

A Fourier derivative collocation method for the solution of the Navier–Stokes problem

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August 29, 2024

A proposed solution to the millennium problem on the existence and smoothness of the Navier–Stokes equations.

1. Introduction

The Navier–Stokes equations are thought to govern the motion of a fluid in \mathbb{R}^3 , see [1–5]. Let $\mathbf{u} = \mathbf{u}(\mathbf{x}, t) \in \mathbb{R}^3$ be the fluid velocity and let $p = p(\mathbf{x}, t) \in \mathbb{R}$ be the fluid pressure, each dependent on position $\mathbf{x} \in \mathbb{R}^3$ and time $t \geq 0$. I take the externally applied force acting on the fluid to be identically zero. The fluid is assumed to be incompressible with constant viscosity $\nu > 0$ and to fill all of \mathbb{R}^3 . The Navier–Stokes equations can then be written as

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = \nu \nabla^2 \mathbf{u} - \nabla p, \quad (1)$$

$$\nabla \cdot \mathbf{u} = 0 \quad (2)$$

with initial condition

$$\mathbf{u}(\mathbf{x}, 0) = \mathbf{u}^\circ \quad (3)$$

where $\mathbf{u}^\circ = \mathbf{u}^\circ(\mathbf{x}) \in \mathbb{R}^3$. In these equations

$$\nabla = \left(\frac{\partial}{\partial \mathbf{x}_1}, \frac{\partial}{\partial \mathbf{x}_2}, \frac{\partial}{\partial \mathbf{x}_3} \right) \quad (4)$$

is the gradient operator and

$$\nabla^2 = \sum_{i=1}^3 \frac{\partial^2}{\partial \mathbf{x}_i^2} \quad (5)$$

is the Laplacian operator. When $\nu = 0$ equations (1), (2), (3) are called the Euler equations. When $\nabla p = 0$ equations (1), (3) are called the Burgers equations. Solutions of (1), (2), (3) are to be found with

$$\mathbf{u}^\circ(\mathbf{x} + e_i) = \mathbf{u}^\circ(\mathbf{x}) \quad (6)$$

for $1 \leq i \leq 3$ where e_i is the i^{th} unit vector in \mathbb{R}^3 . The initial condition \mathbf{u}° is a given C^∞ divergence-free vector field on \mathbb{R}^3 . A solution of (1), (2), (3) is then accepted to be physically reasonable [3] if

$$\mathbf{u}(\mathbf{x} + e_i, t) = \mathbf{u}(\mathbf{x}, t), \quad p(\mathbf{x} + e_i, t) = p(\mathbf{x}, t) \quad (7)$$

on $\mathbb{R}^3 \times [0, \infty)$ for $1 \leq i \leq 3$ and

$$\mathbf{u}, p \in C^\infty(\mathbb{R}^3 \times [0, \infty)). \quad (8)$$

2. Solution of the Navier–Stokes problem

Theorem. Take $\nu > 0$. Let \mathbf{u}° be any smooth, divergence-free vector field satisfying (6). Then there exist smooth functions \mathbf{u}, p on $\mathbb{R}^3 \times [0, \infty)$ that satisfy (1), (2), (3), (7), (8).

Proof. A Fourier derivative collocation method is as follows. Let \mathbf{u}, p be given by

$$\mathbf{u} = \sum_{\mathbf{L}=-\infty}^{\infty} \mathbf{u}_{\mathbf{L}} e^{i\mathbf{kL}\cdot\mathbf{x}}, \quad (9)$$

$$p = \sum_{\mathbf{L}=-\infty}^{\infty} p_{\mathbf{L}} e^{i\mathbf{kL}\cdot\mathbf{x}} \quad (10)$$

