq-20}\) is

\[
E_{\phi,j} = Hm^* f_j = h_{\text{gen}} f_j
\]

for \(\text{eq-24}\)

Where \(H\) is the pseudo-Planck constant, \(f_j\) is the planet matter wave’s circular orbital moving frequency, and \(m^*\) is the mass of the planet. Again notice that \(E_n = E_{n,\phi\theta} = -H * m^* \ f_n\) is \(E_n\) in \(r\phi\theta\) 3D-dimension, and \(E_{\phi,j} = H * m^* \ f_j\) is \(E_j\) in \(\phi, 1D\)-
dimension. Also notices that values of $h_{gen}$, $H$, $f_n$ depend on where you choose $r_1$, while values of $E_n$, $r_n$, $v_n$ are $r_1$ position independent value (see SunQM-2 for explanation). Statement-2: Use $\varphi$-1D-dimension’s $E_\varphi$, with $V_\varphi = 0$, then $E_\varphi = K_\varphi$. So the $\omega_{n,ph}$ is only for $\varphi$-1D-dimension.

So, a planet matter wave’s (not the planet’s) time portion is

$$T(t) \propto e^{-\frac{E_{eq,j}}{h_{gen}}} = e^{-\frac{h_{gen}f_j}{h_{gen}}} = e^{-i2nf_jt} \tag{eq-25}$$

where

$$2\pi f_j = \omega_{j,ph} \tag{eq-26}$$

where $\omega_{j,ph}$ is the phase angular frequency/velocity. Notice that eq-26 is the planet matter wave’s angular frequency/velocity (because its meaning is more like angular velocity, although in Giancoli’s text book it is called angular frequency), it does not equal to the planet’s circular orbit moving angular frequency/velocity. Later on we will see that the former one is related to matter wave packet’s phase velocity, and the later one is related to matter wave packet’s group velocity (for better understanding, see SunQM-3s11’s Figure 1, and SunQM-4s1). Also notice that for a circular moving planet matter wave’s $f_j$ (or $\omega_{n,ph}$), the quantum number $j$ (in $\varphi$-1D-dimension) is inter-changeable with the quantum number $n$ (in 3D-dimension) through eq-18. It is much convenient to use $n$ rather than $j$ because we can directly compare with the previous results (e.g., in SunQM-2 Table 1) for $f_n$, $\omega_{n,ph}$, (except for $E$ because eq-20 does not equal to eq-21). For example, after re-write eq-26 as

$$2\pi f_n = \omega_{n,ph} \tag{eq-27}$$

then we can directly use SunQM-2 Table 1 column 15 to show that Earth’s matter wave orbit frequency $f_n$ at $n=30$ (or $r_1$ at \{0,1\}) is 4.44E-7 (cycle/sec = Hz), so its $2\pi \times f_n = 2.79E-6$ arc/s. It does not equal to Earth planet orbit’s $\omega = \omega_{eqr} = 2\pi / (365.2$ days) = 1.99E-7 arc/s. Similarly, after switch $j$ for $n$, eq-25 can be re-written as

$$T(t) \propto e^{-\frac{E_{eq,j}}{h_{gen}}} = e^{-i2nf_n t} \tag{eq-28}$$

Now let’s use particle QM’s method to deduce the relationship between $E_{n,j}/h_{gen}$, $r_n$, $v_{n,gr}$, $v_{n,ph}$, $\omega_{n,gr}$ and $\omega_{n,ph}$. From SunQM-2 Table 1, $H = h/m' = 2\pi * \sqrt{G*M*r_1}$, $n = \sqrt{r_n/r_1}$, combining two we obtain $H*n = 2\pi *\sqrt{G*M*r_n}$ for the G-central-forced circular movement. So $H$ is a $n$ related value, and should be written as $H_n$

$$H_n = \frac{2\pi}{n} \sqrt{GMr_n} \tag{eq-29}$$

Then, using Newton’s $F = m'*a = m' * v_n^2 / r_n$, $F = G*M*m'/r_n^2$, $m'^*v'^2/r_n = G*M*m'/r_n^2$, $r_n * v_n^2 = G*M$, $r_n * v_n = \sqrt{G*M*r_n}$, we have

$$\frac{E_{eq,j}}{h_{gen}} = \frac{2m'm'v_n^2}{2\pi} = \frac{nv_n^2}{h_n} = \frac{nv_n^2}{2\pi \sqrt{GMr_n}} = \frac{n}{2} \frac{v_n^2}{r_n} = \left(\frac{n}{2}\right) \frac{v_n}{r_n} \tag{eq-30}$$

Combining eq-28 and eq-30, we have

$$2\pi f_n = \left(\frac{n}{2}\right) \frac{v_n}{r_n} = \omega_{n,ph} \tag{eq-31}$$

Here we define $\omega_{n,ph}$ as the planet/particle’s matter wave’s orbital angular (phase) frequency/velocity, and $f_n$ is the planet/particle’s matter wave’s orbital frequency. Then, eq-19 can be written as
\[
\Psi_k(x,t) = Ae^{i(kx - \frac{E}{\hbar} t)} = Ae^{i(kx) e^{-i\omega_{n,ph} t}} = Ae^{i(kx) e^{-i\omega_{n,ph} t}}
\]

where

\[T(t) \propto e^{-i\omega_{n,ph} t}\]

Now let’s define

\[\frac{v_n}{r_n} = \omega_n = \omega_{n,gr}\]

Since \(r_n\) and \(v_n\) is planet/particle’s orbit \(r\) and orbit \(v\), so we know that \(\omega_n = v_n/r_n\) must equal to planet/particle’s orbital (group) angular frequency/velocity. Notices that in eq-31 and eq-34, values of \(n, \omega_{n,ph}\), and \(f_n\) depend on where we choose \(r_1\), while values of \(r_n, v_n, \omega_n = \omega_{n,gr}\) are \(r_1\) position independent values.

QM text books tell us that a particle’s matter wave is actually a wave packet that has group velocity (\(v_{gr}\)) and phase velocity (\(v_{ph}\)). Our results in SunQM-3s11 section III-c also revealed that a planet’s wave function in the \(\phi\)-dimension is composed by a group of wave functions (which are the Schrodinger equation’s solution) that further forms a wave packet (or group wave) out of the phase waves (see SunQM-3s11 Figure 1, and SunQM-4s1). According to text books [18, 19], the relationship between the classical particle velocity and matter wave \(v_{gr}\) and \(v_{ph}\) is:

\[v_{classical} = v_{gr} = 2v_{ph}\]

Notice that eq-35 equivalent to

\[v_{classical} = v_n = v_{n,gr} = v_{gr} = 2v_{ph} = 2v_{n,ph}\]

For any wave (including matter wave), it always has

\[v_{ph} = \lambda f\]

where \(\lambda\) is the wavelength and \(f\) is the wave frequency. Using eq-36 or \(v_{n,ph} = v_n/2\), we can re-write eq-31 as

\[\omega_{n,ph} = 2\pi f_n = \left(\frac{n}{2}\right) \frac{v_n}{r_n} = \frac{n v_{n,ph}}{r_n}\]

where \(v_{n,ph}\) is the phase velocity of planet/particle’s orbit velocity. Eq-38 clearly shows the physical meaning of \(\omega_{n,ph}\) and its relationship with \(v_{n,ph}\).

For a planet/particle moving in circular orbit with quantum number \(n\), eq-34 is always correct, or

\[\omega_n = \omega_{n,gr} = \frac{v_{n,gr}}{r_n} = \frac{v_n}{r_n}\]

where \(\omega_{n,gr}\) is the group angular frequency/velocity of planet/particle’s orbital movement. Eq-39 clearly shows the physical meaning of \(\omega_{n,gr}\) and its relationship with \(v_{n,gr}\). Notice here that values of planet/particle’s \(\omega_n = \omega_{n,gr}, v_n = v_{n,gr}\), and \(r_n\) are \(r_1\) position independent. Using eq-39 (or eq-34), we can re-write eq-31 as

\[\omega_{n,ph} = 2\pi f_n = \left(\frac{n}{2}\right) \frac{v_n}{r_n} = \left(\frac{n}{2}\right) \omega_{n,gr} = \left(\frac{n}{2}\right) \omega_n\]

\[eq-40\]
Notice here that values of \( n, f_n, \) and \( \omega_{n,\text{ph}} \) depend on where we choose \( r_1 \), while values of \( r_n, \) \( v_n \) and \( \omega_n \) are \( r_1 \) position independent. So eq-40 tells us that planet matter wave’s \( \phi \)-angular phase velocity \( \omega_{n,\text{ph}} \) equals to planet circular moving \( \phi \)-angular (group) velocity \( \omega_n \) times \( n/2 \). For example, after eq-27, we showed that Earth’s matter wave orbit frequency \( f_n \) at \( n=30 \) (or \( r_1 \) at \( \{0,1\} \)) is \( 4.44\times 10^{-7} \) cycle/sec = Hz, so its \( \omega_{30,\text{ph}} = 2\pi * f_n = 2.79\times 10^{-6} \) arc/s. Then, according to eq-40, \( \omega_n = \omega_{n,\text{ph}} * 2 \) \( \forall n\), which (closely) equals to Earth planet orbit’s \( \omega = 2\pi/(365.2 \text{ days}) = 1.99\times 10^{-7} \) arc/s (Note: the difference comes from that Earth’s orbit is a little bit deviated from \( \{1,5/6\} \)).

According to eq-35, my (and probably most readers’) original thought was that a planet matter wave’s angular (group) frequency/velocity must equal to \( 2 \times \) of planet’s (phase) angular frequency/velocity, or

\[
\omega_{n,\text{ph}} = \frac{\omega_n}{2} \quad \text{eq-41}
\]

But eq-41 is correct only under a special condition. This is because \( \omega_{n,\text{ph}} \) is a \( r_1 \) dependent value while \( \omega_n \) is a \( r_1 \) independent value. Only if we choose the \( r_1 \) position at planet’s orbit (that is, when forcing \( r_1 = r_n \), then eq-41 will exist (see statement-1 in section I-b). This is exactly the same meaning as that in eq-40-SunQM-3s10’s deduction where we chose \( k = 2\pi/\lambda = 1 \), or \( \lambda = 2\pi, \) or equivalent to in eq-17, we choose \( j=1 \). I emphasize it here because it is easily to be confused, and it had costed me many weeks to figure out.

I-d. For a planet doing circular orbital movement, the time-dependency of probability density \([T(t)]^2 \propto \exp(-\imath t \omega_{n,\text{ph}} t)]^2 = [\exp(-\imath t \omega_{n,\text{ph}} t)]^2 = \exp(-\imath t \omega_{n,\text{ph}} t)]^2 = \imath = 1 \) and define

\[
|T(t)|^2 \propto \left| e^{-\imath \omega_{n,\text{ph}} t} \right|^2 = e^{-\imath \omega_{n,\text{ph}} t} e^{-\imath \omega_{n,\text{ph}} t} = e^{-\imath 2 \omega_{n,\text{ph}} t} = e^{-\imath \omega_{n,\text{gr}} t} = e^{-\imath \omega_n t} \quad \text{eq-44}
\]

Then \([\Phi(\phi) \ast T(t)]^2 \) represents the probability density (of a planet’s matter wave packet) that is doing orbital movement in Solar system’s \( \phi \)-dimension, which equivalent to this planet is doing orbital movement in Solar system’s \( \phi \)-dimension. At \( n=1 \), eq-44 can also be written as

\[
|T(t)|^2 \propto \left( e^{-\imath \omega_{n,\text{ph}} t} \right)^2 = e^{-\imath \omega_{n=1,gr} t} \neq 1 \quad \text{eq-45}
\]

Using eq-9 and eq-44, we have
\[
|\Phi(\varphi)|^2 |T(t)|^2 \propto e^{im\varphi} e^{-i\frac{m}{2} \omega_n t} = e^{i \left\{ \left( \frac{1}{2} m - \frac{1}{2} n \right) \omega_n t \right\}} \left( \cos \left( \frac{\varphi - \frac{n}{m} \omega_n t}{2} \right) \right)^2 \to \left\{ \cos \left( \frac{\varphi - \frac{n}{m} \omega_n t}{2} \right) \right\}^m = \left[ 1 + \cos \left( \varphi - \frac{n}{m} \omega_n t \right) \right]^m \\
\]

where \( m = n - 1 \) because the whole deduction is based on that the planet is in \( |n \ell m> = |nLl> = |n, n > \) QM state. Notice that in eq-46 we again used eq-10 type (citizen-scientist level) method to alter a complex numbered probability to a real numbered probability. Readers should be able to use the strict math (as shown in eq-12 through eq-16) to obtain the same result. The last step in eq-46 is from the general trigonometric relationship in eq-11. Note: for the step-by-step explanation of the physical meaning of each item in eq-46, see SunQM-4s1.