respectively. Here $\mathbf{u}_{\mathbf{L}} = \mathbf{u}_{\mathbf{L}}(t) \in \mathbb{C}^3$, $p_{\mathbf{L}} = p_{\mathbf{L}}(t) \in \mathbb{C}$, $i = \sqrt{-1}$, $k = 2\pi$, and $\sum_{\mathbf{L}=-\infty}^{\infty}$ denotes the sum over all $\mathbf{L} \in \mathbb{Z}^3$. The initial condition \mathbf{u}° is a Fourier series [2] of which is convergent for all $\mathbf{x} \in \mathbb{R}^3$. Equations (1), (2) can be written as

$$\frac{\partial \mathbf{u}_i}{\partial t} + \sum_{j=1}^3 \mathbf{u}_j \frac{\partial \mathbf{u}_i}{\partial \mathbf{x}_j} = \nu \sum_{j=1}^3 \frac{\partial^2 \mathbf{u}_i}{\partial \mathbf{x}_j^2} - \frac{\partial p}{\partial \mathbf{x}_i} \quad \text{for } i = 1, 2, 3, \quad (11)$$

and

$$\sum_{j=1}^3 \frac{\partial \mathbf{u}_j}{\partial \mathbf{x}_j} = 0 \quad (12)$$

respectively. In this method we have for a quantity q that

$$\left[\frac{\partial q}{\partial \mathbf{x}_j} \right] = [G_j] [q] \quad (13)$$

valid at $\mathbf{x} = \mathbf{x}_n^*$ for $n = 1, 2, \dots, N$. Here $[G_j]$ is a known constant $N \times N$ matrix with $[G_j]_{m,n} = G_{j,m,n}$ and $[r]$ means to vectorise r where the components are equal to $r|_{\mathbf{x}=\mathbf{x}_n^*}$, $n = 1, 2, \dots, N$. We denote $q|_{\mathbf{x}=\mathbf{x}_n^*} = [q]_n = q_n$. Then

$$\left[\frac{\partial \mathbf{u}_i}{\partial \mathbf{x}_j} \right]_n = \sum_{\alpha=1}^N G_{j,n,\alpha} u_{i,\alpha}, \quad \left[\frac{\partial p}{\partial \mathbf{x}_i} \right]_n = \sum_{\alpha=1}^N G_{i,n,\alpha} p_{,\alpha}, \quad (14)$$

$$\left[\frac{\partial^2 \mathbf{u}_i}{\partial \mathbf{x}_j^2} \right]_n = \sum_{\alpha=1}^N G_{j,n,\alpha} \left[\frac{\partial \mathbf{u}_i}{\partial \mathbf{x}_j} \right]_{\alpha} = \sum_{\alpha=1}^N \sum_{\beta=1}^N G_{j,n,\alpha} G_{j,\alpha,\beta} u_{i,\beta}, \quad (15)$$

and

$$\left[\frac{\partial \mathbf{u}_i}{\partial t} \right]_n = \frac{\partial}{\partial t} [\mathbf{u}_i]_n = \frac{\partial}{\partial t} u_{i,n}. \quad (16)$$

Equations (11), (12) at $\mathbf{x} = \mathbf{x}_n^*$ imply

$$\frac{\partial}{\partial t} [\mathbf{u}_i]_n + \sum_{j=1}^3 [\mathbf{u}_j]_n \left[\frac{\partial \mathbf{u}_i}{\partial \mathbf{x}_j} \right]_n = \nu \sum_{j=1}^3 \left[\frac{\partial^2 \mathbf{u}_i}{\partial \mathbf{x}_j^2} \right]_n - \left[\frac{\partial p}{\partial \mathbf{x}_i} \right]_n \quad (17)$$

and

$$\sum_{j=1}^3 \left[\frac{\partial \mathbf{u}_j}{\partial \mathbf{x}_j} \right]_n = 0 \quad (18)$$

respectively. Equations (17), (18) imply

$$\frac{\partial}{\partial t} u_{i,n} + \sum_{j=1}^3 \sum_{\alpha=1}^N u_{j,n} G_{j,n,\alpha} u_{i,\alpha} = \nu \sum_{j=1}^3 \sum_{\alpha=1}^N \sum_{\beta=1}^N G_{j,n,\alpha} G_{j,\alpha,\beta} u_{i,\beta} - \sum_{\alpha=1}^N G_{i,n,\alpha} P_{\alpha} \quad (19)$$