So the final full-QM deducted planet’s time-dependent \( \varphi \)-dimensional (non-Born) probability density is

\[
\Phi(\varphi)|^2 |T(t)|^2 \propto \left[ 1 + \cos \left( \varphi - \frac{n}{n-1} \omega_n t \right) \right]^{n-1} \\
\]

Then, at \( n \gg 1 \), eq-47 becomes

\[
|\Phi(\varphi)|^2 |T(t)|^2 \propto \left[ 1 + \cos \left( \varphi - \frac{n}{2} \omega_n t \right) \right]^{n-1} \\
\]

I-e. Further modifying \( 0 \)-dimension’s probability density formula \( |\Theta(\theta)|^2 \) according to \( |\Phi(\varphi)|^2 * |T(t)|^2 \)

From eq-25-SunQM-3s11, we know that for a nLL QM state where \( m = +l = n-1 \), we can write

\[
\Theta(\theta) \propto \sin(\theta)^l = \sin(\theta)^{(n-1)} \\
\]

Then, because the definition of eq-9, and because the projection of a ball-shaped planet in Solar system’s \( \theta \varphi \)-2D dimension has to be a circular shape, we have to define

\[
\Theta(\theta) \propto \sin(\theta)^l = \sin(\theta)^{(n-1)} \\
\]

Notice that it is different than \( |\Theta(\theta)|^2 \propto [\sin(\theta)]^2 (2n-1) \) in eq-26-SunQM-3s11. Then, because eq-50 may have negative values under certain condition, and the probability density has to be a positive value, we need to alter eq-50 to be a forever positive form. Thus we used the same trick (a citizen-scientist level method) as that we have used in eq-10 for the transformation of eq-50:

\[
|\Theta(\theta)|^2 \propto \sin(\theta)^{(n-1)} = \left[ \cos \left( \frac{\pi}{2} - \theta \right) \right]^{(n-1)} = \left[ e^{i \left( \frac{\pi}{2} - \theta \right) / 2} \right]^{2(n-1)} = \left[ \cos \left( \frac{\pi}{2} - \theta \right) / 2 \right]^{2(n-1)} = \left[ \left[ \cos \left( \frac{\pi}{2} - \theta \right) / 2 \right]^{(n-1)} \right] = \left[ \left[ 1 + \cos \left( \frac{\pi}{2} - \theta \right) / 2 \right]^{(n-1)} \right] = \left[ \left[ 1 + \sin(\theta) \right] / 2 \right]^{(n-1)} \\
\]

or

\[
|\Theta(\theta)|^2 \propto \sin(\theta)^{(n-1)} = \left[ \sin \left( \frac{\pi}{2} - \theta \right) \right]^{(n-1)} = \left[ e^{i \left( \frac{\pi}{2} - \theta \right) / 2} \right]^{2(n-1)} = \left[ \sin \left( \frac{\pi}{2} - \theta \right) / 2 \right]^{2(n-1)} = \left[ \left[ \sin \left( \frac{\pi}{2} - \theta \right) / 2 \right]^{(n-1)} \right] = \left[ \left[ 1 + \sin(\theta) \right] / 2 \right]^{(n-1)} \\
\]
\[ |\Theta(\theta)|^2 \propto \left[ \frac{1 + \sin(\theta)}{2} \right]^{(n-1)} \quad \text{eq-52} \]

Thus eq-52 is the final full-QM deduced \( \theta \)-dimensional (non-Born) probability density for an orbital moving planet. Notice that the difference between eq-52 and eq-50 is that we simply lifted \( \sin(\theta) \) curve upward to \( [1 + \sin(\theta)] \) so that its original minimum (\( = -1 \)) value now equals to zero (or there is no negative value), and then normalized its maximum to 1 (by dividing 2 for \( [1 + \sin(\theta)] \)). That is the meaning of eq-11, and this is the standard formula of non-Born probability (NBP). In SunQM-4s1, we will see that NBP applies to many different formed QM.

Combining eq-47 and eq-52, we have the final time-dependent non-Born probability density for a planet in \( \theta \phi \)-2D dimension of Solar system:

\[ |Y(l, m)|^2 = |\Theta(\theta)|^2 |\Phi(\phi)|^2 |T(t)|^2 \propto \left[ \frac{1 + \sin(\theta)}{2} \right]^{(n-1)} \left[ 1 + \cos \left( \phi - \frac{n}{n-1} \omega_n t \right) \right]^{n-1} \quad \text{eq-53} \]

At \( n \gg 1 \), \( n/(n-1) \approx 1 \), eq-53 become

\[ |Y(l, m)|^2 \propto \left[ \frac{1 + \sin(\theta)}{2} \right]^n \left[ 1 + \cos \left( \phi - \omega_n t \right) \right]^n \quad \text{Eq-54} \]

Notice that eq-53 and eq-54 are valid only for nLL QM state. Figure 1 shows a 3D plot of a probability density peak (generated by using eq-54) moving in \( \theta \phi \)-2D-dimension.

Figure 1(a, b). 3D plot of a probability density peak (generated by using eq-54 at \( n=256 \)) moving in \( \theta \phi \)-2D-dimension. Figure 1a (left), with \( \omega_n * t = 0 \). Figure 1b (right), with \( \omega_n * t = 1 \). Using “3D Surface Plotter - An online tool to create 3D plots of surfaces” at: https://academo.org/demos/3d-surface-plotter/

I-f. Further modification of \( r \)-dimension’s probability density formula \( r^2 \cdot |R(r)|^2 \) according to \( |\Theta(\theta)|^2 \cdot |\Phi(\phi)|^2 \cdot |T(t)|^2 \)

For the same reason, because a planet has roughly same diameters in all 3D \( (r, \theta, \phi) \), the traditional QM’s \( r \)-dimensional probability density formula \( r^2 \cdot |R(n, l=n-1)|^2 \propto \frac{1}{r^2} \cdot \exp(1 - r/n) \cdot |(2\pi)^n (2\pi)^n \) (see eq-20-SunQM-3s11) will produce a planet with 50% of the diameter in \( r \)-dimension than that in \( \theta \) or \( \phi \)-dimension. Therefore, by comparing to eq-53, we are forced to define

\[ r^2 |R(n, l = n - 1)|^2 \propto \left[ \frac{r}{r_n} e^{\left( 1 - \frac{r}{r_n} \right)} \right]^n \quad \text{eq-55} \]

Since the result of eq-55 is always greater than zero, we do not need to make further modification (like we what did in eq-51).
I-g. The final full-QM deduced 3D (non-Born) probability density for a circular moving planet in Solar system

Combining eq-55 to eq-53, we obtain a planet’s

\[ r^2 |\Psi(r, \theta, \phi, t)_{\text{planet}}|^2 = r^2 |R(r)|^2 |\Theta(\theta)|^2 |\Phi(\phi)|^2 |T(t)|^2 \propto \left[ \frac{r}{r_n} e^{-\frac{(r-r_n)}{r_n}} \right]^n \left[ \frac{1+\sin(\theta)}{2} \right] \left[ \frac{1+\cos(\phi-n\omega r_\theta)}{2} \right]^{(n-1)} \]

(eq-56)

where \( r_n \) is planet’s orbital radius, and \( \omega_n \) is planet’s orbital angular frequency/velocity. **Eq-56 is the final full-QM deduced 3D (non-Born) probability density for a planet** (or for any particle in nLL QM state), and is valid for both base \( n \) and multiplier \( n' \). For a circular orbit moving planet (which is in nLL QM state), we should use its Eigen \( n' \) (for Eigen \( n' \) concept, see SunQM-3s10 section-IV. Also see SunQM-3s11 Table 1, where planet’s Eigen \( n'_s = n'_e = n'_o \)). Then eq-56 is an Eigen description for a circular orbit moving planet.