and

$$\sum_{j=1}^3 \sum_{\alpha=1}^N G_{j,n,\alpha} u_{j,\alpha} = 0 \quad (20)$$

respectively. Let U be a matrix where $U_{i,n} = u_{i,n}$ and let P be a matrix where $P_{\alpha,n} = p_{\alpha}$. Then equations (19), (20) imply

$$\frac{\partial U}{\partial t} + U(A(n)U) = \nu UB(n) - A(n)^T P \quad (21)$$

and

$$\text{trace}(UA(n)) = 0 \quad (22)$$

respectively. Herein $A(n)$ and $B(n)$ are matrices where

$$A(n)_{\alpha,j} = G_{j,n,\alpha} \quad (23)$$

and

$$B(n)_{\beta,n} = \sum_{j=1}^3 \sum_{\alpha=1}^N G_{j,n,\alpha} G_{j,\alpha,\beta}. \quad (24)$$

The i, n component of (21) recovers (19) since

$$\begin{aligned} [U(A(n)U)]_{i,n} &= \sum_{l=1}^N U_{i,l} [A(n)U]_{l,n} = \sum_{l=1}^N U_{i,l} \left[\sum_{m=1}^3 A(n)_{l,m} U_{m,n} \right] \\ &= \sum_{j=1}^3 \sum_{\alpha=1}^N U_{i,\alpha} A(n)_{\alpha,j} U_{j,n} = \sum_{j=1}^3 \sum_{\alpha=1}^N U_{i,\alpha} G_{j,n,\alpha} U_{j,n}, \end{aligned} \quad (25)$$

$$[UB(n)]_{i,n} = \sum_{l=1}^N U_{i,l}B(n)_{l,n} = \sum_{\beta=1}^N U_{i,\beta}B(n)_{\beta,n} = \sum_{j=1}^3 \sum_{\alpha=1}^N \sum_{\beta=1}^N U_{i,\beta}G_{j,n,\alpha}G_{j,\alpha,\beta}, \quad (26)$$

and

$$\begin{aligned} [A(n)^T P]_{i,n} &= \sum_{l=1}^N A(n)_{i,l}^T P_{l,n} = \sum_{l=1}^N A(n)_{l,i} P_{l,n} = \sum_{l=1}^N G_{i,n,l} P_{l,n} \\ &= \sum_{\alpha=1}^N G_{i,n,\alpha} P_{\alpha,n} = \sum_{\alpha=1}^N G_{i,n,\alpha} p_{\alpha}. \end{aligned} \quad (27)$$

Equation (22) recovers (20) since

$$\begin{aligned} \text{trace}(UA(n)) &= \sum_{j=1}^3 [UA(n)]_{j,j} = \sum_{j=1}^3 \sum_{l=1}^N U_{j,l}A(n)_{l,j} \\ &= \sum_{j=1}^3 \sum_{l=1}^N U_{j,l}G_{j,n,l} = \sum_{j=1}^3 \sum_{\alpha=1}^N U_{j,\alpha}G_{j,n,\alpha}. \end{aligned} \quad (28)$$

Let $Q(n)$ be a matrix such that $(A(n)^T P)Q(n) = 0$. We have

$$\begin{aligned} [(A(n)^T P)Q(n)]_{i,j} &= [A(n)^T (PQ(n))]_{i,j} = \sum_{l=1}^3 A(n)_{i,l}^T (PQ(n))_{l,j} \\ &= \sum_{l=1}^3 \sum_{m=1}^N A(n)_{i,l}^T P_{l,m} Q(n)_{m,j} = 0. \end{aligned} \quad (29)$$