At \( n >> 1, n/(n-1) \approx 1 \), eq-56 become

\[ r^2 |\Psi(r, \theta, \phi, t)_{\text{planet}}|^2 \propto \left[ \frac{r}{r_n} e^{-\frac{(r-r_n)}{r_n}} \right]^n \left[ \frac{1+\sin(\theta)}{2} \right] \left[ \frac{1+\cos(\phi-n\omega r_\theta)}{2} \right] \]

(eq-57)

Both eq-56 and eq-57 are not only valid for a planet, but should also be valid for any object that is in nLL QM state and doing orbital movement (even for an electron). It is valid for both base \( n \) and multiplier \( n' \).

I-h. More discussions on eq-56 and its deduction

1) For eq-7, we have explained it as \( \exp(+i\phi) \) correlates to planet’s eastward orbital rotation’s \( \omega_{n,\text{ph}} \), while \( \exp(-i\phi) \) correlates to the westward orbital rotation’s \( \omega_{n,\text{ph}} \). Because for an eastward rotational planet, it can be thought as the westward orbital rotation’s \( \omega_{n,\text{ph}} = 0 \), so \( \exp(-i\phi) \propto \exp[i*\omega_{n,\text{ph,west}}*t] = \exp[i*0*t] = 1 \). So we have a good physical meaning for the \( \phi \)-dimension’s non-Born probability calculation in eq-9. For the physical meaning of \( \theta \)-dimension’s eq-50 and \( r \)-dimension’s eq-55, see SunQM-4s1.

2) When looking into the formula of eq-56, we are amazed by how simple the formula is and how straightforward the meaning it is: \( r^2 * |R(r,n,l=n-1)|^2 \propto \left[ r/r_n \right] \left[ \exp(1 - r/r_n) \right] \left[ \exp[i*\omega_{n,\text{ph,west}}*t] \right] \left[ \exp[i*0*t] \right] = 1 \). \( \text{eq-56} \) produces an exponential rising curve times an exponential declining curve, with the peak always at \( r = r_n \) and the higher the \( n \), the narrower the peak. \( |\Theta(\theta)|^2 \propto \left[ (1 + \sin(\theta)/2) \right]^{(n-1)} \) produces a peak at \( \theta = \pi/2 \), and the higher the \( n \), the narrower the peak. At \( n >> 1 \), \( |\Phi(\phi)|^2 \propto |T(t)|^2 \propto \left[ (1 + \cos(\phi - \omega n t))/2 \right] \) produces a peak at \( \phi = \omega_n t \), and the higher the \( n \), the narrower the peak.

3) Comparing to the classical physics in SunQM-3s11’s eq-45, the QM formed eq-56 has the angular frequency/velocity \( n/(n-1) * \omega_n = \omega_n + 1/(n-1)*\omega_n > \omega_n \). So there should be a positive precession for all circular orbital moving planets (if the time-dependent non-Born probability density is correct). At \( n \to 1 \), eq-47’s angular frequency/velocity \( n/(n-1)*\omega_n \to \infty \). In one of the future SunQM papers (probably SunQM-4s5), we will use this character to explain how RF is increased when a free particle is gradually trapped in a central (gravity or electric) force field. In another one of the future SunQM papers (probably SunQM-4s2), we will explore several possible applications of this positive precession in the Solar \( \{N,n\} \) QM.

II. To build a full-QM deduced non-Born probability density 3D map to describe the whole Solar system with time-dependent circular orbital movement
II-a. Full-QM deduced (Asteroid or the cold-KBO) belt’s 3D probability density $r^2 \left| \Psi(r, \theta, \phi, t)_{\text{belt}} \right|^2$

As mentioned before, eq-56 is valid for any object (in belt) that is in nLL QM state of the Solar \{N,n\} QM structure and doing orbital movement. But if the whole belt (or all mass of this belt) that is in nLL QM state (like Asteroid belt or the cold-KBO, see SunQM-3s10), then this belt’s orbital rotation (in $\phi$-dimension) can be described by using the alternated eq-46

$$|\Phi(\phi)_{\text{belt}}|^2 |T(t)|^2 \propto \left( e^{i m \phi/2} \right)^2 \left( e^{-i (\frac{m}{n}) \omega_n t} \right)^2 = e^{i \left[ \left( \frac{m}{n} \right) \omega_n t \right]} \left( \cos \left( \frac{m}{n} \left( \phi - \frac{n}{(n-1)} \omega_n t \right) \right) \right)^2$$

$$= 1 + \cos \left( (n-1) \left( \phi - \frac{n}{(n-1)} \omega_n t \right) \right)$$

Eq-58

where $\omega_n = \omega_{\text{rotation}} = v_n / r_n$ (same as eq-39), and $m=n-1$. When $n \gg 1$, $m=n-1$=n, eq-58 can be simplified as:

$$|\Phi(\phi)_{\text{belt}}|^2 |T(t)|^2 \propto \frac{1 + \cos(\phi - \omega_n t)}{2}$$

Eq-59

Eq-59 is valid for both base $n$ and multiplier $n'$. Figure 2 shows the plot of eq-59 with $n=1$ and $n=8$. We can see that at $n=1$, eq-59 in the form of $[1 + \cos(\phi - \omega_n t)]$ produces one probability density peak traveling eastward with $\omega_{\text{rotation}}$. At $n=8$, eq-59 in the form of $[1 + \cos(8(\phi - \omega_n t))]$ produces eight probability density peaks traveling eastward with $\omega_{\text{rotation}}$. Therefore, eq-59 is perfect for describing a belt that is made of $n$ pieces of equal mass, equal size and equal distance, and each piece is doing orbital movement with $\omega_{\text{rotation}}$. Then, after combining eq-58 with eq-52 and eq-55, we have $r^2 \left| \Psi(r, \theta, \phi, t) \right|^2 \propto r^2 \left| R(r) \right|^2 \left| \Theta(\theta) \right|^2 \left| \Phi(\phi)_{\text{belt}} \right|^2 \left| T(t) \right|^2$ for a belt as

$$r^2 \left| \Psi(r, \theta, \phi, t)_{\text{belt}} \right|^2 \propto \left( r / r_n \right)^n \left( 1 + \sin \theta / 2 \right)^{(n-1)} \left( 1 + \cos \left( \frac{(n-1)(\phi - \frac{n}{(n-1)} \omega_n t)}{2} \right) \right)$$

Eq-60

at $n \gg 1$, $m=n-1=n$, eq-58 can be simplified as:

$$r^2 \left| \Psi(r, \theta, \phi, t)_{\text{belt}} \right|^2 \propto \left( r / r_n \right)^n \left( 1 + \sin \theta / 2 \right)^n \left( 1 + \cos \left( \frac{n(\phi - \omega_n t)}{2} \right) \right)$$

Eq-61

where $n \gg 1$. Eq-60 (and eq-61) is the final full-QM deduced 3D probability density for a $\phi$-rotating belt. It can be used for Asteroid belt (with $n=48$) and the cold-KBO (with $n=192$).

Figure 2. Plot eq-59 at $n=1$ with $\omega_n t = 0$, at $n=8$ with $\omega_n t = 0$, and at $n=8$ with $\omega_n t = 1$. 
II-b. QM deduced $|\Phi(\phi)|^2 * |T(t)|^2$ orbital rotation for Oort cloud

For Oort cloud’s mass, if it is in $|nLL>$ QM state, then the orbital movement of probability density 3D map can be described by either eq-60 (or eq-61). If it is not in $|nLL>$ QM state, then its probability density’s $\phi$-dimensional macro movement is more complicated, but still can roughly be described by eq-60 (or eq-61). For simplicity, let’s use eq-61. Then, combining eq-61 to SunQM-3s11’s eq-17, we have

$$r^2 |\Psi(r, \theta, \phi, t)_{\text{OortCloud}}|^2 \propto r^2 \left[ |d_1 R(1 \times 5.33 \times 6^3, l) Y(l, m)|^2 \left\{ \frac{1 + \cos[n(\phi - \omega_4,1)]}{2} \right\} + |d_2 R(2 \times 5.33 \times 6^3, l) Y(l, m)|^2 \left\{ \frac{1 + \cos[n(\phi - \omega_4,2)]}{2} \right\} + |d_3 R(3 \times 5.33 \times 6^3, l) Y(l, m)|^2 \left\{ \frac{1 + \cos[n(\phi - \omega_4,3)]}{2} \right\} + |d_4 R(4 \times 5.33 \times 6^3, l) Y(l, m)|^2 \left\{ \frac{1 + \cos[n(\phi - \omega_4,4)]}{2} \right\} \right]$$

or eq-62

where $l = 0 \ldots n-1$, $m = -1 \ldots +1$, and $d_1 \ldots d_n$ are the linear combination coefficients, and $\omega_{4,1}$, $\omega_{4,2}$, $\omega_{4,3}$, $\omega_{4,4}$, and $\omega_{4,5}$ are the (averaged) angular frequencies/velocities for the orbital moving mass in $\{4, n=1..5/6\}$ orbital spaces. Notice that $n$ in eq-62 can be any value (as long as $>> 1$), here we can choose $n=6^4=1296$.

II-c. QM deduced $|\Phi(\phi)|^2 * |T(t)|^2$ for a self-spinning Sun

Sun’s self-spinning can also be roughly described by eq-59. At the lowest resolution, let’s suppose the Sun spins like a solid ball with a single angular frequency/velocity $\omega_{\text{sunSpin}}$. By combining eq-5-SunQM-3s11 and eq-61, we can have a QM deduced Sun’s probability density function with self-spin

$$r^2 |\Psi(r, \theta, \phi, t)_{\text{Sun}}|^2 \propto \left[ \text{eq.} - 5 - \text{SunQM} - 3s11 \right] \left\{ \frac{1 + \cos[n(\phi - \omega_{\text{sunSpin}})]}{2} \right\}$$

or eq-63

At a higher resolution, according to Solar $\{N,n\}$ QM theory, Sun core spins faster than its outer shell. Let’s assume that Sun’s $\{-1,1\}$ inner core has $\omega_{\text{sunSpin}}$, Sun’s $\{-1,n=1..5\}$ orbit shells has $\omega_{\text{sunSpin}}$, and Sun’s $\{0,1\}$ orbit shell has $\omega_{\text{sunSpin}}$, with $\omega_{\text{sunSpin1}} > \omega_{\text{sunSpin2}} > \omega_{\text{sunSpin3}}$. Since eq-5-SunQM-3s11 only includes Sun’s $\{-1,n=1..5\}$ orbit shells and $\{0,1\}$ orbit shell for the enough accuracy, so we don’t need to consider $\omega_{\text{sunSpin1}}$. Then, at a higher resolution, we can have a QM deduced self-spinning Sun’s probability density function

$$r^2 |\Psi(r, \theta, \phi, t)_{\text{Sun}}|^2 \propto r^2 \left[ |a_1 R\left(\frac{1}{6}, 0, Y(0,0)\right)|^2 + |a_2 R\left(\frac{2}{6}, l = 0.1, Y(l, m)\right)|^2 + |a_3 R\left(\frac{2}{6}, l = 0.1, Y(l, m)\right)|^2 \right] \left\{ \frac{1 + \cos[n(\phi - \omega_{\text{sunSpin}})]}{2} \right\} +$$

$$\left. + |a_4 R\left(\frac{4}{6}, l = 0.3, Y(l, m)\right)|^2 + |a_5 R\left(\frac{5}{6}, l = 0.4, Y(l, m)\right)|^2 \right\} \left\{ \frac{1 + \cos[n(\phi - \omega_{\text{sunSpin}})]}{2} \right\} +$$

$$r^2 \left[ \left\{ \frac{1 + \cos[n(\phi - \omega_{\text{sunSpin}})]}{2} \right\} \right]$$

or eq-64

where $l = 0 \ldots n-1$, $m = -1 \ldots +1$, and $a_1 \ldots a_6$ are the linear combination coefficients. Notice that $n$ in eq-64 can be any value (as long as $>> 1$), here we can choose $n=6^4=1296$. Notice that previous results showed that the convective zoom (see SunQM-3s8) and Ylm cycle (see SunQM-3s9) effects will overcome $\omega_{\text{sunSpin}}$ in some local regions.
II-d. To build a full-QM deduced non-Born probability density 3D map to describe the whole Solar system with time-dependent circular orbital movement

It has the similar form as that in SunQM-3s11’s eq-62, except that the probability density formulas of planet (or belt) are now following eq-56 (or eq-60). The Eigen n’ (from SunQM-3s11’s Table 1) is used in eq-56 for all planets so that it is the Eigen description of the 3D probability density for the orbital moving planets (meaning besides the orbit $r_n$, it also describes the planet’s size). Because all planets have extremely high Eigen n’ value (>1E+9), so we can use eq-57 (instead of eq-56). However, for n/(n-1) * $\omega_n$ in eq-56, we don’t know which n we should use. For example, for Earth at 1.5/6, if we use Earth planet’s Eigen n’ = 1.81E+9, then n/(n-1) = 1, or eq-57 is accurate. But if we use n’ = 5*6^1=30, then n/(n-1) = 30/29 = 1.03 will produce huge amount of positive procession. For this uncertainty, we use eq-57 but still retain the n/(n-1) * $\omega_n$ in eq-56. So the actual 3D probability density formula for all eight planets is:

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{planet}}^2 \propto \left( \frac{r}{r_n} e^{(1-\frac{r}{r_n})} \left[ 1 + \sin \theta \left( \frac{n + \cos \left( \frac{n - 1}{n - 1} \omega_n t \right)}{2} \right) \right]^n \right) $$

eq-65

where n’ is planet’s Eigen n’ (>1E+9), and n is one of planet’s multiplier n (uncertain which one).