Here nonzero $Q(n)$ is possible to construct because p at one arbitrary \mathbf{x} point is arbitrary and can be set to zero without loss of generality. Then (21) implies

$$\left(\frac{\partial U}{\partial t} \right) Q(n) + (U(A(n)U))Q(n) = (vUB(n))Q(n). \quad (30)$$

Equation (30) is the same as we would get for the Burgers equations. Now we consider a matrix Riccati equation problem.

$$\frac{\partial X}{\partial t} = aX + bY, \quad (31)$$

$$\frac{\partial Y}{\partial t} = cX + dY, \quad (32)$$

with

$$X = (U\lambda)Y. \quad (33)$$

Then we get

$$\left(\frac{\partial U}{\partial t}\lambda\right)Y + (U\lambda)\frac{\partial Y}{\partial t} = a(U\lambda)Y + bY \quad (34)$$

which implies

$$\left(\frac{\partial U}{\partial t}\lambda\right)Y + (U\lambda)[c(U\lambda)Y + dY] = a(U\lambda)Y + bY \quad (35)$$

implying

$$\frac{\partial U}{\partial t}\lambda + (U\lambda)c(U\lambda) + (U\lambda)d = a(U\lambda) + b. \quad (36)$$

We then let $a = b = 0$, $\lambda = Q(n)$, $c = Q(n)^{-1}A(n)$, $d = -\nu Q(n)^{-1}B(n)Q(n)$ to recover (30). Then (31) implies

$$X = X|_{t=0}. \quad (37)$$

Equation (32) implies

$$\frac{\partial Y}{\partial t} = cX|_{t=0} + dY \quad (38)$$

and so

$$\frac{\partial}{\partial t}(e^{-dt}Y) = e^{-dt}cX|_{t=0} \quad (39)$$

which integrating with respect to t yields

$$e^{-dt}Y = \int_0^t e^{-d\tau}cX|_{t=0}d\tau + Y|_{t=0} \quad (40)$$

which implies

$$e^{-dt}Y = \left[e^{-d\tau}(-d)^{-1}cX|_{t=0}\right]_0^t + Y|_{t=0} \quad (41)$$

to obtain

$$Y = (e^{-dt})^{-1} \left\{ (e^{-dt} - I)(-d)^{-1}cX|_{t=0} + Y|_{t=0} \right\}. \quad (42)$$

Equation (33) then implies

$$\begin{aligned} U\lambda &= X|_{t=0}Y^{-1} \\ &= (U|_{t=0}\lambda)Y|_{t=0} \left[\left\{ (e^{-dt} - I)(-d)^{-1}c(U|_{t=0}\lambda)Y|_{t=0} + Y|_{t=0} \right\}^{-1} e^{-dt} \right] \\ &= (U|_{t=0}\lambda) \left[\left\{ (e^{-dt} - I)(-d)^{-1}c(U|_{t=0}\lambda) + I \right\}^{-1} e^{-dt} \right]. \end{aligned} \quad (43)$$

No blowup is possible since the Burgers equations are regular. \square

For the Euler equations we have

$$U\lambda = (U|_{t=0}\lambda) \{c(U|_{t=0}\lambda)t + I\}^{-1} \quad (44)$$

and blowup is possible since for odd N the equation

$$\det (c(U|_{t=0}\lambda)t + I) = 0 \tag{45}$$

can have a solution t where $0 < t < \infty$.

References

- [1] Batchelor G. 1967. *An introduction to fluid dynamics*. Cambridge U. Press, Cambridge.
- [2] Doering C. 2009. The 3D Navier–Stokes problem. *Annu. Rev. Fluid Mech.* **41**: 109–128.
- [3] Fefferman C. 2000. Existence and smoothness of the Navier–Stokes equation. *Clay Mathematics Institute*. Official problem description.
- [4] Ladyzhenskaya O. 1969. *The mathematical theory of viscous incompressible flows*. Gordon and Breach, New York.
- [5] Tao T. 2013. Localisation and compactness properties of the Navier–Stokes global regularity problem. *Analysis and PDE*. **6**: 25–107.