Here are eight known planets’ full-QM deduced non-Born probability density formulas (all $r_n$, n’(s) and $\omega_n$) are obtained from SunQM-3s11’s Table 1, and the initial $\varphi$ position is set on Aug. 14, 2019:

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Mercury}}^2 \approx \left( \frac{r}{5.64 \times 10^{10}} e^{(1-\frac{r}{5.64 \times 10^{10}})} \left[ 1 + \cos \left( 1.08 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{1.09 \times 10^9} $$

eq-66

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Venus}}^2 \approx \left( \frac{r}{1.00 \times 10^{11}} e^{(1-\frac{r}{1.00 \times 10^{11}})} \left[ 1 + \cos \left( 3.14 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{1.45 \times 10^9} $$

eq-67

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Earth}}^2 \approx \left( \frac{r}{1.57 \times 10^{11}} e^{(1-\frac{r}{1.57 \times 10^{11}})} \left[ 1 + \cos \left( 0.16 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{1.81 \times 10^9} $$

eq-68

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Mars}}^2 \approx \left( \frac{r}{2.25 \times 10^{11}} e^{(1-\frac{r}{2.25 \times 10^{11}})} \left[ 1 + \cos \left( 3.33 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{1.31 \times 10^9} $$

eq-69

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Jupiter}}^2 \approx \left( \frac{r}{7.12 \times 10^{11}} e^{(1-\frac{r}{7.12 \times 10^{11}})} \left[ 1 + \cos \left( 5.31 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{7.26 \times 10^8} $$

eq-70

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Saturn}}^2 \approx \left( \frac{r}{1.60 \times 10^{12}} e^{(1-\frac{r}{1.60 \times 10^{12}})} \left[ 1 + \cos \left( 5.74 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{6.53 \times 10^9} $$

eq-71

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Uranus}}^2 \approx \left( \frac{r}{2.85 \times 10^{12}} e^{(1-\frac{r}{2.85 \times 10^{12}})} \left[ 1 + \cos \left( 1.27 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{5.22 \times 10^9} $$

eq-72

$$ r^2 |\Psi (r, \theta, \varphi, t)|_{\text{Neptune}}^2 \approx \left( \frac{r}{4.45 \times 10^{12}} e^{(1-\frac{r}{4.45 \times 10^{12}})} \left[ 1 + \cos \left( 0.52 \frac{n - 1}{n - 1} \omega_n t \right) \right] \right)^{6.53 \times 10^{10}} $$

eq-73

Following eq-60, we have the full-QM deduced non-Born probability density formulas (in Eigen description) for the two belts:
\[ r^2 |\Psi(r, \theta, \phi, t)_{\text{Asteroid Belt}}|^2 \propto \left( \frac{r}{4.01 \times 10^{11}} \right)^{48} \left( 1 + \sin(\theta) \right)^{47} \left( 1 + \frac{c_1}{2} \right) \text{eq-64} \]

\[ r^2 |\Psi(r, \theta, \phi, t)_{\text{KBO}}|^2 \propto \left( \frac{r}{6.40 \times 10^{12}} \right)^{192} \left( 1 + \sin(\theta) \right)^{191} \left( 1 + \frac{c_1}{2} \right) \text{eq-75} \]

where Asteroid belt's orbital rotation \( \omega_0 = 4.54E-8 \text{ arc/s} \), and Kuiper belt's \( \omega_0 = 7.11E-10 \text{ arc/s} \) are obtained from SunQM-3s11's Table 1. Notice that we are also not sure which multiplier \( n' (n' >> 1) \) is the correct one for \( n/(n-1) \) in eq-74 and eq-75.

For the four undiscovered \( \{3, n=2.5/6\} \) planets/belts, here we treat them as planets, so they have the same 3D probability density formula as eq-65:

\[ r^2 |\Psi(r, \theta, \phi, t)_{\{3,2\} \text{Planet}}|^2 \approx \left( \frac{r}{2.56 \times 10^{13}} \right)^{5.64 \times 10^{12}} \text{eq-76} \]

\[ r^2 |\Psi(r, \theta, \phi, t)_{\{3,3\} \text{Planet}}|^2 \approx \left( \frac{r}{5.76 \times 10^{13}} \right)^{5.08 \times 10^{13}} \text{eq-77} \]

\[ r^2 |\Psi(r, \theta, \phi, t)_{\{3,4\} \text{Planet}}|^2 \approx \left( \frac{r}{1.02 \times 10^{14}} \right)^{4.06 \times 10^{14}} \text{eq-78} \]

\[ r^2 |\Psi(r, \theta, \phi, t)_{\{3,5\} \text{Planet}}|^2 \approx \left( \frac{r}{1.60 \times 10^{14}} \right)^{5.08 \times 10^{14}} \text{eq-79} \]

where \( \phi_{\{N=3, n=2.5\}} \) are the (unknown) initial \( \phi \)-positions for the four undiscovered planets.

Now we can use a matrix production to constitute a complete probability density function for Solar \( \{N,n\} \) QM structure (as shown in eq-80). Notice that in comparison with SunQM-3s11’s eq-62, eq-80 not only produces a 3D map of probability density peaks for a complete Solar system (including a Sun, eight known planets, two known belts, four undiscovered planets, and Oort cloud (all in circular orbital movement)), but also includes Sun’s self-spin, and may even reflect planet’s orbital precession movement. Notice that in eq-80, the coefficient matrix is a diagonal only matrix, all non-diagonal cells have values of zero. The (most right) vector space column is composed by the probability density functions (in bold) with the equation numbers of eq-64, eq-66 through eq-79, and eq-62.
Similar as that in SunQM-3s11, here we name eq-80 as the “Solar $r^2 *|R(n,l)|^2 *|Y(l,m)|^2$ master matrix formula”. Eq-80 is the Eigen description (at planet or belt level) of our Solar system using Schrodinger equation’s solution. It can also be written as the integration form:

$$\text{Mass}(r, \theta, \phi, t) = 1.99 \times 10^{30} (\text{kg}) = \int_0^{1500 \times 1.49 \times 10^{11}} \int_0^{2\pi} \int_0^{\pi} \text{[eq - 80]} \times \sin(\theta) \, dr \, d\theta \, d\phi$$

**eq-81**

Therefore, each coefficient in eq-80’s diagonal matrix can be obtained because each item (Sun, planet, belt, cloud)’s integration should equal to this item’s mass.

III. More discussion on the Solar $r^2 *|R(n,l)|^2 *|Y(l,m)|^2$ master matrix formula eq-80

1) The calculated r-dimensional probability density distribution of eq-80 is shown in Figure 3. Notice that the previous low-resolution diagram of probability density r-distribution (in SunQM-3s1 Figure 3 where all eight planets’ probability density peak widths were very broad) is now updated to a high-resolution diagram (in which all eight planets’ probability density peak widths are close to planets’ true diameters). Check SunQM-3s11 section V for detailed explanation.

![Solar (N,n) QM model's probability density r-distribution from inside Sun to Oort cloud (as shown in eq-80)](image)

Figure 3. Probability density distribution in r-dimension calculated from eq-80 (where all eight planets’ probability density peak widths are close to planets’ diameters).

2) The real time-dependency of eight planets in $\theta=\pi/2$ plane (or x-y plane in Solar system) is illustrated in Figure 4 with the initial $\phi$-dimensional positions set on Aug. 14, 2019. Check SunQM-3s11 section V for detailed explanation.
Figure 4. Illustration of eq-80 generated eight planets in $\theta = \pi/2$ plane with initial $\phi$ positions on Aug. 14, 2019, and the $\phi$-ot positions after 60 days of orbital movement.

3) Can we use 3D probability density $r^2 \cdot |R(n,l)|^2 \cdot |Y(l,m)|^2$ map to calculate out the $\phi$-positions of the four undiscovered planets at $\{3,n=2...5/6\}$ orbits? (see SunQM-3s11 section VIII for detailed discussion).

4) A prediction that all mass entities (from the whole universe to a single quark) can be described by Schrodinger equation and solution (see SunQM-3s11 section IX for detailed discussion).

5) A (kind of) wrap-up discussion on the phase-1 study of Solar $\{N,n\}$ QM modeling (see SunQM-3s11 section X for details).

IV. The non-Born probability density method is valid (at least) for the macro movement in $\{N,n\}$ QM

This study revealed that at least for the macro movement in $\{N,n\}$ QM, the probability density function can be directly proportional to its matter wave function (here we named it as the non-Born probability density). More detailed explanations and evidences have been moved to the next paper SunQM-4s1. Although the non-Born probability (NBP) method was developed in 2019, I did not dare to publish it until I finished SunQM-4s1.

V. Eq-80 describes a Solar system not only at planet’s Eigen description level, but also at any level of resolution (down to proton level, or up to the whole universe level)! Therefore “Simultaneous-Multi-Eigen-Description (SMED)” is one of many nature attributes of QM

SunQM-3s10’s Figure 3 revealed that a Kuiper belt at $\{2,6\}$ can be described by $\{N,n\}$ QM at different levels of resolution. In current paper, with the new analysis result, let's further expand this idea with more examples. Example-1:
In SunQM-3s11 section-VII, we had an example to show that the cold-KBO at orbit \(3,1/6\) can be Eigen described as \([0,n=6^*5.33/6] = (0,192/6), or n=192, or in |192,191,192> QM state, with the radial wave function \(r^{1/2} \times |R(n=6^3,l=n-1)\|^2\) with \(r_1\) at \((0,1/6)\). It can also be described with a linear combination of radial wave functions \(r^{1/2} \times |R(n=6^18,l=n-1)\|^2\) with \(r_1\) at \((-15,1/6)\), which means many proton’s matter wave packets are running in the cold-KBO’s matter wave resonance chamber, or, the cold-KBO is made of many protons (and neutrons). Furthermore, it can also be described by radial wave function \(r^{1/2} \times |R(n=6^(-5),l=n-1)\|^2\) with \(r_1\) at \((8,1/6)\), which means the Milky way’s matter wave packet is also running in the cold-KBO’s matter wave resonance chamber.

Example-2:

Figure 5 demonstrated that a planet in Solar system can be described by \([N,n]\) QM’s r-dimensional probability density at different levels of resolution. At low resolution level, this planet can be described by a single \(r^{1/2} \times |R(n,l)\|^2\) curve which describes the whole Solar system’s \((N=1..4,n/6)\) region (as shown in Figure 5’s thick grey line). At median resolution level, this planet can be described by a single \(r^{1/2} \times |R(n,l)\|^2\) curve that describes only one N super-shell in Solar system’s \((N=1..4,n/6)\) region (as shown in Figure 5’s dot line) in which this planet is located. At median-high resolution level, this planet can be described by a single \(r^{1/2} \times |R(n,l)\|^2\) curve that describes only one n shell in Solar system’s \((N=1..4,n/6)\) region (as shown in Figure 5’s grey thin line) and this curve covers the whole \([N,n/6]\) orbit shell (r-dimensional) space from \(r_1\) to \(r_{n+1}\). At high resolution level, this planet can be described by a single \(r^{1/2} \times |R(n,l)\|^2\) curve at the Eigen n that describes not only the orbit \(r\) of this planet, but also the body size of this planet (as shown in Figure 5’s solid thin line).

Furthermore, at a very high resolution level, this planet can be described by a combination of many \(r^{1/2} \times |R(n,l)\|^2\) curves at the Eigen n that describing for a proton/neutron. So while each single \(r^{1/2} \times |R(n,l)\|^2\) describes not only the size of a single proton/neutron, but also the orbit \(r\) of this proton/neutron located (in Solar system and inside this planet), the combined \(r^{1/2} \times |R(n,l)\|^2\) describes not only the orbit \(r\), but also the body size of this planet (not shown in Figure 5). This example explains a planet in Solar system can be described by \([N,n]\) QM’s at different levels of resolution.

Example-3:

Here let us show another similar example: Earth at orbit \(1,5/6\) can be Eigen described as \([-10,n=5^*6^11//6] = [-10,1.81E9/6], or n=1.81E+11, or in |1.81E+11,n-1> QM state, with the radial wave function \(r^{1/2} \times |R(n=5^6^11,n=1.81E9,l=n-1)\|^2\) with \(r_1\) at \([-10,1/6]\), (see SunQM-3s11 Table 1). It can also be described with a linear combination of radial wave functions \(r^{1/2} \times |R(n=5^6^11,l=1.81E9,n=1-1)\|^2\) with \(r_1\) at \((-15,1/6)\), (notice that here \(l\) is no longer limited to be n-1, but includes n-2, n-3, …), which means many proton’s matter wave packets are running in Earth’s matter wave resonance chamber, or, the Earth is made of many protons (and neutrons). Furthermore, it can also be described by radial wave function \(r^{1/2} \times |R(n=5^6^11,(-7),l=n-1)\|^2\) with \(r_1\) at \((8,1/6)\), which means the Milky way’s matter wave packet is also running in the Earth’s matter wave resonance chamber.

Example-4:

Suppose Sun at size of \([0,2/6]\) contains (Sun’s mass / proton’s mass = 1.9885E+30 kg / 1.67E-27 kg) ~1.19E+57 of proton/neutrons. Also suppose our universe has a size of \([13,1/6]\) (see SunQM-1s2 Table 1). Then Sun can be Eigen described as \([0,2/6]\), or in [1,0] QM state with the radial wave function \(r^{1/2} \times |R(n=1,l=0)\|^2\) with \(r_1\) at \((0,1/6)\), (see SunQM-3s11 Table 1). It can also be described with a linear combination of radial wave functions \(r^{1/2} \times |R(n=1^6^15,l)\|^2\) with \(r_1\) at \((-15,1/6)\), (notice that here \(l\) is no longer limited to be n-1, but includes n-2, n-3, …) which means ~1.19E+57 of proton/neutron’s matter wave packets are running in Sun’s matter wave resonance chamber, or, the Sun is made of ~1.19E+57 protons/neutrons. Furthermore, it can also be described by radial wave function \(r^{1/2} \times |R(n=1^6^15,(-13),l=n-1)\|^2\) with \(r_1\) at \((13,1/6)\), which means that our universe’s matter wave packet is also running inside the Sun’s matter wave resonance chamber.

We can expand this explanation to a new concept: any object in our universe can be described by \([N,n]\) QM at different levels of resolution. So based on the \([N,n/6]\) QM structure model, eq-80 describes a solar system not only at the Eigen description (of a planet or a belt) level, but can be (simultaneously) at any possible levels of resolution (down to atom level, proton level, or up to galaxy level, even to the whole universe level)! In other words, the Eigen level is changeable in eq-80. In the current eq-80, the Eigen level is set to at Solar system’s planet/belt level. We can reset the Eigen level down to the hydrogen atom level (by setting \(r_1\) to the size of a H-atom at \((-12,1/6)\)), or to proton/neutron level (by setting \(r_1\) to the size of a proton at \((-15,1/6)\)), or to quark level (by setting \(r_1\) to the size of a quark at \((-17,1/6)\)), or up to galaxy level (by setting \(r_1\) to the size of a galaxy at \((8,1/6)\), or even to the whole universe level (by setting \(r_1\) to the size of our universe at
relative mass density

\[ \text{Conclusion} \]

Figure 5. Illustration of a planet in a Solar system can be described by \(\{N,n\}\) QM at different levels of resolution (copied from SunQM-3s11’s Figure 5a with some modifications). Note: these curves are artificially drawn to mimic the \(r^2|R(n,j)|^2\) calculated curve, they are not the true calculated result.
In this paper, we developed a non-Born probability density calculation for a planet in nLL QM state doing circular orbital movement: its $\phi$-dimensional probability density function is directly proportional to its matter wave function $|\Phi(\phi)|^2 \propto \Phi(\phi) \propto \exp(i\phi)$, and its time-dependency is also a non-Born calculation $|T(t)|^2 \propto \exp(-i \omega_{\text{phase}} t)$.

Then we calculated out (a full-QM deduced) $|\Phi(\phi)|^2 \times |T(t)|^2$ for a planet’s time-dependent probability density in $\phi$-dimension. We believe that by adding the non-Born calculation to Born calculation, the QM will become more self-consistent and more complete.

References

http://vixra.org/pdf/1805.0102v2.pdf  (original submitted on 2018-05-03)


[3] Yi Cao, SunQM-1s2: Comparing to other star-planet systems, our Solar system has a nearly perfect $\{N,n/6\}$ QM structure.

[4] Yi Cao, SunQM-1s3: Applying $\{N,n\}$ QM structure analysis to planets using exterior and interior $\{N,n\}$ QM.


[7] Yi Cao, SunQM-3s1: Using 1st order spin-perturbation to solve Schrodinger equation for nLL effect and pre-Sun ball’s disk-lyzation.

[8] Yi Cao, SunQM-3s2: Using $\{N,n\}$ QM model to calculate out the snapshot pictures of a gradually disk-lyzing pre-Sun ball.

[9] Yi Cao, SunQM-3s3: Using QM calculation to explain the atmosphere band pattern on Jupiter (and Earth, Saturn, Sun)’s surface.

[10] Yi Cao, SunQM-3s6: Predict mass density r-distribution for Earth and other rocky planets based on $\{N,n\}$ QM probability distribution.

http://vixra.org/pdf/1812.0302v2.pdf  (submitted on 2019-03-08)

[12] Yi Cao, SunQM-3s8: Using $\{N,n\}$ QM to study Sun’s internal structure, convective zone formation, planetary differentiation and temperature r-distribution.


[17] David J. Griffiths, Introduction to Quantum Mechanics, 2nd ed., 2015, pp61, free particle wave function, eq-2.94 and eq-2.95 combined, also see wiki “Schroedinger equation”


[20] A series of my papers that to be published (together with current paper):
SunQM-4: Using full-QM deduction and \{N,n\} QM’s non-Born probability density 3D map to build a complete Solar system with orbital movement.
SunQM-4s1: More explanations on the non-Born probability (NBP) in \{N,n\}QM.
SunQM-4s2: More explanations on non-Born probability (NBP)’s positive precession in \{N,n\}QM.
SunQM-4s3: Using \{N,n\} QM to analyze Earth’s atmosphere pattern and its effect on the weather.
SunQM-4s4: Schroedinger equation and \{N,n\} QM.
SunQM-4s7: Addendums, Updates and Q/A for SunQM series papers
SunQM-5: A new version of QM based on interior \{N,n\}, multiplier n’, |R(n,l)|^2 * |Y(l,m)|^2 guided mass occupancy, and RF, and its application from string to universe (drafted in April 2018).
SunQM-5s1: White dwarf, neutron star, and black hole re-analyzed by using the internal \{N,n\} QM (drafted in April 2018).

[21] Major QM books, data sources, software I used for this study:
Douglas C. Giancoli, Physics for Scientists & Engineers with Modern Physics, 4th ed. 2009.
Wikipedia at: https://en.wikipedia.org/wiki/
(Free) online math calculation software: WolframAlpha (https://www.wolframalpha.com/)
(Free) online spherical 3D plot software: MathStudio (http://mathstud.io/)
(Free) offline math calculation software: R
Microsoft Excel, Power Point, Word.
Public TV’s space science related programs: PBS-NOVA, BBC-documentary, National Geographic-documentary, etc.
Journal: Scientific American.